

# Exercises for Conformal Field Theory (MD4)

## Problem set 4, due November 20, 2019

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### 1 Vertex operators

Let us consider the free boson CFT from the lecture. Next to the current  $j(z) = i\partial X(z, \bar{z})$  there is another conformal primary

$$V_\alpha(z, \bar{z}) = : e^{i\alpha X(z, \bar{z})} : \quad (1)$$

called vertex. In the following we will verify that  $V(z, \bar{z})$  is indeed a conformal primary and compute its eigenvalues under the current  $j(z)$ . This is most easily done using Wick's theorem. In case you are unfamiliar with that we collected all necessary information at the end of the sheet.

- A) Let us compute the momentum of the vertex operator. As we have a free theory  $p = \pi$  and due to  $\pi = j_0$  this amounts to computing the commutator  $[j_0, V]$ . To do so recall the integration of the mode expansion of the current

$$X(z, \bar{z}) = x_0 - i j_0 \log z + i \sum_{n \neq 0} \frac{j_n}{n} z^{-n} + \text{antiholomorphic} \quad (2)$$

from the lecture and use  $[j_m, j_n] = m\delta_{m, -n}$  as well as the Heisenberg commutation relation.

- B) Given a conformal primary  $\phi(z, \bar{z})$  with momentum  $p_\phi$

$$[j_m, \phi_n] = p_\phi \phi_{m+n} \quad (3)$$

derive the singular part of the OPE  $j(z)\phi(w, \bar{w})$ . State the result you expect for the OPE  $j(z)V_\alpha(w, \bar{w})$ .

- C) Verify the guess for  $j(z)V_\alpha(w, \bar{w})$  from the last exercise using Wick's theorem. As you can see, the momentum can be read off from the OPE.
- D) Show that  $V_\alpha(w, \bar{w})$  is a conformal primary by computing the OPE  $T(z)V_\alpha(w, \bar{w})$  using Wick's theorem. What is the conformal dimension?
- E) The Lagrangian of the free boson  $\mathcal{L} = \frac{1}{2}\partial X \bar{\partial} X$  is invariant under shifts  $X(z, \bar{z}) \rightarrow X(z, \bar{z}) + a$  for constant  $a$ . Use this symmetry to deduce for which values of  $\alpha$  and  $\beta$  the correlator of two vertex operators

$$\langle V_\alpha(z, \bar{z})V_\beta(w, \bar{w}) \rangle \quad (4)$$

is non-zero. Interpret this result in terms of the symmetry and state the result of the above correlator.

- F) By the operator-state correspondence there is a state

$$|\alpha\rangle := \lim_{|z| \rightarrow 0} V_\alpha(z, \bar{z})|0\rangle \quad (5)$$

Furthermore the primary  $j(z) = i\partial X(z, \bar{z})$  gives rise to

$$\lim_{|z| \rightarrow 0} i\partial X(z, \bar{z})|0\rangle = j_{-1}|0\rangle \quad (6)$$

As such a general state in the free boson CFT is the combination of both

$$\lim_{|z| \rightarrow 0} : (i\partial X(z, \bar{z}))^{n_1} (i\partial^2 X(z, \bar{z}))^{n_2} \dots V_\alpha(z, \bar{z}) : |0\rangle = j_{-1}^{n_1} j_{-2}^{n_2} \dots |\alpha\rangle, \quad (7)$$

where we omitted the antiholomorphic piece. State its conformal weight and its momentum.

G) We specialize to momenta  $\alpha = \sqrt{2}m$  with  $m \in \mathbb{Z}$ . Denote the Hilbert space spanned by the above vectors with  $V_{\mathbb{Z}}$ . Compute the character

$$\chi(\tau, z) = \text{Tr}_{V_{\mathbb{Z}}} (q^{L_0 - \frac{c}{24}} x^{j_0}), \quad x = e^{-2\pi iz} \quad . \quad (8)$$

(Note: This is also the character for the  $\hat{\mathfrak{su}}(2)_1$  vacuum representation.)

