

Problem set 6

Problem 1. Drude model

$$\langle \dot{p}(t) \rangle = -\gamma \langle p(t) \rangle + f(t) \quad , \quad \text{apply Laplace transform on both sides, } \frac{\dot{\bar{p}}(t)}{\bar{p}(t)} + \gamma \bar{p}(t) = f(t) \quad ,$$

$$\text{L.H.S} = s \cdot \bar{p}(s) - \bar{p}(0) + \gamma \cdot \bar{p}(s) = \text{RHS} = f(s)$$

In s-domain, EOM is solved by $\bar{p}(s) = \frac{1}{s+\gamma} (\bar{p}(0) + f(s)) \xrightarrow{\mathcal{L}^{-1}} \bar{p}(t) = e^{-\gamma t} \bar{p}(0) + \int_0^t dt e^{-\gamma(t-\tau)} f(\tau)$

Let $f(t)$ be constant f ,

$$\bar{p}(t) = e^{-\gamma t} \bar{p}(0) + \left(-\frac{f}{\gamma}\right) e^{-\gamma(t-\tau)} \Big|_0^t = e^{-\gamma t} \bar{p}(0) + \frac{f}{\gamma} (1 - e^{-\gamma t})$$

When $t \gg \gamma^{-1}$ one gets steady condition $\bar{p} = \frac{f}{\gamma} \quad \gamma = \tau^{-1} \quad J = \frac{ne^2 \tau}{me} E$

Let $f(t)$ be $f(t) = f \cdot e^{i\omega t}$, $f(s) = f \cdot \frac{1}{s-i\omega}$

$$\bar{p}(s) = \frac{1}{s+\gamma} \bar{p}(0) + \frac{f}{(s+\gamma)(s-i\omega)}$$

$$\bar{p}(t) \xrightarrow{\mathcal{L}^{-1}} e^{-\gamma t} \bar{p}(0) + \frac{f}{i\omega + \gamma} (e^{i\omega t} - e^{-\gamma t})$$

$$\text{Let } t \gg \gamma^{-1}, \quad \bar{p}(t) = \frac{f(t)}{i\omega + \gamma} \quad \rightarrow \quad J(\omega) = \frac{ne^2 \tau}{me} \cdot \frac{i\omega \tau + 1}{i\omega} E(\omega)$$

* complex $\sigma_0(\omega)$ means a dynamical response to the oscillatory field. (current oscillating \sim field oscillating)

* when $\omega \gg \tau^{-1}$, $\sigma_0(\omega)$ becomes purely imaginary \rightarrow only oscillation no damping
 \rightarrow high frequency oscillation prevents collision

Problem 2. Hellmann-Feynman theorem

$$\hat{H}_\lambda |\psi_\lambda\rangle = E_\lambda |\psi_\lambda\rangle, \quad E_\lambda = \langle \psi_\lambda | \hat{H}_\lambda | \psi_\lambda \rangle \quad (\langle \psi_\lambda | \psi_\lambda \rangle = 1)$$

$$\begin{aligned} \frac{\partial E(\lambda)}{\partial \lambda} &= \left\langle \frac{\partial}{\partial \lambda} \psi_\lambda \middle| \hat{H}_\lambda \middle| \psi_\lambda \right\rangle + \langle \psi_\lambda | \frac{\partial}{\partial \lambda} \hat{H}_\lambda | \psi_\lambda \rangle + \langle \psi_\lambda | \hat{H}_\lambda \middle| \frac{\partial}{\partial \lambda} \psi_\lambda \rangle \\ &= E_\lambda \left\langle \frac{\partial}{\partial \lambda} \psi_\lambda \middle| \psi_\lambda \right\rangle + \langle \psi_\lambda | \frac{\partial}{\partial \lambda} \hat{H}_\lambda | \psi_\lambda \rangle + E_\lambda \langle \psi_\lambda | \hat{H}_\lambda | \psi_\lambda \rangle \\ &= E_\lambda \frac{\partial}{\partial \lambda} \langle \psi_\lambda | \psi_\lambda \rangle + \langle \psi_\lambda | \frac{\partial}{\partial \lambda} \hat{H}_\lambda | \psi_\lambda \rangle \\ &\quad \uparrow \\ &\quad 0 \end{aligned}$$

Problem 3 Van Hove Singularities

$$E_{\vec{k}} = -2t \sum_{\mu=1}^d \cos(k_{\mu}a)$$

$$dE = \frac{\partial E}{\partial k_x} dk_x + \frac{\partial E}{\partial k_y} dk_y + \frac{\partial E}{\partial k_z} dk_z = \vec{\nabla} E \cdot d\vec{k} = \vec{\nabla} E \cdot (dk_{\perp} + dk_{\parallel}) = |\nabla E| dk_{\perp}$$

$$N = \int_{1.BZ} g(\vec{k}) d^3k = \frac{N}{\Omega_r} \int_{1.BZ} d^3k = \frac{N}{\Omega_r} \int_{1.BZ} dk_{\perp} \int_{\partial \Omega} dS$$

$$= \frac{N}{\Omega_r} \int dE \int_{S_E} \frac{dS}{|\nabla E|} = \int dE D(E) \rightarrow D(E) = \frac{N}{\Omega_r} \int_{S_E} \frac{dS}{|\nabla E|}$$

