

Neutrino Mass

and

Grand Unification

Lecture XIV

3/12/2021

LMU

Fall 2021



SU(5) and Unification!

Can it be solved?

$$\frac{1}{\alpha'(E_2)} = \frac{1}{\alpha'_2(E_1)} + \frac{b}{2\pi} \ln \frac{E_2}{E_1}$$

$$b = \frac{11}{3} T_{GB} - \frac{2}{3} T_F - \underbrace{\frac{1}{3} T_S}_{\text{complex scalars}}$$

↓

$$+ \gamma_H \quad (AF)$$

1-2

$$\frac{1}{\alpha_U} = \frac{1}{\alpha'_2(M_W)} + \frac{b_2}{2\pi} \ln \frac{M_X}{M_W}$$

$$\frac{1}{\alpha_J} = \frac{1}{\alpha_1(M_W)} + \frac{b_1}{2\pi} \ln \frac{\mu_x}{M_W}$$



$$2\pi \left(\frac{1}{\alpha_1} - \frac{1}{\alpha_2} \right) | = (b_2 - b_1) \ln \frac{\mu_x}{M_W}$$

$T_{GB} = N$

for $SU(N)$, $N \geq 2$

$$T_{GB} = 0, \quad \text{for } U(1)$$

$$\Rightarrow b_2 = \frac{22}{3} - \frac{4}{3} u_F - \frac{1}{6}$$

A Higgs

$$b_1 = 0 - \frac{4}{3} u_F - \frac{1}{10} \nearrow$$

$$(b_1 = \frac{3}{5}, b' = \frac{3}{5} b_1)$$



$$M_x^{1-2} \approx 10^3 \text{ GeV}$$

$$M_x^{2-3} \approx 10^{17} \text{ GeV}$$

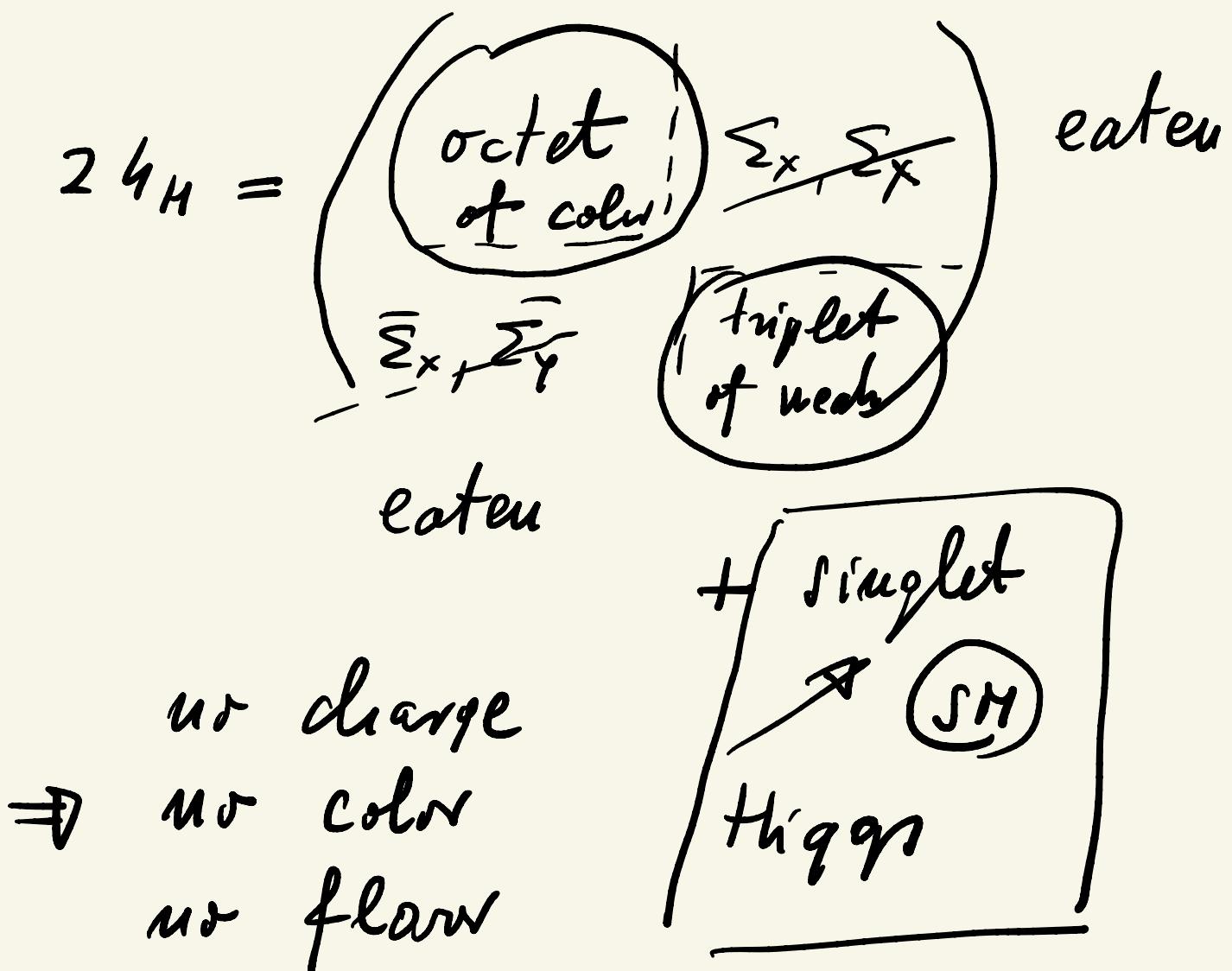
Who runs below M_x ?

- SM fermions
- W, Z, γ
- Higgs

but:

$$5_H = \left(\begin{array}{c} T \\ \dots \\ D = \Phi \end{array} \right) \rightarrow M_T > 10^2 \text{ GeV}$$

Higgs



$S_c, 3_{W^-}$ = physical

$$M_g = \frac{5}{4} v_x^2 b \quad \left\{ \begin{array}{l} M_x = \dots \vartheta_x \\ M_y = \dots \end{array} \right.$$

$$\frac{b}{2} T_r (2 \gamma_H)^4$$

$$m_3 = 5 b v_x = 4 m_f$$

$M_x = \pm g v_x$ $f \approx 0(1)$

X is heavy!

$b = ?$ what if b small?

SM : electron $w_e = y_e v_w$

$w_e \neq v_w$ \uparrow
 $\sim 100 \text{ GeV}$

$$= 10^{-5} v_w !!!$$

Q. What if $b \sim y_e \ll 1$?

- \underline{m}_3 , \underline{w}_ℓ = ?

z_w : $y = 0 \Rightarrow b_1(z_w) = 0$

$f_{b_1} \equiv b_1(z_w) = 0$

$$b_2(\beta_w) = -\frac{1}{3} \cdot \frac{1}{2} \cdot 2 = -\frac{1}{3}$$

↑
real

$b_2(\beta_w) = fb_2$

$$\frac{1}{\alpha_V} = \frac{1}{\alpha_1(M_W)} + \frac{b_1}{2\pi} \ln \frac{\mu_x}{M_W}$$

$$\frac{1}{\alpha_V} = \frac{1}{\alpha_2(M_W)} + \frac{b_2}{2\pi} \ln \frac{m_3}{M_W} +$$

$$+ \frac{(b_2 + fb_2)}{2\pi} \ln \frac{\mu_x}{m_3}$$

$$= \frac{1}{\alpha_2(M_W)} + \frac{b_2}{2\pi} \ln \frac{\mu_x}{M_W} + \frac{fb_2}{2\pi} \ln \frac{\mu_x}{m_3}$$

$$\left[\frac{1}{\alpha_1(M_w)} - \frac{1}{\alpha_2(M_w)} \right] 2\pi = (b_2 - b_1) \ln \frac{M_x}{M_w}$$

$$-\frac{1}{3} \ln \frac{M_x}{M_3}$$

negative

{

$$2\pi \left(\frac{1}{\alpha_1} - \frac{1}{\alpha_2} \right) \frac{1}{M_w} = (b_2 - b_1) \ln \frac{M_x^0}{M_w}$$

$$(M_3 = M_x)$$

$$= (b_2 - b_1) \ln \frac{M_x}{M_w} - \frac{1}{3} \ln \frac{M_x}{M_3}$$

↓

best: $M_3 \simeq M_W$



$$M_X^{(1-2)} \leq 10^{14} \text{ GeV}$$

ruled out!