H) (optional) In string theory the free boson takes the role of the coordinate of the target space. As such we need  $D$  free bosons  $X^\mu$  with  $\mu = 1, \dots, D$  instead of only one  $X$ . As you maybe heard, string theory naturally predicts gravity. Concretely the graviton appears in the state<sup>1</sup>

$$\lim_{|z| \rightarrow 0} \epsilon_{\mu\nu} : i\partial X^\mu(z, \bar{z}) i\bar{\partial} X^\nu(z, \bar{z}) V_k(z, \bar{z}) : |0\rangle = \epsilon_{\mu\nu} j_{-1}^\mu \bar{j}_{-1}^\nu |k\rangle \quad . \quad (9)$$

Analogously to (1) compute the OPE

$$T(z) \epsilon_{\mu\nu} : i\partial X^\mu(w, \bar{w}) i\bar{\partial} X^\nu(w, \bar{w}) V_k(w, \bar{w}) : \quad . \quad (10)$$

Next to the expected terms you will find a cubic term such that the above state is not a conformal primary. What needs to be imposed to get rid of this anomalous term? Interpret this relation!

Of course, the graviton should be massless. What condition onto the momentum  $k$  ensures the masslessness? What is the conformal dimension of a massless graviton?

## Wick's theorem for free fields

The contraction of two operators  $A$  and  $B$  is defined to be

$$A^\bullet B^\bullet \equiv AB - :AB: \quad . \quad (11)$$

Recalling the full OPE for a scalar  $X$

$$X(z, \bar{z})X(w, \bar{w}) = -\log|z-w|^2 + :X(z, \bar{z})X(w, \bar{w}): \quad (12)$$

one sees that the contraction picks out the singular part of the OPE. For free theories Wick's theorem reads

$$ABCDEF \dots = :ABCDEF \dots: + \sum_{\text{singles}} :A^\bullet B^\bullet CDEF \dots: + \sum_{\text{doubles}} :A^\bullet B^\bullet C^\bullet D^\bullet EF \dots: + \dots \quad (13)$$

The first sum is over all possibilities to have a single contraction while the second sum is over all possibilities to have two contractions and so forth. Let us give examples to clarify how to perform the first sum

$$\begin{aligned} X(z_1, \bar{z}_1)X(z_2, \bar{z}_2)X(z_3, \bar{z}_3) &= :X(z_1, \bar{z}_1)X(z_2, \bar{z}_2)X(z_3, \bar{z}_3): + \\ &+ :X^\bullet(z_1, \bar{z}_1)X^\bullet(z_2, \bar{z}_2)X(z_3, \bar{z}_3): \\ &+ :X^\bullet(z_1, \bar{z}_1)X(z_2, \bar{z}_2)X^\bullet(z_3, \bar{z}_3): \\ &+ :X(z_1, \bar{z}_1)X^\bullet(z_2, \bar{z}_2)X^\bullet(z_3, \bar{z}_3): \\ &= :X(z_1, \bar{z}_1)X(z_2, \bar{z}_2)X(z_3, \bar{z}_3): + \\ &- \log|z_1 - z_2|^2 :X(z_3, \bar{z}_3): \\ &- \log|z_1 - z_3|^2 :X(z_2, \bar{z}_2): \\ &- \log|z_2 - z_3|^2 :X(z_1, \bar{z}_1): \quad . \end{aligned} \quad (14)$$

To shorten the next example we introduce the notation  $X_i := X(z_i, \bar{z}_i)$  and  $z_i - z_j := z_{ij}$ . Furthermore we do not write out the single contractions as they should be clear from the last example.

$$\begin{aligned} X_1 X_2 X_3 X_4 &= :X_1 X_2 X_3 X_4: + \sum_{\text{singles}} + :X_1^\bullet X_2^\bullet X_3^\bullet X_4^\bullet: + :X_1^\bullet X_2^\bullet X_3^\bullet X_4^\bullet: + :X_1^\bullet X_2^\bullet X_3^\bullet X_4^\bullet: \\ &=: X_1 X_2 X_3 X_4: + \sum_{\text{singles}} + \log|z_{12}|^2 \log|z_{34}|^2 + \log|z_{13}|^2 \log|z_{24}|^2 + \log|z_{14}|^2 \log|z_{23}|^2 \end{aligned} \quad (15)$$

<sup>1</sup>The symmetric and traceless part of  $\epsilon_{\mu\nu}$  is the fluctuation  $h_{\mu\nu}$  of the metric  $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ .