

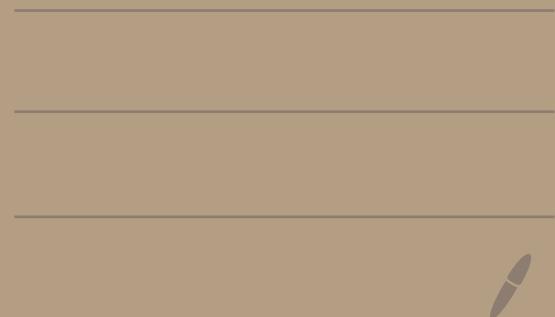
LMU GUT Course

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Lecture XV

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22/12/2020



Symmetry breaking in  $SU(5)$ :

Higgs and Yukawa sectors

$$SM = SU(3) \times \underbrace{SU(2) \times U(1)}_{V(1)}$$

$$M_W \propto \langle H \rangle \quad \text{Higgs } H = \begin{matrix} \text{doublet} \\ V(1) \end{matrix}$$

•  $SU(5) \Rightarrow (x, y)$  gauge bosons

$$m_x = m_y > 10^{15} \text{ GeV}$$

$$SU(5) \xrightarrow{\Sigma} SU(3) \times SU(2) \times U(1)$$

$$M_{GUT} = M_x$$

Who should  $\Sigma$  be?

example  $SU = SU(2) \times U(1)$

$$l_2^+ = (\bar{e})_L, e_R \quad \downarrow \\ U(1)_{\text{em}}$$

$$\mathcal{L}_Y = \bar{l}_L H e_R \gamma_e$$

↑      ↑      ↑  
doublet      doublet      singlet

•  $\Sigma$  in  $SU(5)$

$\Sigma$  should not couple  
to fermions

$f = \bar{5}_L, 10_L$  (anti sym)

$(5_R = \begin{pmatrix} d^{\alpha} \\ \bar{e}^c \\ \bar{u}^c \end{pmatrix})$  color  
weak doublet

$\bullet 5_H = \begin{pmatrix} T^{\alpha} \\ - \\ H \end{pmatrix}$  weak doublet

$\langle 5_H \rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ - \\ 0 \end{pmatrix} \quad \overbrace{\quad}^3 \text{ weak}$

$SU(5) \rightarrow \cancel{SU(3) \times SU(2) \times U(1)}$

$\langle 5_H \rangle$

NO

$SU(5) \xrightarrow{\langle 5_H \rangle} SU(4)$

~~$10_A$~~

$$(10_F = \begin{pmatrix} u^c & & & \\ & d & & \\ & & e & c \end{pmatrix})$$

$\langle 10_A \rangle$   
breaks  $Q_{em}$ !

→  
no neutral field

•  $15_H \leftarrow$  does not work

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$$SU(5) \xrightarrow{\langle \Sigma \rangle} SU(3) \times SU(2) \times U(1)$$

$\gamma = 4 \quad 2 \quad 1 \quad 1$

$\boxed{\langle \Sigma \rangle \text{ does not break}} \\ \text{reality}$

Adjoint:  $\Sigma \rightarrow U \Sigma U^+$

$$\boxed{T_a \Sigma = 0}$$

$$\Sigma^+ = \Sigma$$

$$U U^+ = 1$$

$$\Rightarrow \det U = 1$$

defining  $5 \times 5$  unitary

$$\langle \Sigma \rangle \rightarrow \boxed{U \langle \Sigma \rangle U^+ = \text{diagonal Hermitian}}$$

$$\langle \Sigma \rangle = \text{diagonal}$$

$$\Rightarrow \left[ \Sigma, T_a \in C \right] = 0$$

↑  
Cartan

$$\Sigma \rightarrow U\Sigma U^+$$

$$= \Sigma + i \sum_a [T_a, \Sigma] + \dots$$

$\langle \Sigma \rangle$  preserves rank

$$\Sigma = T_a \phi_a \quad a = 1, \dots, 24$$

$T_V \Sigma^2,$   
 $\dots T_X \Sigma^3,$   
 $T_Y \Sigma^4,$

$$\Sigma^2 \rightarrow U\Sigma^2 U^+$$

$$\Sigma^3 \rightarrow U\Sigma^3 U^+$$

$$\Sigma^4 \rightarrow \dots$$

$$m \Sigma^3 \text{ i.u } V \\ (m = 0)$$

$$V = -\frac{\mu^2}{2} \text{Tr} \Sigma^2 \leftarrow \text{mass}$$

$$+ \frac{a}{4} (\text{Tr} \Sigma^2)^2 + \frac{b}{2} \text{Tr} \Sigma^4$$

$\langle \Sigma \rangle = \text{diagonal}$

$$\langle \Sigma \rangle = v_x \text{ diag } (1, 1, 1, -\frac{3}{2}, -\frac{3}{2})$$

↑

preserve  $SU(3)_C, SU(2)_L, U_Y^{(1)}$

let us prove that this is  
a minimum

Li '1974

Posted

↑  
global minimum

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SH       $H = \text{doublet}$

$$V = -\frac{\mu_H^2}{2} H^+ H + \frac{\lambda}{4} (H^+ H)^2$$

- $\lambda > 0 \leftarrow \text{boundedness}$
- $\mu_H^2 > 0 \leftarrow \text{sgn. breaking}$

range of parameters that  
ensures (proper) sgn. breaking



$$\langle \Sigma^2 \rangle = v_x^2 \operatorname{diag} (1, 1, 1, \frac{9}{4}, \frac{9}{4})$$

$$\langle \Sigma^4 \rangle = v_x^4 \operatorname{diag} (1, 1, 1, \frac{81}{16}, \frac{81}{16})$$

$$T_v \Sigma^2 = v_x^2 \frac{15}{2}$$

$$T_v \Sigma^4 = v_x^4 \frac{105}{8}$$



$$V = -\frac{\mu}{2} \frac{15}{2} v_x^2 + \frac{a}{4} v_x^4 \left(\frac{15}{2}\right)^2$$

$$+ b/2 \frac{105}{8} v_x^4$$

$$\frac{\partial V}{\partial v_x} = 0$$



$v_x \neq 0$  not physical

$$\frac{\partial^2 V}{\partial \vartheta_k^2} \Big|_0 = -\mu < 0 \quad \text{not a}$$

extremum

minimum

$$\mu^2 = \frac{15a + 7b}{2} \vartheta_{x^2}$$

step 1

$$(w \ll \mu; \quad w T_V \Sigma^3)$$

~~$T_V \Sigma, T_V \Sigma^2, T_V \Sigma^3, T_V \Sigma^4$~~

$$T_V \Sigma^5 \leftarrow T_V \Sigma^2 T_V \Sigma^3$$

↑ not new

Guth, Welzberg  
'82

deal with

$$w \neq 0$$

$\Downarrow$  show  $\langle \Sigma \rangle = \text{local}$   
minimum

Study 2<sup>nd</sup> derivatives matrix

$\Downarrow$   
show that the eigenvalues  
are positive

$\frac{\partial^2 V}{\partial \Sigma_i \partial \Sigma_j} \Big|_{\langle \Sigma \rangle}$  = matrix at  
particle  
masses

(scalar:  $\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{m^2}{2} \phi^2$ )

$$m^2 = \frac{\partial^2 V}{\partial \phi^2} \Big|_{\min}$$

I need to compute the  
masses in  $\Sigma$  ( $\phi_i \in \Sigma_i$ )

$24 \times 24$  matrix

$G \rightarrow H$   
 $SU(5) \rightarrow SU(3) \times SU(2) \times U(1)$

use the symmetries of

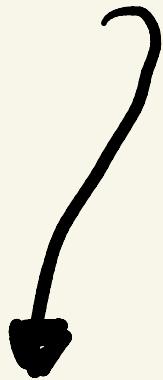
$H$

particles connected by  
symmetry  $\rightarrow$  the same mass

$$\bar{\Sigma} \rightarrow V \bar{\Sigma} V^+ \sim 5 \times \overline{5}$$

$$5 - V5$$

$$\overline{5} - \overline{5} V^+$$



$$Q(z_{ij}) = (Q_i - Q_j)$$

$$Q(10_{ij}) = (Q_i + Q_j)$$

$$\Sigma = \begin{pmatrix} \Sigma_8 & \begin{matrix} \downarrow & \downarrow \\ \bar{\Sigma}_x & \bar{\Sigma}_y \end{matrix} \\ (\sim \text{gluons}) & | \\ \hline \bar{\Sigma}_x & \left\{ \begin{matrix} \Sigma_3 \\ (\sim \vec{w}) \end{matrix} \right. \end{pmatrix} \left. \begin{array}{l} \{ \text{SU}(3) \\ \{ \text{SU}(2) \end{array} \right.$$

$$+ \bar{\Sigma}_0 \text{ (singlet)} \\ (\sim B \text{ of } U(1)_Y)$$

$$Q(\Sigma_{14}) = Q_1 - Q_4 = -\frac{1}{3} - 1 \\ = -\frac{4}{3}$$

$$\Sigma_F = \begin{pmatrix} d \\ -e \\ ec \\ rc \end{pmatrix} / R$$


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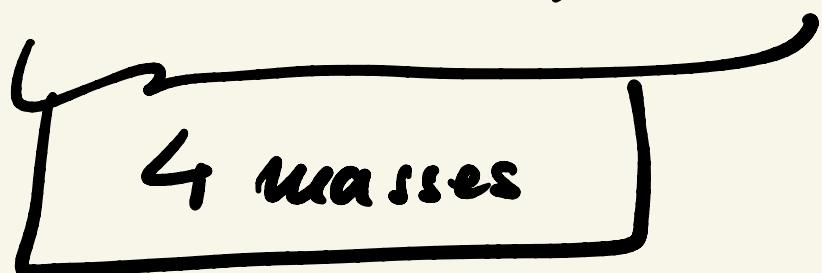
Masses :  $\Sigma_8$  — same mass (color)

$\Sigma_3$  — same mass (weak)

$\bar{\Sigma}_0$  — mass

$\underbrace{\Sigma_x^{\alpha}, \Sigma_y^{\alpha}}$  — same mass  
weak

masses:  $\Sigma_8, \Sigma_3, \Sigma_0, \Sigma_X$



$$SU(5) \rightarrow SU(3) \times SU(2) \times U(1)$$

$$24 \quad \underbrace{8 + 3 + 1}_{12}$$

$$24 - 12 = 12 \text{ broken}\\ \text{generations}$$

Higgs  $\Rightarrow$  12 massive gauge  
 $(x^\alpha, y^\alpha)$

$\Rightarrow$  12 scalars get "eaten"

potential  $\Rightarrow$  global  
symmetry

- $\mathcal{V}(\Sigma) \leftarrow$  masses

12 broken glu.  $\Rightarrow$  12 NG bosons

$NG \equiv$  Nbroken  
- Goldstone

$$m = 0$$

Guess which fields become

NG bosons



$$\sum_{x_i} \sum_y = NG$$

\* check

$m=0$  ( $\text{no } \Sigma^3$ )  $\Rightarrow D: \Sigma \rightarrow -\Sigma$

$\langle \Sigma \rangle \neq 0 \Rightarrow \cancel{\text{spont.}}$

$\Rightarrow$  domain walls

$\Sigma_8, \Sigma_3, \Sigma_0$

Notes:  
Beyond

$$m_{\Sigma_0}^2 = \frac{15a + 7b}{2} v_x^2$$

$$\mu = \frac{15a + 7b}{2} v_x^2$$

$$m_{\Sigma_0} = \mu^2 \text{ (Higgs mechanism)}$$

$\Sigma_g \rightarrow$  one states (G)

$\Sigma_g \rightarrow -11-$  (A)

$$\Sigma = \langle \Sigma \rangle + \begin{pmatrix} G/2 & & & \\ & -G/2 & & \\ & & 0 & \\ & & & 0 \\ & & & 0 \end{pmatrix}$$

$$+ \begin{pmatrix} 0 & 0 & 0 & +A/2 \\ & 0 & 0 & -A/2 \\ & & 0 & \\ & & & 0 \end{pmatrix}$$



$$\Sigma = \text{diag} \left( \frac{G}{2} + \vartheta_x, -\frac{G}{2} + \vartheta_x, \vartheta_x, A/2 - \frac{3}{2}\vartheta_x, -A/2 - \frac{3}{2}\vartheta_x \right)$$

$$\Sigma^2 = \text{diag} \left( 6\frac{\partial}{\partial t} + \partial_x^4 + 6\partial_x^2, 6\frac{\partial}{\partial t} + \partial_x^4 - 6\partial_x^2, \right. \\ \left. \partial_x^2, A^2\frac{\partial}{\partial t} + 9\frac{1}{4}\partial_x^2 - 3\frac{1}{2}\partial_x^4, \right. \\ \left. A^2\frac{\partial}{\partial t} + 9\frac{1}{4}\partial_x^4 + 11 \right)$$

$\boxed{\Sigma^4 = \text{keep } 6^2, A^2 \text{ terms}}$



$$-\frac{\mu^2}{2} T_V \Sigma^2 \rightarrow -\frac{\mu^2}{2} \left( \frac{G^2}{2} + \frac{A^2}{2} \right)$$

$$\frac{a}{q} (T_V \Sigma^2)^2 \rightarrow \frac{15}{2} a \frac{1}{4} \left( G^2 + A^2 \right) \partial_x^2$$

$$\frac{b}{2} T_V \Sigma^4 \rightarrow \frac{b}{4} \left[ G^2 6 \partial_x^2 (?) \right. \\ \left. + A^2 \frac{27}{2} \partial_x^2 (?) \right]$$

$$u = \frac{15a + 7b}{2} \vartheta x^2 \quad (1)$$



$$15a + 7b > 0$$

$$m_B^2 = \frac{5}{4} b \vartheta x^2$$

$$m_A^2 = 5b \vartheta x^2$$

$$b > 0$$

↑

(2)

ensure that

$\langle \Sigma \rangle = \text{local}$   
 $\text{minimum}$

$$\cdot b = 0 \Rightarrow m_B = m_A = 0$$

$$\Rightarrow m_{\Sigma_8} = m_{\Sigma_2} = 0$$

why massless in  $L=0$  limit?