Q. What is the minimal
realistic extension of minimal
 $SU(5)$?

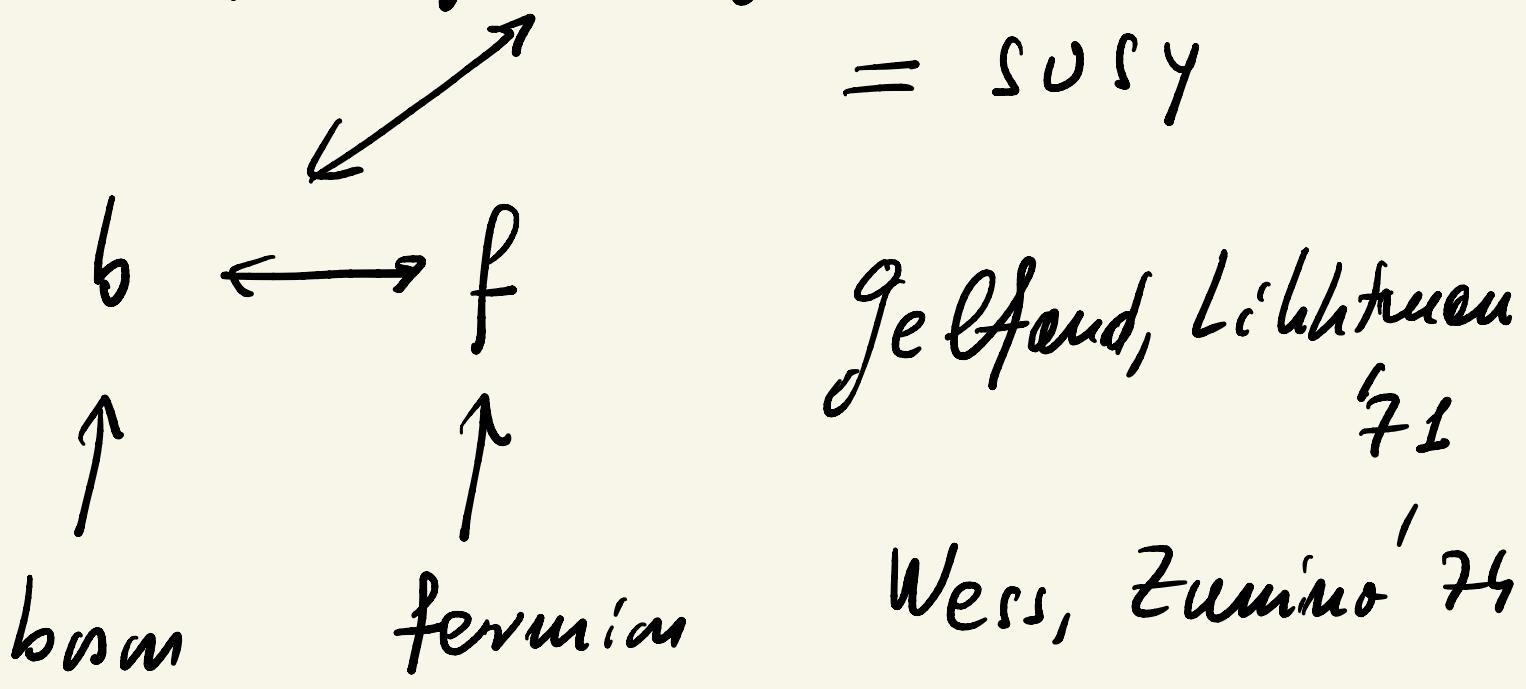
A. $24_F \rightarrow (24_V, 24_H)$

SUPER SYMMETRIC

$SU(5)$

Super Symmetry = Su Sy

= SUSY



(fl e \longleftrightarrow $\tilde{e}(b)$ (selectron))

$w \longleftrightarrow \tilde{w}$ (Wino)
(b) (f)

(b) $h \longleftrightarrow \tilde{h}$ (Higgsino) (f)

(f) $u \longleftrightarrow \tilde{u}$ (up squark) (b)

SUSY = broken

Theory (prejudice) :

$$\Lambda_{\text{SUSY}} = m_{\tilde{p}} - m_p \simeq \text{TeV}$$

$\begin{matrix} \uparrow & \uparrow \\ (\text{s-particle}) & \text{particle} \end{matrix}$



