

GUT Course 22/23

Lecture XVII

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23/12 /2022

LMU  
Winter 2022

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$SU(5)$  $\downarrow \langle 24_H \rangle$  $G_{SM}$  $\downarrow \langle 5_H \rangle$  $U(1)_{em} \times SU(3)_c$ 

but if there was only  $5_H$   
 $\Rightarrow SU(4)$  at the end

$$\cancel{SU(3) \sim SU(2) \times U(1)}$$

$\delta_{\text{gen}}$

$(3+1)_{\text{gen}}$

$$SU(n)$$

$$\downarrow \langle F_H \rangle = N_H$$

$$SU(N-1) \times U(1)$$

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$$SU(5) \quad (\tau=4)$$

maximal  
subgroups

$$SU(3) \times SU(2) \times U(1) \\ (\tau=4)$$

$$SU(4) \times U(1) \\ (\tau=4)$$

but

$$\langle \begin{pmatrix} 5_H \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ -\frac{1}{e_W} \end{pmatrix} \} \quad \begin{matrix} SU(4) \\ good \end{matrix}$$

+

$$\mathcal{L}_Y = f(\bar{5}_F, 10_F, 5_H)$$

∴

$$\langle 5_H \rangle \neq 0 \Rightarrow$$

$SU(4)$  inv. velocities!

$m_d = m_e$

Add  $d > 4$  terms

$$\Delta \mathcal{L}_Y = \bar{5}_F^i 10_F j 5_H^{*k} \frac{(24_H)^j}{\Lambda}$$

$$\Lambda \gg M_{GUT} = \langle 24_H \rangle$$

check:

$$(1) \quad \bar{5}_F^i 10_{Fij} 5_H^{*k} \frac{(24_H)_k^j}{\Lambda}$$

$$\rightarrow \bar{5}_F^i 10_{Fij} 5_H^{*k} \frac{\langle 24_H \rangle_u^j}{\Lambda}$$

?

breaks  $SU(4)_C$

$$= \bar{5}_F^i 10_{Fij} {5_H^*}^\mu \frac{V_j S_\mu^j}{\Lambda}$$

$$\langle 24_h \rangle = \text{diag } V_{\text{GUT}} (1, 1, 1, -\frac{3}{2}, -\frac{3}{2})$$

$$(\alpha=1,3) \quad V_\alpha = V_{\text{GUT}}$$

$$(\alpha=4,5) \quad V_\alpha = -\frac{3}{2} V_{\text{GUT}}$$

↓

$$\Delta S_y = \bar{5}_F^i 10_{Fiy} {5_H^*}^\mu \frac{V_{\text{GUT}} S_\mu^y}{\Lambda}$$

$$\langle {5_H^*}^\mu \rangle^\mu = 0$$

$$+ \bar{S}_F^1 10_{F12} S_H^{*^b} (-\frac{3}{2} V_{OUT}) \delta_L^q$$



$\Lambda$

$$\boxed{\bar{S}_F^i 10_{Fis} \frac{v_w (-\frac{3}{2} V_{OUT})}{\Lambda}}$$



$\Lambda$

$$-\frac{3}{2} \frac{V_{OUT}}{\Lambda} \left\{ \begin{array}{l} \bar{S}_F^\alpha 10_{F\alpha S} + \\ + \bar{S}_F^4 10_{F45} + \\ + \cancel{\bar{S}_F^5 10_{F55}} \end{array} \right\} v_w$$

$$= -\frac{3}{2} \left( \frac{V_{OUT}}{\Lambda} \right) v_w \left\{ d^c c^T c d + e^T c e^c \right\}$$

↑

keep our relations  
 $m_\delta = m_e$

but also

$$(2) \bar{5}_F^i 24_H^k 10_{Fkj} 5_H^{*j} \frac{1}{\Lambda}$$



$$\rightarrow \bar{5}_F^i \langle 24_H \rangle_k^i 10_{Fkj} 5_H^{*j} \frac{1}{\Lambda}$$



breaks  $SU(4)$  accidental



 transfer \$SU(4)\$ baryons  
 to \$d^c\$ and \$e\$ in \$\bar{S}\_F\$

$$\bar{S}_F^i \bar{V}_k S_i^k 10_{Fkj} S_H^{*j} \frac{1}{\Lambda}$$

$$\langle S_u \rangle^* = v_w \epsilon_{j5}$$

$$\rightarrow \bar{S}_F^k \bar{V}_h 10_{Fkj5} v_w \frac{1}{\Lambda}$$

$$\Rightarrow d^c d (v_{aut}/\Lambda v_w) +$$

$$e e^c (-\frac{3}{2} v_{aut}/\Lambda v_w)$$

↑

splits up, we

↓

no relation between  
 $m_d$  and  $m_e$

if  $\lambda = \mu_p \simeq 10^{19}$  GeV

$$\Rightarrow \left( \frac{V_{0.0T}}{\lambda} \right) \simeq 10^{-3}$$

$$+ \Delta Z_F = 10_F 10_F 5_F \frac{2^L H}{\lambda} \quad \checkmark$$

Suggestion: stick in  
indices

Step back

$$L_y = \overline{5}_F \gamma_1 10_F \overline{5}_H^* \quad (d=4)$$

$$\delta L_y = \overline{5}_F 10_F \overline{5}_H^* \frac{24_H}{\Lambda}$$



$$\boxed{\overline{5}_H^* 24_H}$$

$$\overline{5}^* {}^i 24_H^j$$

→ Irreducible  
 $(H_{\mu\nu})^{ij}_{\alpha\beta} \rangle$

$$\rightarrow (H_{\text{new}})^{[ij]}_k - T_V(H_{\text{new}})$$

↗

$$(H_{\text{new}}^{ij})_{(5)}$$

$$[ij] = (AS)_{ij}$$

$$\Rightarrow [ij] = \frac{5 \cdot 4}{2} = 10$$

$$\Rightarrow 50^{[ij]}_u - 5^{[ij]}_i =$$

$= 45 \text{ components}$

Instead, let's stick

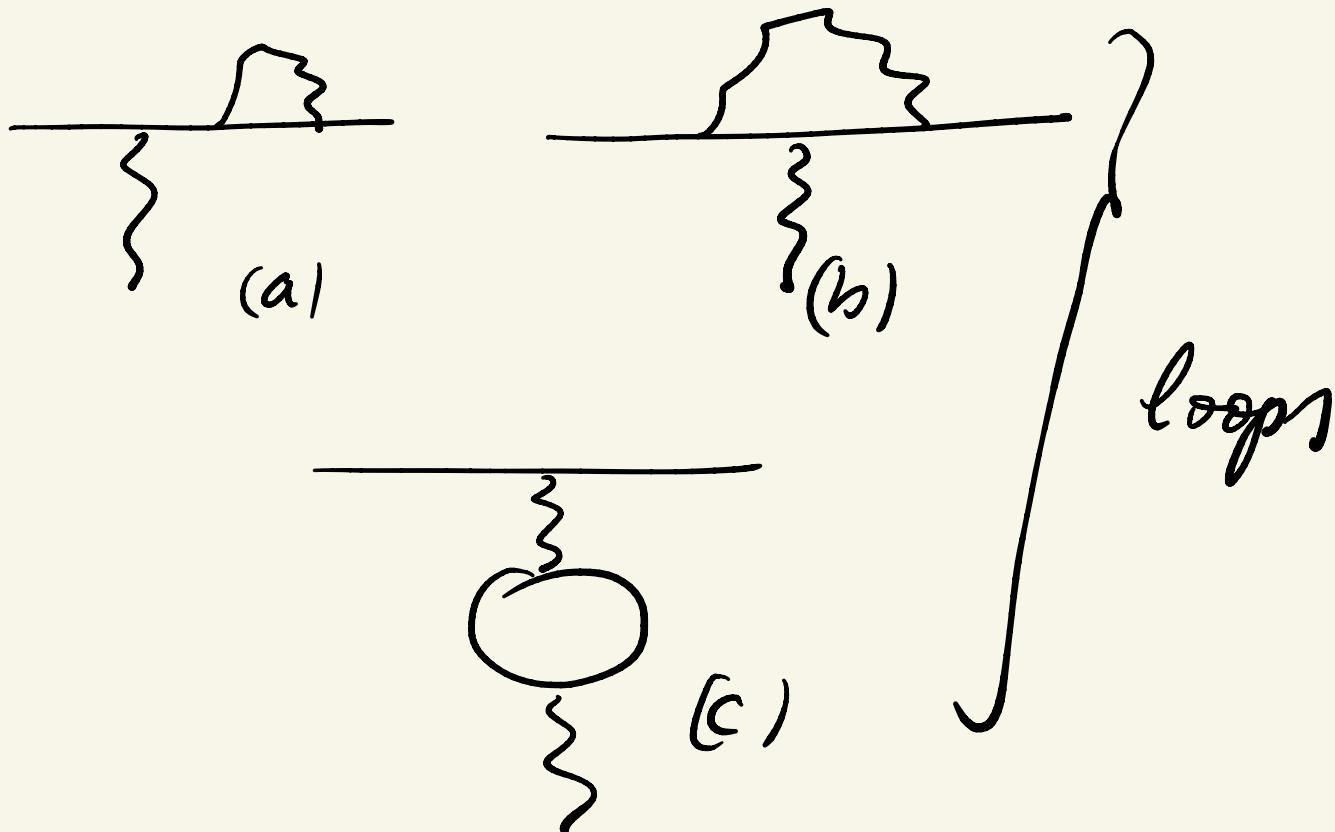
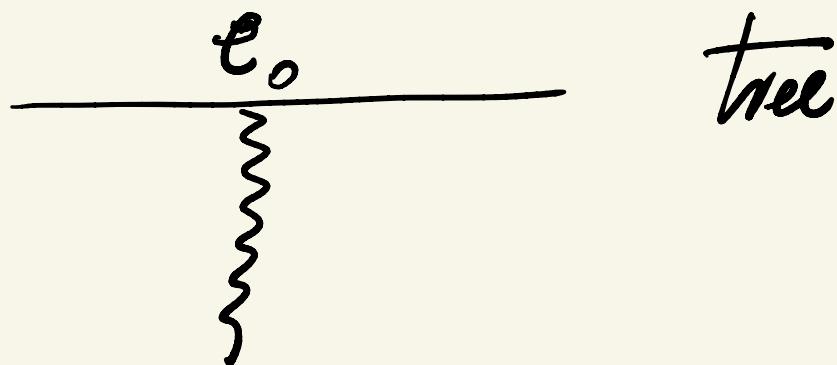
good old minimal  $SO(5)$

(+  $d=5 \Delta L_y$ )

# UNIFICATION of

## GAUGE COUPLINGS

QED



$$e(1) = e^0$$

$$e(E) = e_0 - \frac{e_0^3}{16\pi^2} \ln \frac{1}{E} \quad \leftarrow$$

$$(b) \quad \alpha \int \frac{d^4 h}{(2\pi)^4} \frac{1}{h^2} \frac{1}{\lambda} \frac{1}{\lambda}$$

$$\alpha \int \frac{d^4 h}{h^4} \propto \ln 1$$

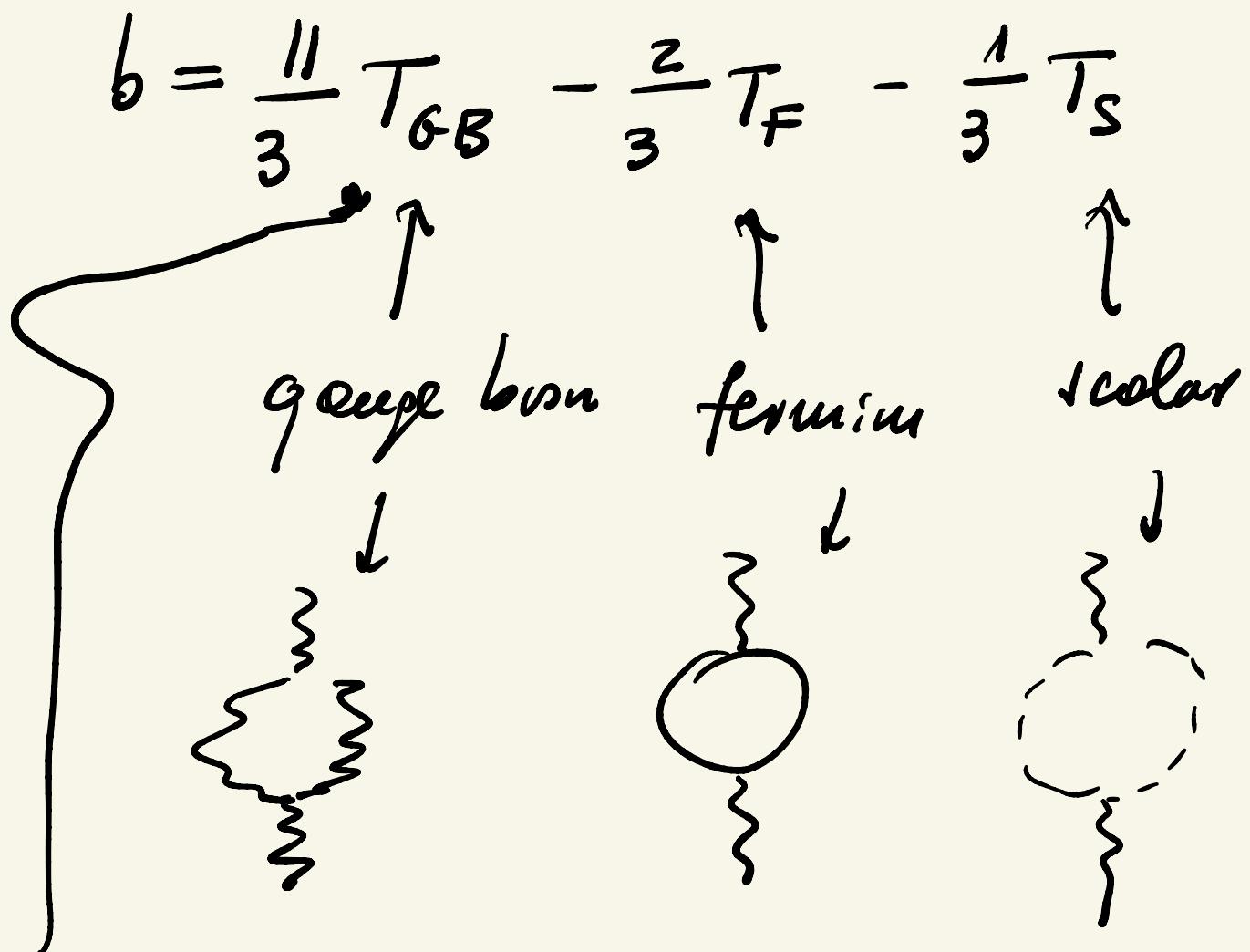
$$e(E_2) = e(E_1) + \frac{e^3}{16\pi^2} \ln \frac{E_2}{E_1}$$

$\begin{array}{c} + \\ \dots - \end{array} \rightarrow$  wave charge  
large  $E$   $\begin{array}{c} - \\ + \\ \dots - \end{array}$

↓  
general formula

$$\frac{1}{\alpha(E_2)} = \frac{1}{\alpha(E_1)} + \frac{b}{2\pi} \ln \frac{E_2}{E_1}$$

$$b = \frac{11}{3} T_{GB} - \frac{2}{3} T_F - \frac{1}{3} T_S$$


  
 gauge boson      fermion      scalar

} gauge boson  $\Rightarrow$   
 Asymptotic Freedom (AF)

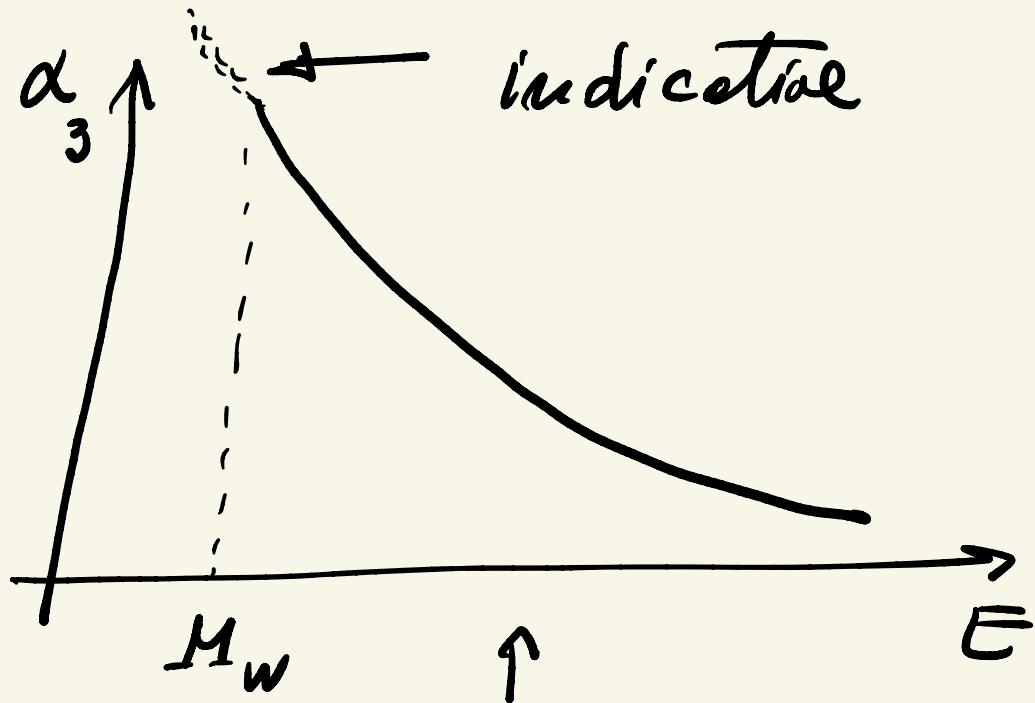
$$\overline{T} \delta^{ab} = T_r T_a T_b$$

recall:  $T_{fund} = \overline{T}_f$  (fundamental  
 repn.)

$$T_r \overline{T}_f^a \overline{T}_f^b = \frac{1}{2} \delta^{ab}$$

•  $b > 0$  (q.b. there)

$\Rightarrow$  AF



example : QCD  $\leftrightarrow$  SU(3)

$$\alpha_3(M_w) \simeq 1/10 \quad \text{SU}(3)_c$$

$$\alpha_2(M_w) \simeq 1/30 \quad \text{SU}(2)_L$$

$$v_{ew}(M_w) \simeq 1/130 \quad \text{U}_{\text{color}} v_{ew}$$

Deep Inelastic  
Scattering

(1969)