Hint: NG mechanism

$$m_A = 2 m_\phi$$

$$(m_3 \equiv m_{\Sigma_3}, m_8 \equiv m_{\Sigma_8})$$

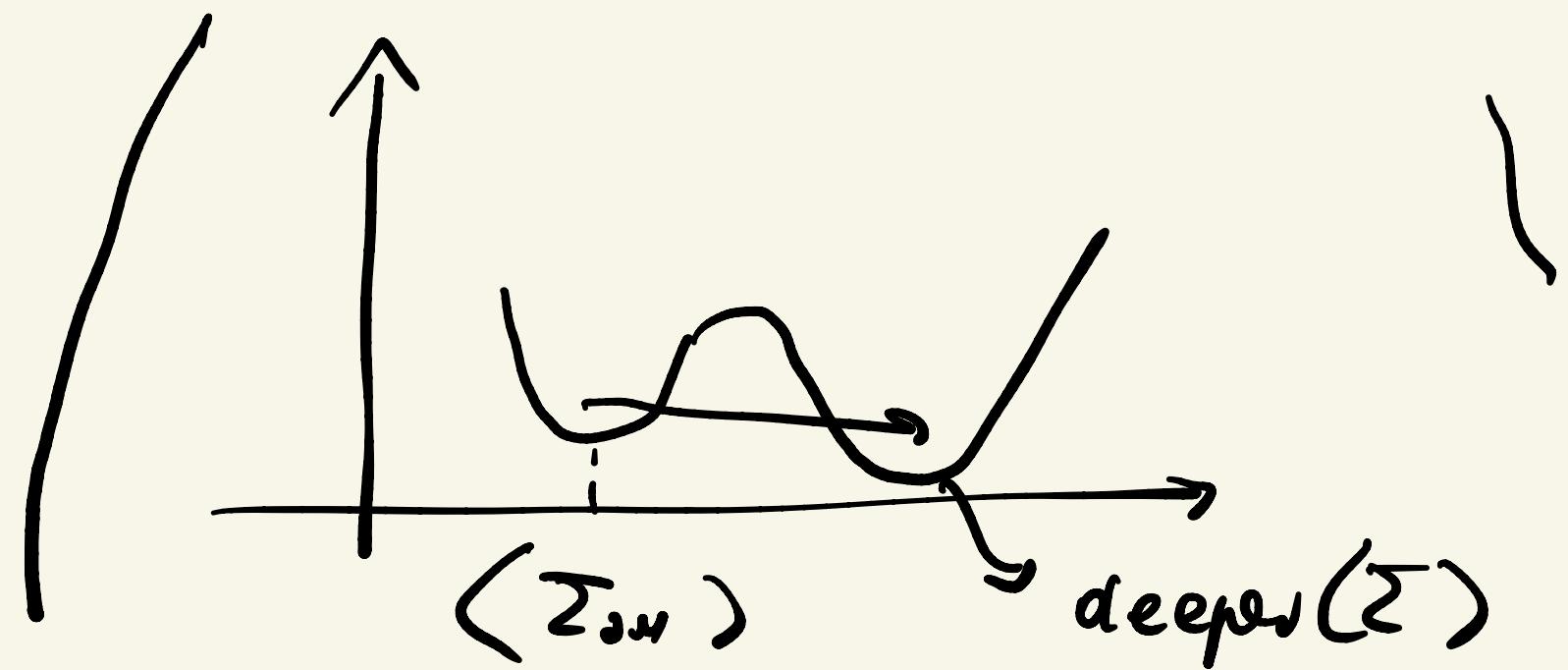
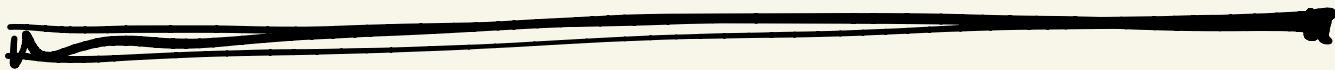
$$m_3 = 2 m_8$$

$$L_i : \langle \Sigma \rangle_{SM} = \partial_x (1, 1, 1, -3/2, -3/2)$$

= local minimum  $\Rightarrow$   
global minimum

$$\cdot \langle \Sigma' \rangle = \underbrace{\partial_{x'}(1, 1, 1, 1, -4)}$$

$SU(5)$  maximum  
 $\rightarrow SU(3) \times U(1)$



as long as  $T_{\text{tun}} > T_J$

side  
comment

age of  
universe

$$\bullet \text{SU}(I) \longrightarrow \text{SU}(3) \times \text{SU}(2) \times \text{U}(1)$$

$$\langle \Sigma \rangle = v_x \text{diag} (1, 1, 1, -\frac{3}{2}, -\frac{3}{2})$$

$$b > 0, \quad 15a + 7b > 0$$

$$\mu^2 > 0$$

$$\bullet \text{SU}(2) \times \text{U}(1) \quad \frac{1}{\langle \bar{\psi} \psi \rangle} \quad V(1)$$

$$\overline{\rho} = ?$$

$\Phi \supseteq H$  (must!)

↓  $\alpha$  SM doublet

$$\boxed{\Phi = S_H}$$

$$\Phi = \begin{pmatrix} T^1 \\ T^2 \\ T^3 \\ \phi^+ \\ \phi^0 \end{pmatrix} \begin{array}{l} \text{col} \\ \text{weak} \end{array}$$

↓

$$\langle \bar{\Phi} \rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} \boxed{\text{must}}$$

How to achieve this?

$$\Sigma = \Sigma_8 + \Sigma_3 + (\underbrace{\Sigma_x, \Sigma_y}_{\text{eaten}})$$