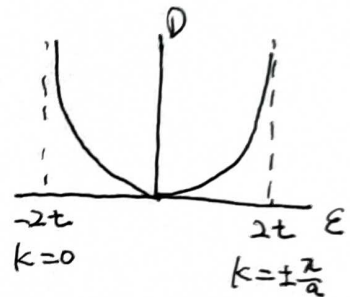
1-d case

$$\frac{\partial}{\partial k} E = 2at \sin(ka) \xrightarrow{H-F} \langle v_k \rangle$$

$$= 0 \Rightarrow k = 0, \pm \frac{\pi}{a}$$

$$|\nabla E| = 2at |\sin(ka)| = 2at \sqrt{1 - \cos^2 ka} = 2at \sqrt{1 - (E/2t)^2}$$

$$D(E) \propto [1 - (E/2t)^2]^{-\frac{1}{2}}$$



2-d case

$$E(k_x, k_y) = -2t [\cos(k_x a) + \cos(k_y a)]$$

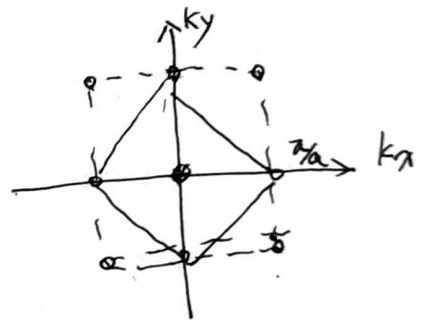
$$\frac{\partial}{\partial k_x} E = 2t \sin(k_x a), \quad \frac{\partial}{\partial k_y} E = 2t \sin(k_y a)$$

$$\vec{\nabla} E = 0 \rightarrow k_x = 0, \pm \frac{\pi}{a}, \quad k_y = 0, \pm \frac{\pi}{a}$$

$$\frac{\partial^2}{\partial k_x^2} E = 2a^2 t \cos(k_x a), \quad \frac{\partial^2}{\partial k_y^2} E = 2a^2 t \cos(k_y a)$$

$$\text{maximum } (0, 0) \frac{\partial^2}{\partial k^2} = 2a^2 t, \quad \text{minimum } (\frac{\pi}{a}, \frac{\pi}{a}) \frac{\partial^2}{\partial k^2} = -2a^2 t$$

$$\text{saddle point } (\frac{\pi}{a}, 0), \frac{\partial^2}{\partial k^2} = \pm 2a^2 t$$



vicinity of $(0,0)$, $x = \delta k_x$, $y = \delta k_y$, $A = 2at^2$

$$\epsilon = \epsilon_{\min} + \frac{1}{2}A(x^2 + y^2)$$

$$\frac{\delta \epsilon}{A} = \frac{x^2 + y^2}{2} \quad \delta \epsilon = \epsilon - \epsilon_{\min} = \frac{1}{2}A(x^2 + y^2) \rightarrow \begin{cases} x = R \cdot \cos \theta \\ y = R \cdot \sin \theta \end{cases} \quad \delta \epsilon = \frac{1}{2}AR^2$$

$$\frac{\partial \epsilon}{\partial x} = Ax, \quad \frac{\partial \epsilon}{\partial y} = Ay, \quad |\nabla \epsilon| = A \cdot \sqrt{x^2 + y^2} = A \cdot R$$

$$D(\epsilon) = \int_{\delta \epsilon} \frac{ds}{|\nabla \epsilon|} = \frac{2\pi R}{A \cdot R} = \frac{2\pi}{A} \rightarrow \text{constant}$$

vicinity of $(\frac{\pi}{a}, \frac{\pi}{a})$, $\epsilon = \epsilon_{\max} - \frac{1}{2}A(x^2 + y^2)$, $-\delta \epsilon = \epsilon - \epsilon_{\max} = -\frac{1}{2}A(x^2 + y^2)$

Similarly, $D(\epsilon) = \text{constant}$.

vicinity of saddle point $(\frac{\pi}{2a}, \frac{\pi}{2a})$ $(0, \frac{\pi}{2a})$

$$\delta \epsilon = \epsilon = 0 + \frac{1}{2}A(x^2 - y^2) \rightarrow \begin{cases} x = R \cdot \cosh \alpha \\ y = R \cdot \sinh \alpha \end{cases}$$

$$\frac{\partial \epsilon}{\partial x} = Ax, \quad \frac{\partial \epsilon}{\partial y} = -Ay, \quad |\nabla \epsilon| = AR \sqrt{\cosh^2 \alpha + \sinh^2 \alpha}$$

$$ds = R d\alpha \sqrt{\cosh^2 \alpha + \sinh^2 \alpha}$$

$$D(\epsilon) \propto \int_{-\alpha_c}^{+\alpha_c} d\alpha = 2\alpha_c$$

Consider our truncation on $[-\alpha_c, +\alpha_c]$

$$\alpha_c = R \cdot \cosh \alpha_c = \frac{R}{2}(e^{\alpha_c} + e^{-\alpha_c})$$

For large enough α_c , $\alpha_c \approx \frac{R}{2} e^{\alpha_c}$ so $\alpha_c \approx \ln \frac{2\alpha_c}{R}$

$$\text{Note that } R^2 = \frac{2\delta \epsilon}{A}$$

$$D(\epsilon) \propto 2 \times \frac{1}{2} \ln \frac{2\alpha_c^2 A}{\delta \epsilon} = \ln(2\alpha_c^2 A) - \ln \delta \epsilon$$

Hence $D(\epsilon)$ diverge logarithmically