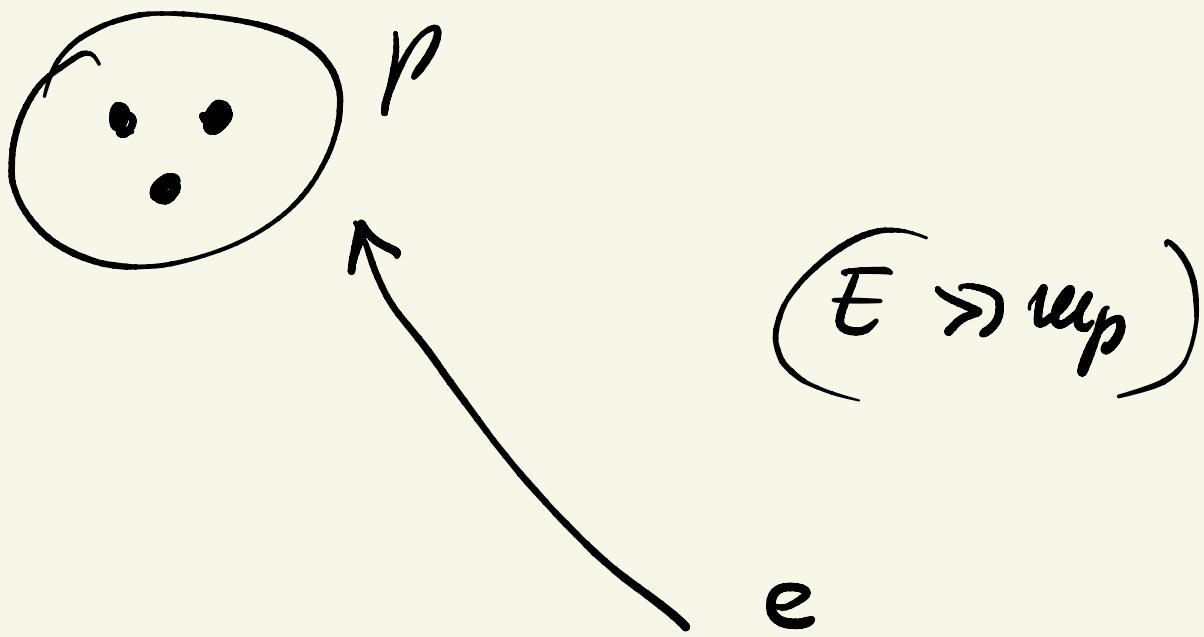
3 quarks  $\leftarrow$  Jell - Mann 1964

(aces)  $\leftarrow$  \* \* Zweig 1964

1967 - 1969

Bjorken

↙ Rutherford  
picture



$$e + p \rightarrow e + X$$

(Scaling )

( $E, q^2 = \text{exchanged}$   
momentum)

$\Leftrightarrow$  "free" quarks



$$\boxed{\alpha_3 = \kappa_3(E)}$$

high precision

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$$\bullet T(\text{fund}) = \frac{1}{2}$$

$$T(\text{Adjoint}) = ?$$

$$A = F \times \overline{F}$$

$$\Leftrightarrow F \rightarrow UF$$

$$A \rightarrow UAU^+$$

$$\underline{SU(2)} \quad F = f = D \Rightarrow T(D) = \frac{1}{2}$$

↗

$$T_a = \sigma_a / 2$$

$$\therefore T_1 (\sigma_3 / 2)^2 = 1/2$$

$A = \text{triplet}$

$$T_3 = \text{diag}(1, -1, 0)$$

$$\Rightarrow T_1 T_3^2 = 2 = T_{\text{GS}}$$

$$= T(A)^{\text{ant}}$$

$$SU(3) \quad F = 3 \Rightarrow T(F) = \frac{1}{2}$$

$$T_3 = \frac{1}{2} \begin{pmatrix} 1 & & \\ & -1 & \\ & & 0 \end{pmatrix} \Rightarrow$$

$$A = \text{octet} = 8$$

$$8 = \underbrace{3 + 2 + 2 + 1}_{\text{}} \quad$$

$$SU(3) \quad \uparrow \quad SU(2)$$

"pions" ( $\pi^+$ ,  $\pi^-$ ,  $\pi^0$ )

"kaons" ( $k^+$ ,  $k^0$ ) }  $SU(2)$   
 "anti-kaons" ( $\bar{k}^0$ ,  $k^0$ ) } doublet

$\eta'$  singlet

$$3 = \underbrace{2 + 1}_{\mathfrak{su}(3)} \quad \underbrace{\mathfrak{su}(2)}$$

$$f = 3 \times \bar{3} = (\bar{2} + 1) \times (2 + 1)$$

$$= \overbrace{3 + 2 + \bar{2} + 1}^{\text{---}} + 1$$

$$\underbrace{T(f)}_{\mathfrak{su}(3)} = T(3) + \underbrace{T(2) + T(\bar{2}) + T(1)}_{\mathfrak{su}(2)}$$

$$= 2 + \frac{1}{2} + \frac{1}{2} + 0 = 3$$

$$T(A)_{\mathfrak{su}(2)} = 2$$

$$T(A)_{\mathfrak{su}(3)} = 3$$

$$T(A)_{\text{sub}} = N$$

Proof

Try doing holiday

$$\frac{1}{\alpha(E_2)} = \frac{1}{\alpha(E_1)} + \frac{b}{2\pi} \ln \frac{E_2}{E_1}$$



particles that run:  
(p)

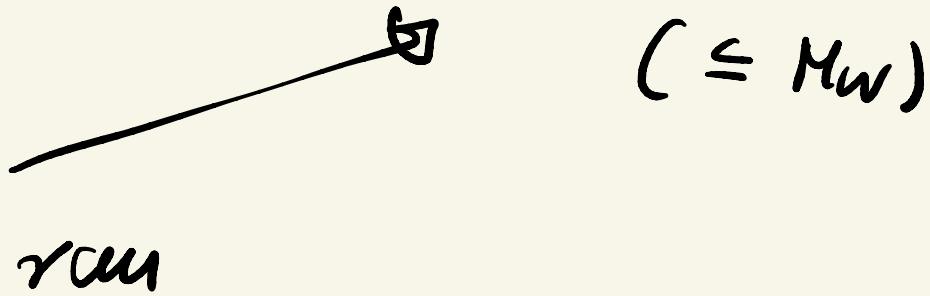
$$m_p \leq E_1 \quad (\text{center } b)$$

$$\left( \ln \frac{E_2}{E_1} \gg 1 \right)$$

$$M_x = M_y = v_{\text{GUT}}$$

S<sub>SB</sub> :

$$M_W = M_T = M_A = 0$$



$$( \leq M_W )$$

(x, y) don't "run"

b  $\not\Rightarrow$  cancellation only

from "light" particles

$$(m \leq E_1)$$

$E \gg V_{\text{out}}$

$\Rightarrow x, q$  "run" (euler b)

Kaons:  $\begin{pmatrix} k^+ \\ h^0 \end{pmatrix}, \begin{pmatrix} \bar{h}^0 \\ h^- \end{pmatrix}$

{

$$h^0 \neq \bar{h}^0$$

(CP viol.,  $\Delta m_K$ )

$$\boxed{\Delta m_K = \frac{m_{h^0} - m_{\bar{h}^0}}{m_{h^0} + m_{\bar{h}^0}} \simeq 10^{-14}}$$



Computed in SM

$\Leftrightarrow$  GIM mechanism

$$k^+ = u \bar{s} \quad \bar{h}^0 = \bar{d} s$$

$$k^0 = d \bar{s} \quad h^- = \bar{u} s$$

$$\pi^0 = (\bar{u}u - \bar{d}d) \frac{1}{\sqrt{2}}$$

$$h^0 \longleftrightarrow \bar{h}^0$$

CP

$$CP: \quad k_+ = \frac{h_0 + \bar{h}_0}{\sqrt{2}}, \quad k_- = \frac{h_0 - \bar{h}_0}{\sqrt{2}}$$

$$K_S = K_+ + \varepsilon K_- \quad (\varepsilon \approx 1/10^3)$$

$$K_L = K_- - \varepsilon K_+$$

$$\begin{aligned} K_+ &\rightarrow \pi \pi \\ K_- &\rightarrow \pi \pi \pi \end{aligned} \quad \left. \right\} CP$$

$$\boxed{\Gamma_+ \gg \Gamma_- \Leftrightarrow \tau_+ \ll \tau_-}$$