s particles run above TeV

$$b = \frac{11}{3} T_{GB} - \frac{2}{3} T_F - \frac{1}{3} T_S$$

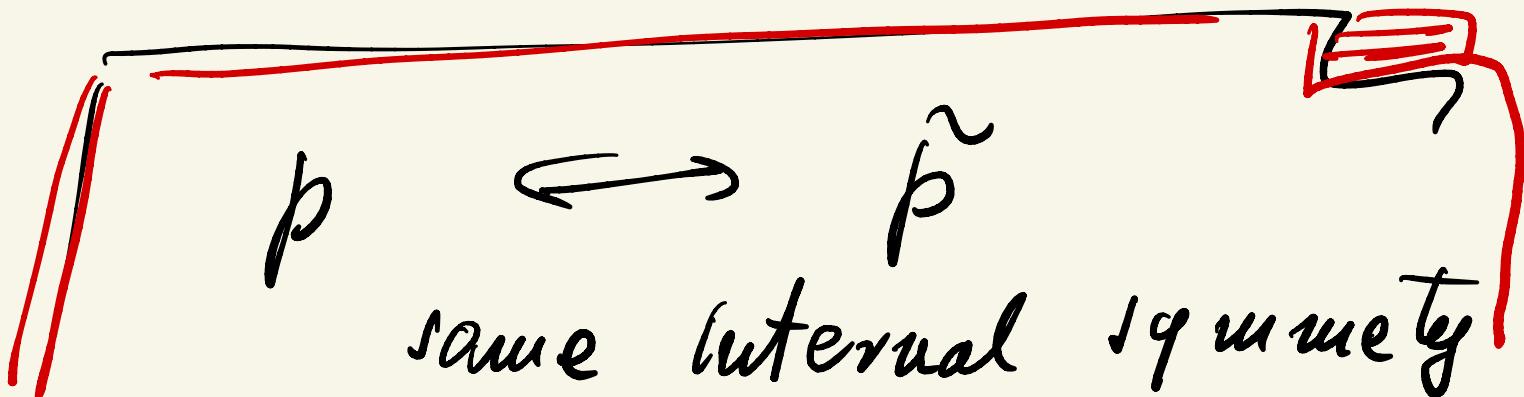


$$b^{\text{sky}} = ?$$

$$(g.\text{boson}) A \xleftarrow{} \tilde{A} \begin{matrix} (s=1/2) \\ (\text{gaugino}) \end{matrix} (f)$$

$$(\text{fermion}) f \xleftarrow{} \tilde{f} \begin{matrix} (s\text{fermion}) \\ (b) \\ \uparrow \\ (s=0) \end{matrix}$$

$$(s=0) h \xleftarrow{} \tilde{h} \begin{matrix} (s=1/2) \end{matrix}$$



LHC

$$b = \frac{11}{3} T_{GB} - \frac{2}{3} T_F - \frac{1}{3} T_S \quad (*)$$

↓

gauge boson A

gaugeino \tilde{A} → how much?

$$-\frac{2}{3} \times \overline{T}_{GB}$$

↑ fermions

sfermions

↗ gaugeino



$$b^{susy} = \left(\frac{11}{3} - \frac{2}{3} \right) \overline{T}_{GB} - \left(\frac{2}{3} + \frac{1}{3} \right) \overline{T}_F$$

$$- \left(\frac{1}{3} + \frac{2}{3} \right) T_S$$

≈ Higgsino

$$\boxed{b^{\text{susy}} = 3T_{GB} - T_F - T_S}$$

$$b_3 = 3 \cdot 3 - 2ug - 0 = 9 - 2ug$$

$$b_2 = 3 \cdot 2 - 2ug - \frac{1}{2} = 6 - 2ug - \frac{1}{2}$$

$< 0 \quad (ug=3)$

$$\boxed{\alpha'_2 \neq AF \text{ in susy}}$$

$$b_1 = -2ug - \left(\frac{3}{5}\right) \cdot \frac{1}{2}$$

normalization

$$\Rightarrow \sin^2 \theta_W = 0.23$$

(susy)

Marciano, G. S.

81

$$(\sin^2 \theta_W)^{\text{exp}} = 0.20$$

$$M_t \simeq 200 \text{ GeV}$$

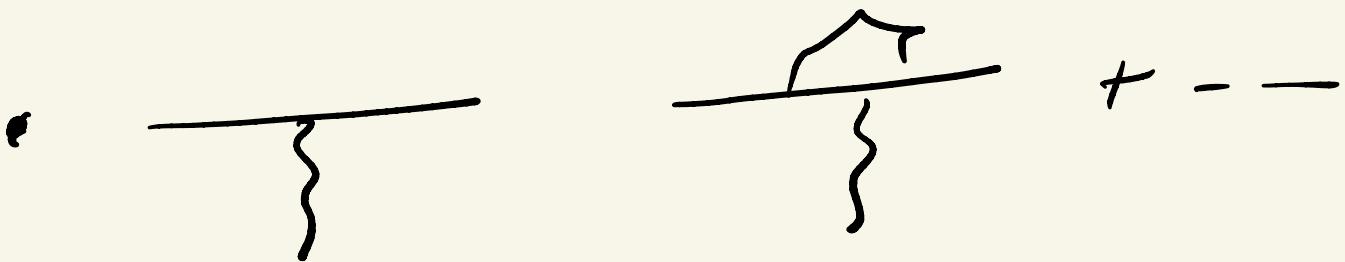
Scale of susy breaking:

$$\Lambda_{\text{susy}} = ?$$

Why a prejudice $\Lambda_{\text{susy}} \simeq \text{Tev}$?

↓

Hierarchy Problem



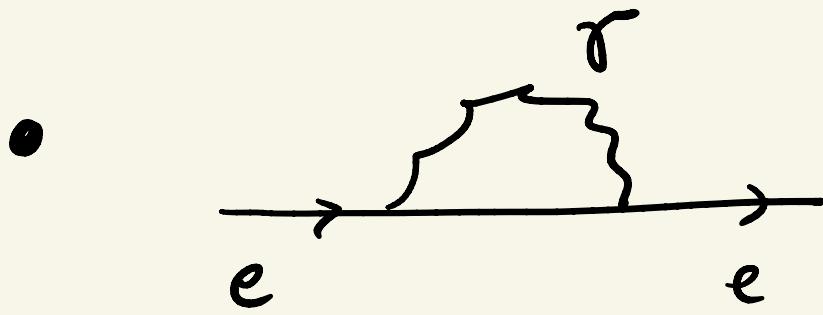
$$e^{(E)} = e_0 \left(1 + \frac{e_0^2 \ln^1 E}{4\pi} \right)$$

$\underbrace{\hspace{10em}}$

small "infinity"

$$\Lambda \lesssim M_p \simeq 10^{19} \text{ GeV}$$

$$\Rightarrow d \ln^1 E < 1$$



Verishoff 1930

$$m_e^{(E)} = m_e^0 + ad\alpha + b dm_e^0 \ln \gamma_E$$

Q1. Is $a \neq 0$ possible?

A. $a = 0$

Q2. Why $a = 0$?

$$m_e^0 \bar{e} e = m_e^0 (\bar{e}_L e_R + \bar{e}_R e_L)$$

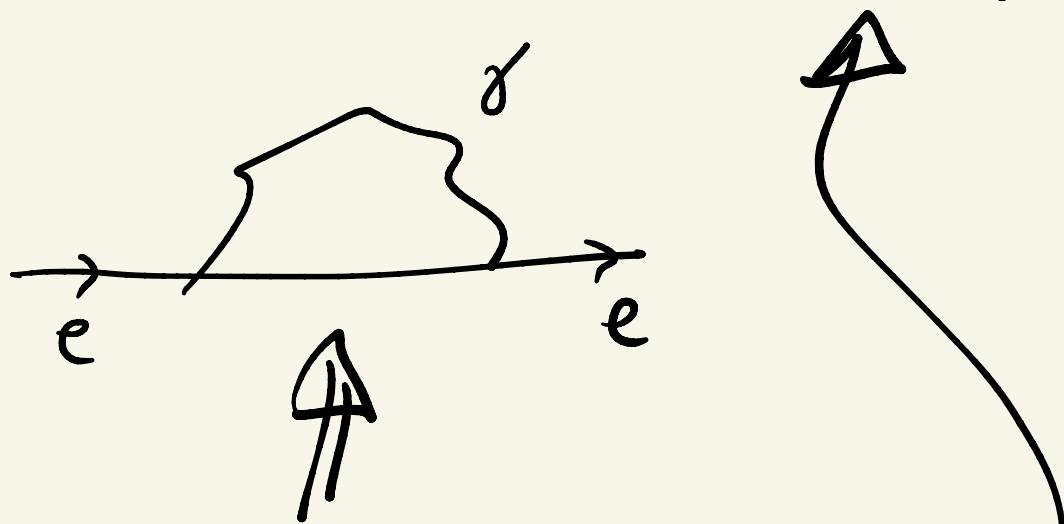
\Rightarrow if $m_e^0 = 0$

$$\Rightarrow e_L \rightarrow e^{i\alpha} e_L, \quad e_R \rightarrow e_R$$

chiral symmetry

$$e_L \rightarrow e_L, \quad e_R \rightarrow e^{i\beta} e_R$$

- $\underline{M_e^0 = 0} \Rightarrow \text{chiral symmetry}$



loops must preserve



$M_e \neq M_e^0$, so that

$$m_e^0 \rightarrow 0 \Rightarrow m_e \rightarrow 0$$

(diral symmetry)



$$m_e^{loop}(E) = m_e^0 \left(1 + \frac{\alpha}{\pi} \ln \frac{1}{E} \right)$$



to all orders!