↗

$\Sigma_0$

↑      ↓ singlet

Color octet

**NO DOUBLET**

$$V = V_\Sigma + V_{\bar{\Phi}} + V_{\Sigma\bar{\Phi}}$$

(above)

$$V_{\bar{\Phi}} = -\mu_{\bar{\Phi}}^2 \bar{\Phi}^\dagger \bar{\Phi} + \frac{\lambda}{4} (\bar{\Phi}^\dagger \bar{\Phi})^2$$

$$\begin{aligned} V_{\Sigma\bar{\Phi}} &= \alpha \bar{\Phi}^\dagger \bar{\Phi} V \Sigma^2 + \\ &+ \beta \bar{\Phi}^\dagger \Sigma^2 \bar{\Phi} \end{aligned}$$

$$\bar{\Phi} \rightarrow U \bar{\Phi}$$

]

$$\Sigma \rightarrow U \bar{\Sigma} U^+$$

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$$\langle \bar{\Phi} \rangle = ?$$

$\Sigma$  fields - heavy

decouple

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$$\Rightarrow \Sigma \rightarrow \langle \Sigma \rangle$$

↓

$$V_{\bar{\Phi}}(\text{physical}) = V_{\bar{\Phi}} +$$

$$\alpha \bar{\Phi}^+ \bar{\Phi} \langle \nabla \Sigma^2 \rangle +$$

$$+ \beta \bar{\Phi}^+ \langle \Sigma^2 \rangle \bar{\Phi}$$

↓

$$\alpha \text{ term: } -\mu_{\Phi}^2 + \alpha \text{ Tr} \langle \Sigma^2 \rangle \\ = -\bar{\mu}_{\Phi}^2$$

only changes the mass term

only changes the mass term

$$\beta \left[ T^+ T^- \partial_x \gamma + \frac{g}{4} H^+ H^- \partial_x^2 \right]$$

$$\partial_x = \text{diag} (1, 1, 1, -3/2, -3/2)$$

$$T^+ T^- = T_\alpha^* T_\alpha \quad \alpha = 1, 2, 3$$

$H^+ H^-$  = usual 5D invariant

\*  $T^\alpha \longleftrightarrow d$

( $-1/3$  charge)

$\begin{pmatrix} x \leftrightarrow \\ y \end{pmatrix} \overset{\alpha}{\longleftrightarrow} 4/3$   
 $1/3 *$

= T mediates proton decay

Yukawa sector

$\bar{s}_F^{(L)}(s_F(r)), D_F(L)$

$s_u = \bar{\Phi}$



$$\mathcal{L}_y = \bar{S}_F^* \bar{\Omega}_F \bar{S}_H^* \gamma_1$$

$$\bar{\Omega}_F = U \Omega_F U^T$$

$$\bar{S}_F \rightarrow U S_F$$

$$\bar{S}_H \rightarrow U S_H$$

$$\bar{S}_F \underbrace{U^+ U}_{I} \bar{\Omega}_F \underbrace{U^T U^*}_{I} \bar{S}_H^*$$

$= \text{invariant}$

$$T_1 \sim S_H$$

$$T_1^* \bar{S}_F^2 \bar{\Omega}_F^{21}$$

$$(= T_1^* \bar{d}_{R2} u_2^{e3})$$

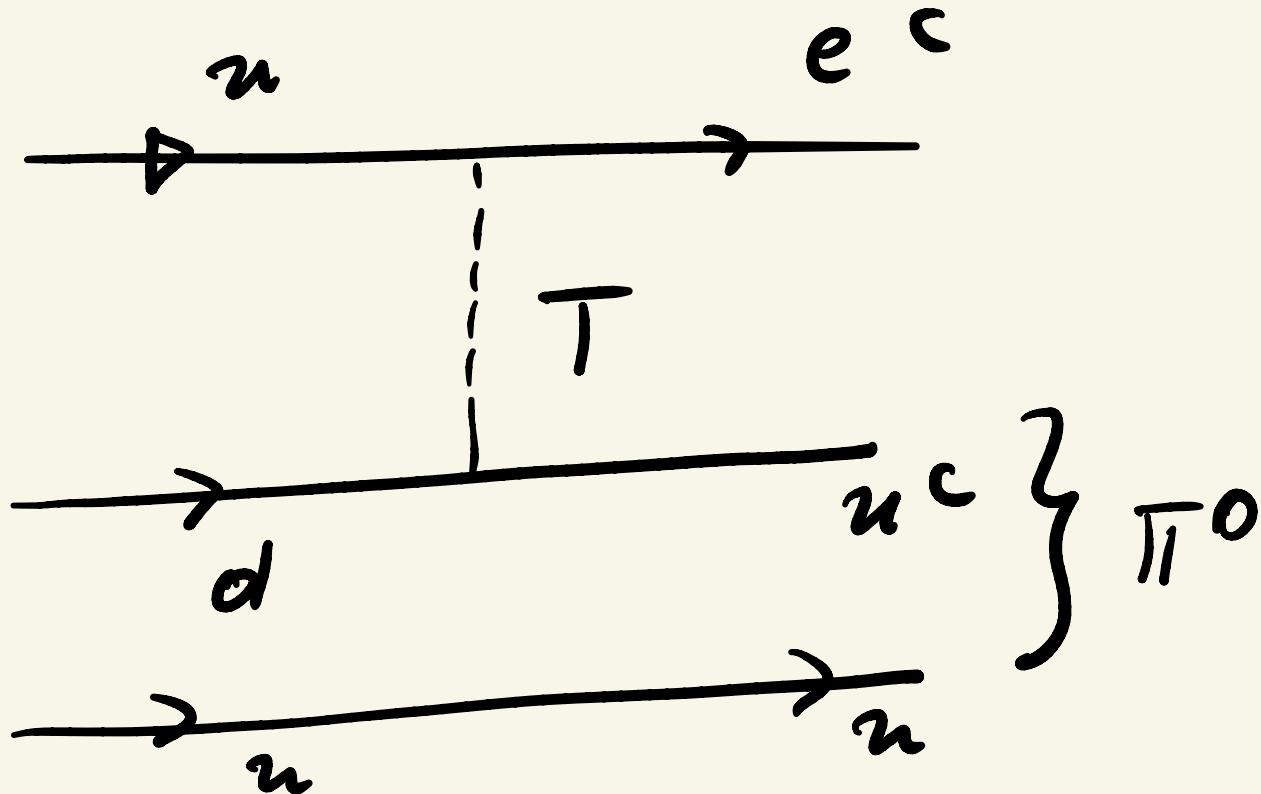
$$+ T_1^* \bar{5}_F^4 10_{41} =$$

$$( = T_1^* \bar{e}_R^c u_{L1} )$$

↓

$$\Rightarrow T \left( \bar{u}_L^c d_R + \bar{u}_L e_R^c \right) y_1$$

Q:  $-\frac{1}{3}$      $\frac{2}{3} - \frac{1}{3}$      $-\frac{2}{3} + 1$



$T \rightarrow \rho$  decay

$m_T > 10^{10} \text{ GeV}$

$$\beta [ T^+ T + \frac{g}{4} H^+ H ] v_x^2$$

$$-\mu^2 (H^+ H + T^+ T) +$$

$$= (-\mu^2 + \beta \frac{g}{4} v_x^2) (H^+ H) +$$

$$(-\mu^2 + v_x^2 \beta) T^+ T$$

positive and large

$$-\mu' + \frac{g}{4}\rho \rightarrow v_x^2 = M_w^2$$

$\approx 0$

$(M_w \ll \mu_x \sim v_x)$

Fine Tuning

$$(-\mu + \frac{g}{4}\rho v_x^2 - \frac{5}{4}\rho v_x^2) T T$$

0                          ↓

$$\boxed{\zeta < 0}$$

T gets a large ( $\zeta$ ) mass term

$$\overbrace{\langle T \rangle = 0}^{\text{H}}$$

$$\Rightarrow \langle h \rangle \neq 0$$

$$-\mu + \frac{g}{4} \rho v_x^2 < 0$$

complete the symmetry  
breaking

$$M_T = M_{\text{cut}}$$