$$m_e^{u\text{loop}}(E) = m_e^0 \left(1 + \# \left(\frac{\alpha}{\pi} \right)^u \ln \frac{1}{E} \right)$$

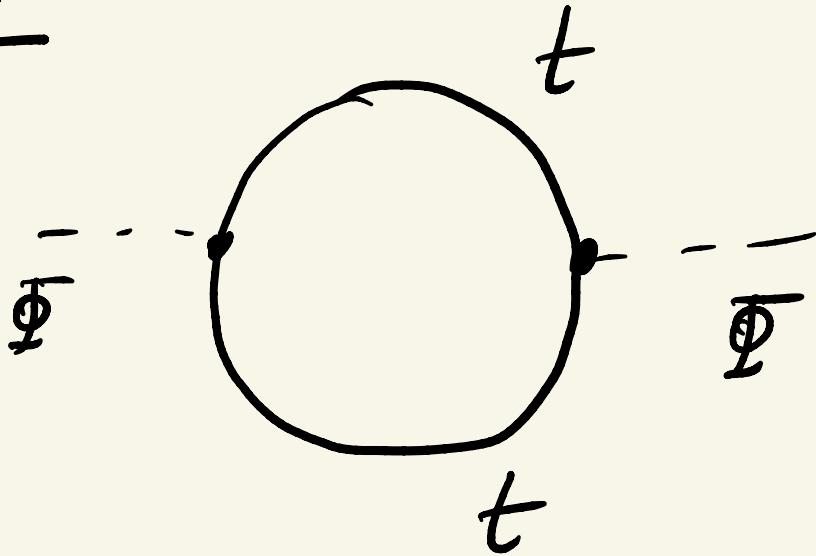
$\boxed{< 1}$

Higgs mass

$m^2 \bar{\Phi}^+ \bar{\Phi}^- \not\Leftarrow \underline{\text{no}} \text{ "chiral" sym.}$



1 loop



$$m_{\bar{\Phi}}^2 (1 \text{ loop}) \simeq \frac{4t^2}{16\pi^2} \left(\equiv \frac{\alpha'}{4\sigma} \right)$$

$$\left(1^2 + m_t^2 \ln \frac{1}{E} \right)$$

↑
"hierarchy problem"

GUT : $\lambda > M_X$!

Fine-Tuning (FT)

$$m_{\Phi}^2 = m_0^2 + \frac{d_t}{4\pi} \left(1^2 + m_t^2 \ln \frac{1}{m_0^2} \right)$$



$$m_0^2 + \frac{d_t}{4\pi} 1^2 \approx (100 \text{ GeV})^2$$

$$(10^{16} \text{ GeV})^2$$

$$(10^{16} \text{ GeV})^2$$

• chiral symmetry

$$\mathcal{L}_D = i \bar{f} \gamma^\mu \partial_\mu f - m_f \bar{f} f$$

• $m_f \neq 0 \Rightarrow \boxed{f \rightarrow e^{i\alpha} f}$

$\Leftrightarrow f_L \rightarrow e^{i\alpha} f_L, f_R \rightarrow e^{i\alpha} f_R$

Why? $\bar{f} f = \bar{f}_L f_R + \bar{f}_R f_L$

• $m_f = 0 \Rightarrow f_L \rightarrow e^{i\beta} f_L, f_R \rightarrow f_R$

$(f_L \rightarrow f_L, f_R \rightarrow e^{i\delta} f_R)$

$\Leftrightarrow f_L \rightarrow e^{i\epsilon} f_L, f_L \neq f_R + \dots$



- $f \rightarrow e^{i\alpha} f$ (vector symmetry)

always

- $m_f = 0 \Rightarrow f \rightarrow e^{i\beta \gamma_5} f$ (chiral)



$$f_L \rightarrow e^{i\beta} f_L, f_R \rightarrow e^{-i\beta} f_R$$

$L \neq R$