$$\Rightarrow \rho v_x \sim M_{\text{cut}}$$

$$(-\mu^2 + \nu M_{\text{cut}}^2 = 0) = -M_w^2$$

$10^{30} \text{ GeV}^2 \quad 10^{30} \text{ GeV}^2 \quad (100 \text{ GeV})^2$

FT

Summary

$\Sigma \Rightarrow \Sigma_8, \Sigma_3, \Sigma_0 (\mu_0)$

$m_3 = 2 m_8 = ?$

$\mu_0 = ?$

$\overline{\Psi}$  :  $T$  :  $m_T \approx M_{\text{soft}}$

$H$  : light



complete the Yukawa  
story

and study fermion  
masses

$$\overline{s}_F^i 10_F^{i'} 5_H^{*\dagger} = \text{inv.}$$

$$G_{ij\mu\nu} 10_F^{i'} 10_F^{\nu*} \overline{s}_H^\mu$$

# Proton decay

$(x, y)$   $\alpha_{1/3}$

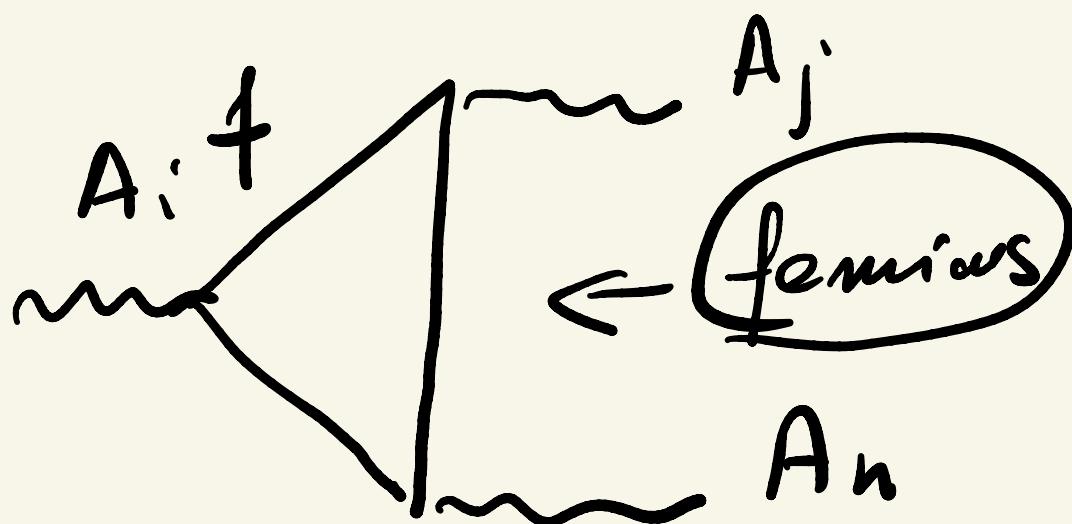
color

$T^\alpha (1/3)$

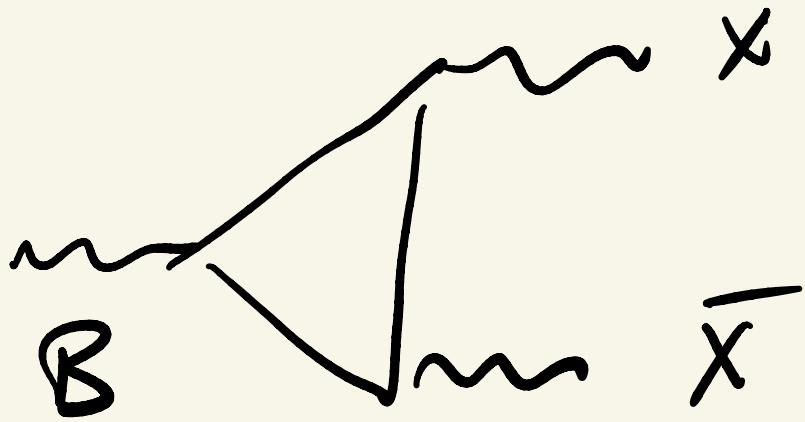
color

fractional charge

$$T_a \bar{u}^a e^c + T_a^* \bar{d}_s u_s^c \text{ color}$$



scalars = irrelevant



anomaly =  $\Theta(?)$

- $m_h = \sqrt{\lambda} / g M_W$

$$\lambda \sim g \Rightarrow m_h \sim M_W$$

- $M_T^2 = \beta v_X^2 \Rightarrow \beta \rightarrow 1$   
(p decay)

$$M_2 = 2w_8 \propto \sqrt{6} \propto$$

$r_0$  arbitrary

a little "rasis"

octet at center; triplet at ends