


Neutrino Physics Course

Lecture XXI

7/7/2023

L MU

Summer 2023



Type II seesaw: a variation on a theme

Aim: a theory of
neutrino mass

- in SM $m_\nu = 0$
- in SM we know the origin
of charged fermion masses:
Higgs mechanism

- Question of u_1 is a window into BSM

theory of mass =

= mass related to new physical processes



Example: LRSDs

self-contained, predictive
theory of both origin and

nature of m_ν

$$\nu_L \xleftrightarrow{LR} \nu_R$$

$$\begin{pmatrix} \nu \\ e \end{pmatrix}_L \xleftrightarrow{LR} \begin{pmatrix} \nu \\ e \end{pmatrix}_R$$

\Downarrow

$$m_\nu \neq 0$$

$$N_i \equiv C \bar{\nu}_R^T \dots$$

$$M_N \propto M_R \text{ (heavy)}$$

with

$$M_D = -M_D^T \frac{1}{M_N} M_D$$

See saw mechanism: Type I

LRSU: $M_D = M_D^T$ (LR=C)

$M_D = M_D^+$ (LR=P)

↓

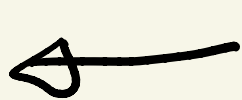
$M_D =$ determined as a f-m
of M_D, M_N

$$M_N = V_R^T u_N V_R$$

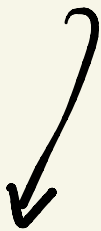
$$M_\nu = V_L^* m_\nu V_L^t$$



ν oscillates, $0\nu 2\beta$, ---



LHC

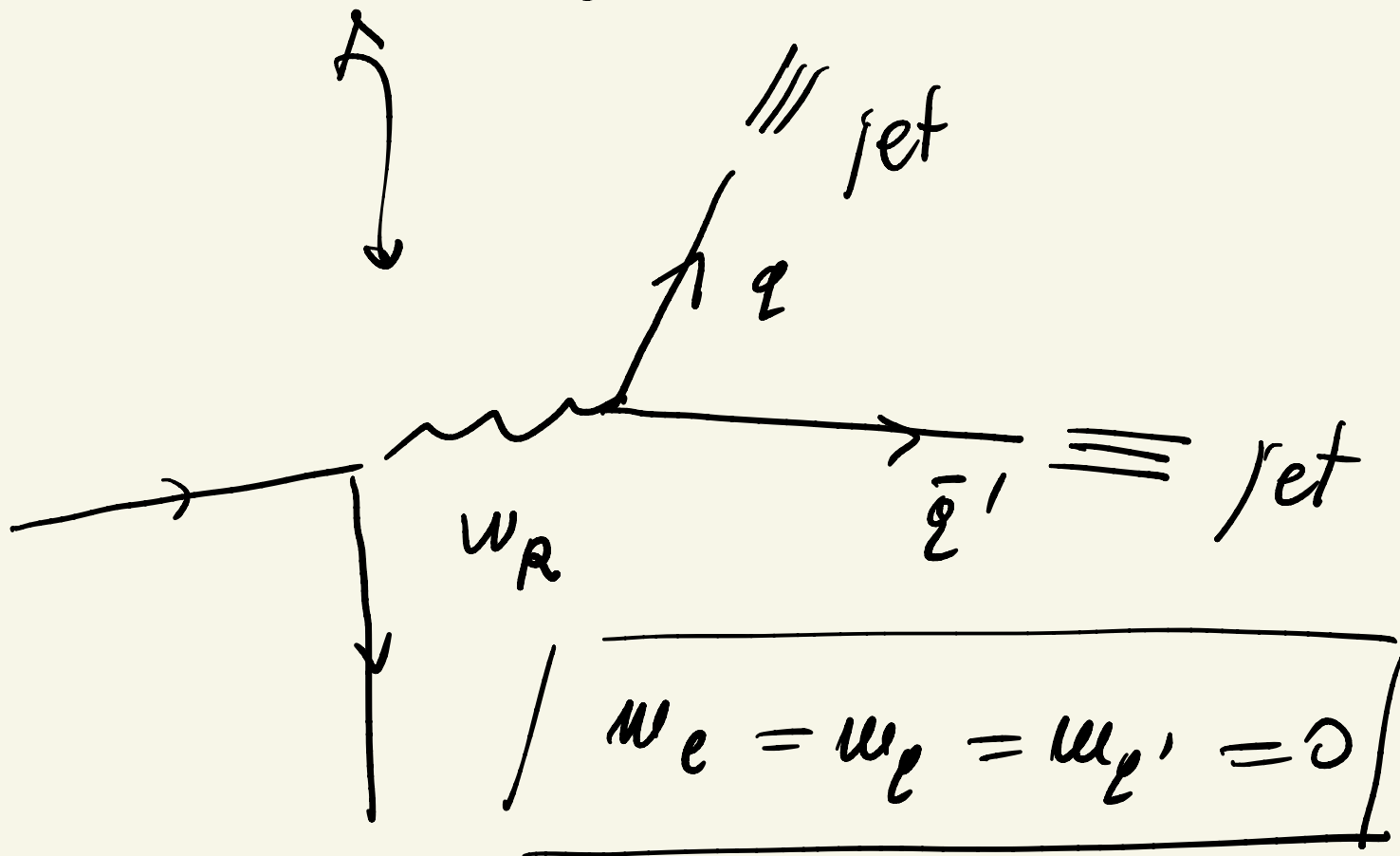


$$\frac{g}{\sqrt{2}} \bar{N} \gamma^\mu e W_{R\mu}^+$$

produced through W_R

and N decays : ← discussed

- $N \rightarrow e_L W_L^+$ (Θ_{LN})
- $N \rightarrow e_R jj$ (W_R)



- $N \rightarrow e_R W_L^+$ ($W_L - W_R$ mixing)
- ?

$$\frac{1}{2} |D_\mu \Phi|^2$$

$$\left. \begin{array}{l} N \rightarrow e + W^+ \\ N^c \rightarrow \bar{e} + W^- \end{array} \right\} \begin{array}{l} N = N^c \\ \text{Mejorona} \end{array}$$

$$\left(\begin{array}{l} N = N_L + C \bar{N}_L^T \\ = N_L + (N^c)_R \end{array} \right)$$

Mejorona

$$N = N^c$$



$$\begin{array}{l}
 N \rightarrow e + w^+ \\
 \rightarrow \bar{e} + w^-
 \end{array}
 \left. \vphantom{\begin{array}{l} N \\ \rightarrow \end{array}} \right\} \begin{array}{l} \text{equal} \\ \text{rates} = \\ \text{Majority} \end{array}$$

Keung, G.S. 1983

$$\textcircled{*} \quad LA = C$$

$$M_D = i M_N \sqrt{\frac{1}{M_N} M_\nu}$$

$$\Theta_{\nu N} = i \sqrt{\frac{1}{M_N} M_\nu}$$



$N \rightarrow e_L + W^+$ predicted



ignore LRSM,
but keep N

seesaw (type I)

or: I don't care about
the theory of ν mass

↓ instead

$$SM + \text{stuff} =$$

$$= \text{neutrino mass}$$

$$\text{stuff} = ?$$

$$(1) \text{ stuff} = N$$

• one $N \Rightarrow$ one massive ν

• two $N \Rightarrow$ two massive ν

minimum

Glashow,

Yenagide, ...

(800)

how to produce N ?

✓ through $\theta_{\nu N}$

• see saw : $\mu_D \ll \mu_N$

$$\Rightarrow \theta_{\nu N} = \frac{\mu_D}{\mu_N} \ll 1$$



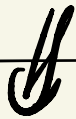
no way of producing N

• $\mu_D = i \sqrt{\mu_N} \quad 0 \quad \sqrt{\mu_D}$

~ not predicted ($0 = \text{arbitrary}$)



we should do better
= verifiable scenario
(model)



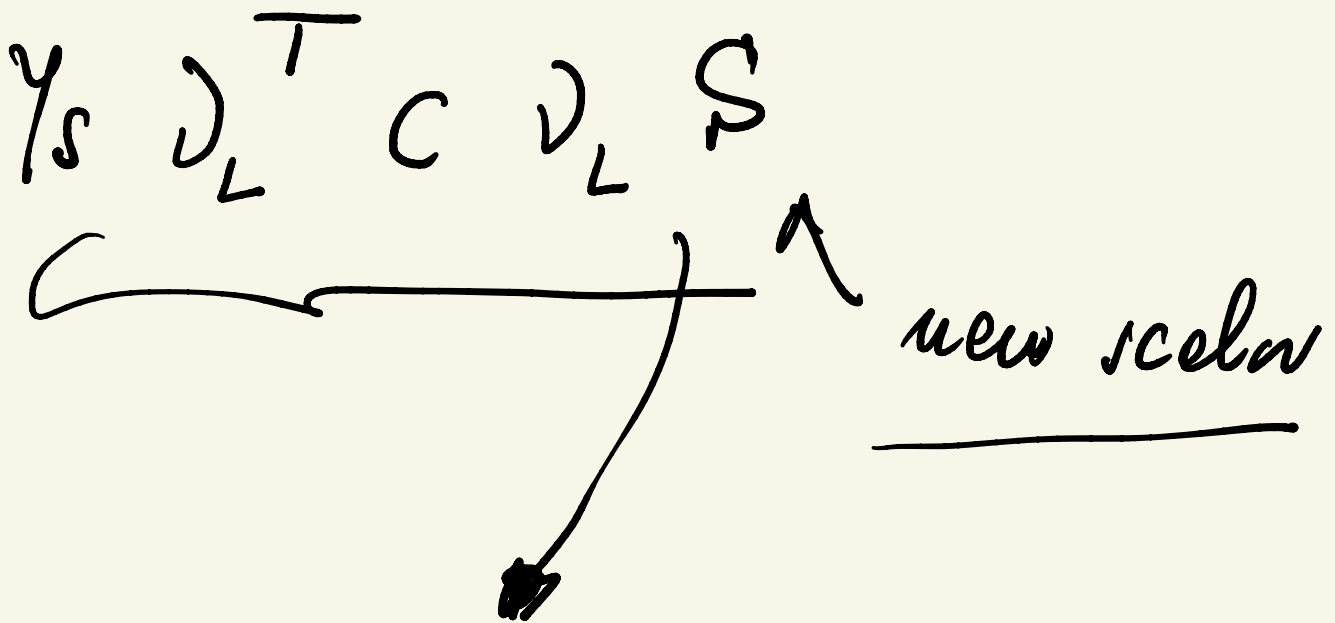
• stuff = scalar



couples directly to v

SM: $l_L = \begin{pmatrix} \nu \\ e \end{pmatrix}_L$, e_R





quantum numbers of S ???

$\Rightarrow S = SU(2)_2$ triplet

$\nu_L \quad C \quad \nu_L \quad S \quad \gamma_s$

$T_3 : \frac{1}{2} + \frac{1}{2}$

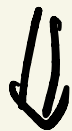
$\underbrace{\hspace{10em}}_1$

\uparrow
 -1

$$Y(s) = +2, \text{ nucle}$$

$$Y(\nu_L) = Y(\ell_L) = -1$$

$$Q \begin{matrix} \uparrow \\ \uparrow \end{matrix} = T_3 + \frac{Y}{2}$$



$S \subseteq \Delta \quad (Y=2 \text{ triplet})$

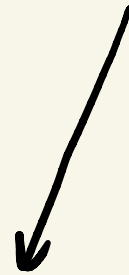


{ Δ_L of LRSM }

↓ Type II error

to be explained

$$\Delta \mathcal{L}_y(\Delta) = l_L^T C' \Delta Y_{\Delta} l_L$$



$$l_L^T C' Y_{\Delta} U^T \underbrace{U \Delta U^T}_{\text{adjoint}} U l_L$$

$\neq \text{inv.}$

as expected!

correct \Downarrow anti-symmetrise

$$\Delta \mathcal{L}_y(\Delta) = \ell_L^T \gamma_\Delta C i\sigma_2 \Delta \ell_L$$

$$\begin{aligned} &\rightarrow \ell_L^T \gamma_\Delta C U^T i\sigma_2 U \underbrace{\Delta U^+ U}_{=} \ell_L \\ &= \ell_L^T \gamma_\Delta C i\sigma_2 U^+ U \Delta U^+ U \ell_L \\ &= \ell_L^T \gamma_\Delta C i\sigma_2 \Delta \ell_L = \text{inv.} \quad \checkmark \end{aligned}$$

Q. E. D.

$$\Delta = \begin{pmatrix} \delta^+ & \delta^{++} \\ \delta^0 & -\delta^+ \end{pmatrix}$$

Exp. (LHC) : $\left[m_{\delta^{++}} \gtrsim 400 \text{ GeV} \right]$

but:

split of masses in $\Delta \simeq 0 (\text{TeV})$

\Downarrow

$\left[m_{\delta^+}, m_{\delta^0} \gtrsim 400 \text{ GeV} \right]$

\Downarrow

$\Delta \neq \text{Higgs field}$



$$\bar{V}_\Delta = \frac{1}{2} \mu_\Delta^2 T_\Delta \Delta^+ \Delta +$$

$$+ \frac{1}{4} \lambda_\Delta (T_\Delta \Delta^+ \Delta)^2 + \dots$$

$$\boxed{\mu_\Delta^2 > 0}$$

$$\left(\lambda_\Delta > 0 \right. \\ \left. \parallel \right. \\ \left. V = \text{bounded} \right)$$



$$\langle \Delta \rangle = 0 !$$

$$\frac{\partial V}{\partial \langle \Delta \rangle} = \mu_\Delta^2 \langle \Delta \rangle + \lambda_\Delta \langle \Delta \rangle^3 = 0$$

$$\parallel = \langle \Delta \rangle (\mu_\Delta^2 + \lambda_\Delta \langle \Delta \rangle^2)$$

$\overbrace{\hspace{10em}} > 0$

$$\Rightarrow \langle \Delta \rangle = 0$$

$$\Rightarrow \mu_\nu = 0 \quad !! \quad ??$$

bwt

$$V = \bar{V}_\phi + \bar{V}_\Delta + \bar{V}_{\phi\Delta}$$

$$\bar{V}_{\phi\Delta} = \mu \phi^\top i \sigma_2 \Delta^\dagger \phi$$

$$\gamma(\phi) = +1$$

$$\begin{pmatrix} \Delta \rightarrow U \Delta U^\dagger \\ \Delta^\dagger \rightarrow U \Delta U^\dagger \end{pmatrix}$$

⇓ next step

$$\langle \Delta \rangle = ?$$

$$\Delta = \begin{pmatrix} f^+ & f^{++} \\ f^0 & -f^+ \end{pmatrix}$$

$$\Delta^\dagger = \begin{pmatrix} f^- & f^{0*} \\ f^{--} & -f^- \end{pmatrix}$$



$$\langle \Delta^+ \rangle = \begin{pmatrix} 0 & v_\Delta \\ 0 & 0 \end{pmatrix}$$

$$\langle \bar{V}_{\phi_0} \rangle = \mu \langle \phi \rangle^T i \sigma_2 \langle \Delta^+ \rangle \langle \phi \rangle$$

$$= \mu \begin{pmatrix} 0 & v_{SM} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 0 & v_\Delta \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ v_{SM} \end{pmatrix}$$

$$= \mu \begin{pmatrix} 0 & v_{SM} \end{pmatrix} \begin{pmatrix} 0 & 0 \\ -v_\Delta & 0 \end{pmatrix} \begin{pmatrix} 0 \\ v_{SM} \end{pmatrix}$$

$$= -\mu v_{SM}^2 v_\Delta$$



$$\langle V \rangle = \langle V_{\phi_A} \rangle + \langle V_{\phi} \rangle + \langle V_H \rangle$$

$$\Rightarrow \langle V \rangle = \dots - \mu v_{SM}^2 v_{\Delta}$$

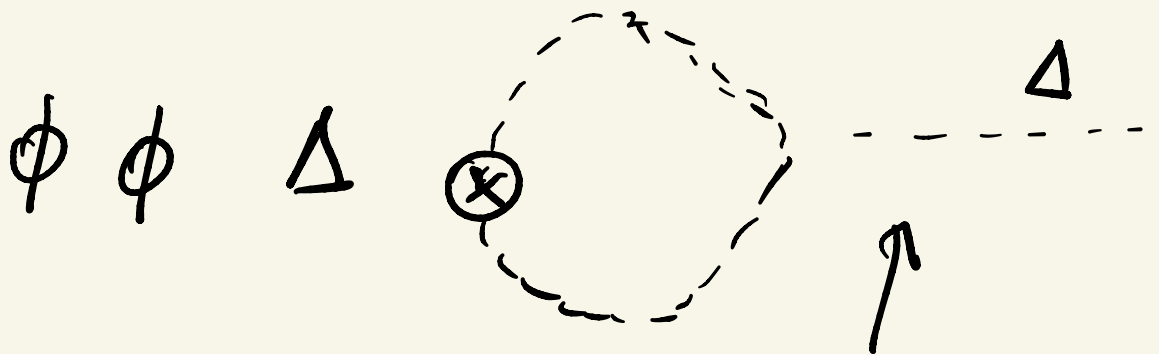
$$+ \frac{1}{2} \mu_{\Delta}^2 v_{\Delta}^2 + \frac{\lambda}{4} v_{\Delta}^4$$

$$\frac{\partial \langle V \rangle}{\partial v_{\Delta}} = -\mu^2 v_{SM}^2 + \mu_{\Delta}^2 v_{\Delta} + \lambda v_{\Delta}^3 = 0$$

$$\boxed{v_{\Delta} \neq 0} \quad (\text{we made it})$$



linear in Δ term



tadpole



(i) $v_D \neq 0$

(ii) $v_D \ll v_{SM}$ ($m_D^2 \gg m_h^2$)

$$M_W^2 = \frac{1}{4} g^2 v_{SM}^2 + g^2 v_\Delta^2$$

$$M_Z^2 = \frac{1}{4} (g^2 + g'^2) v_{SM}^2 + 2(g^2 + g'^2) v_\Delta^2$$

but

$$M_Z^2 \approx \frac{g^2 + g'^2}{g^2} M_W^2 \quad (\text{exp.})$$

$$\Rightarrow v_\Delta \ll v_{SM}$$

$$v_\Delta \lesssim \text{GeV} \quad v_{SM} \approx 200 \text{ GeV}$$

}

↓ compute v_Δ

$$-\mu v_{SM}^2 + \mu_\Delta^2 v_\Delta + \lambda_\Delta v_\Delta^3 = 0$$

neglect

$$\lambda_\Delta \lesssim 1 \Rightarrow$$

$$\lambda v_\Delta^3 \lesssim v_\Delta^3 \ll \mu_\Delta^2 v_\Delta$$

⇓

$$v_\Delta \approx \mu \frac{v_{SM}^2}{\mu_\Delta^2}$$

$\mu \approx m_\Delta$ for illustration

$$v_\Delta = \frac{v_{SM}^2}{m_\Delta}$$



$$m_\nu = Y_\Delta v_\Delta \approx Y_\Delta \frac{v_{SM}^2}{m_\Delta}$$

$$m_\Delta \rightarrow \infty \Rightarrow m_\nu \rightarrow 0$$

⇒ Type II seesaw

analog ⇕ Type I seesaw

$$m_\nu = \frac{m_D^2}{m_N} \approx Y_0^2 \frac{v_{SM}^2}{m_N}$$

$$(m_D = Y_0 v_{SM})$$

↓ new generation

$$M_\nu = Y_\Delta v_\Delta$$



$$\Delta \mathcal{L}_\nu (\Delta) =$$

$$\begin{pmatrix} v^T & e^T \\ \hline \end{pmatrix}_L C Y_{\Delta} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} \delta^{++} & \delta^{++} \\ \delta^0 & -\delta^+ \end{pmatrix} \begin{pmatrix} v \\ e \end{pmatrix}_L$$

$$= \begin{pmatrix} v^T & e^T \\ \hline \end{pmatrix}_L C Y_{\Delta} \begin{pmatrix} \delta^0 & -\delta^+ \\ -\delta^+ & -\delta^{++} \end{pmatrix} \begin{pmatrix} v \\ e \end{pmatrix}_L$$

$$= \begin{pmatrix} v^T & e^T \\ \hline \end{pmatrix}_L C Y_{\Delta} \begin{pmatrix} \delta^0 v & -\delta^+ e \\ -\delta^+ v & -\delta^{++} e \end{pmatrix}_L$$

$$= v_L^T C Y_{\Delta} v_L \delta^0 - v_L^T C Y_{\Delta} e_L \delta^{++} - e_L^T C Y_{\Delta} v_L \delta^+ + e_L^T C Y_{\Delta} e_L \delta^{++}$$

$$\left\{ \begin{array}{l} - \left(e_L^T C Y_{\Delta} \delta^{++} e_L \right) \end{array} \right.$$

⇓

$$\downarrow$$

$$M_\nu = Y_\Delta \nu_\Delta +$$

$$- e_L^T C \frac{-M_\nu}{\nu_\Delta} e_L \delta^{++}$$

neutrino mass matrix

$$M_\nu = V_\ell^* m_\nu V_\ell^t = \text{measured}$$

$$\Downarrow$$

$$\delta^{--} \rightarrow l_i + l_i \quad (-M_\nu)_{ij}$$



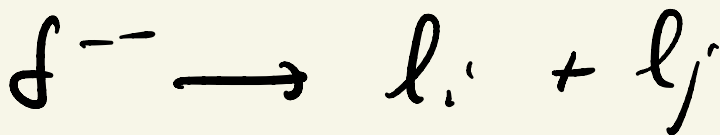
charged leptons

overall decay rate depends

on V_D

but

branching ratios = determined

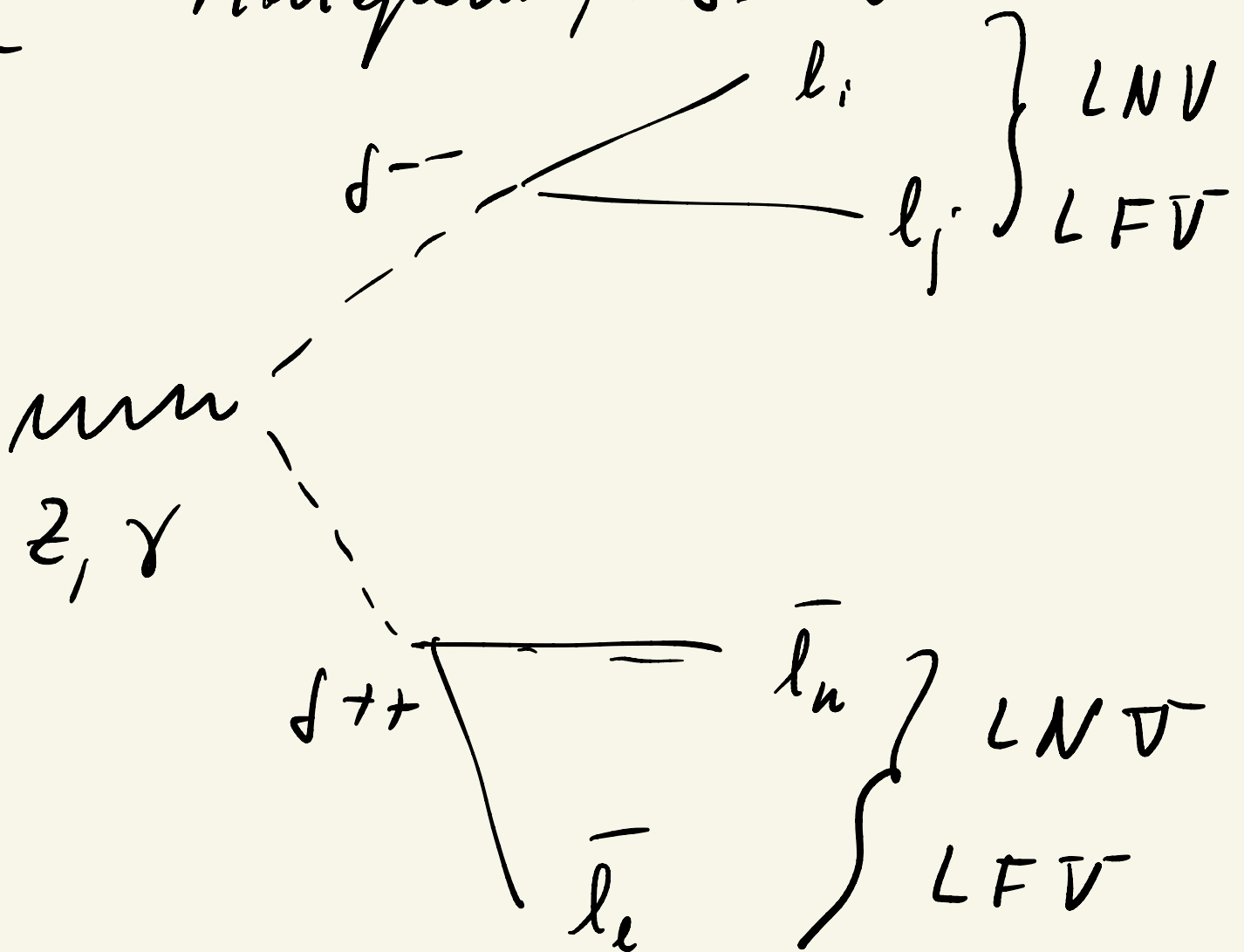


Schnetz, ? 2008

Type II seesaw = 1980-1981

Maggi, Veltman
Lazarides et al

Mohapatra, G.S.



on mixings

quarks: V_{CKM}

$$\begin{array}{ccc} \theta_{12}^q & , & \theta_{13}^e & , & \theta_{23}^q \\ \parallel & & \parallel & & \parallel \\ 13^\circ & & 10^{-3} & & 10^{-2} \end{array}$$

leptons: V_{PMNS}

$$\begin{array}{ccc} \theta_{12}^l & , & \theta_{13}^l & , & \theta_{23}^l \\ \parallel & & \parallel & & \parallel \\ 30^\circ & & 10^\circ & & 45^\circ \end{array}$$

$$M_\nu = V_{PMNS}^* M_\nu V_{PMNS}^\dagger$$

$$LFV = \text{large}$$

μ vs m_Δ ?

$$M_\nu = Y_\Delta \nu_\Delta$$

$$\left\{ \begin{array}{l} \nu_\Delta = \mu \left(\frac{v_{SM}}{m_\Delta} \right)^2 \leq \mu/4 \\ (Y_\Delta \leq O(1)) \end{array} \right.$$

$$m_\nu \leq \nu_\Delta \Rightarrow$$

$$10^{-1} \text{eV} \leq \nu_\Delta \leq \text{GeV}$$

$\Delta \neq \text{Higgs}$

Higgs ϕ

$$V = -\frac{\mu^2}{2} \phi^2 + \frac{\lambda}{4} \phi^4$$

$$\Rightarrow \langle \phi \rangle = \mu^2 / \lambda$$

$$m_\phi^2 = 2\lambda \langle \phi \rangle^2$$

$$m_\phi \simeq \langle \phi \rangle$$

• Δ : $\mu_{\Delta} \gtrsim 400 \text{ GeV}$

$$v_{\Delta} \leq 6 \text{ GeV}$$



$\Delta \neq \text{Higgs}$



$$\langle \Delta \rangle = \mu \frac{v_{SM}^2}{\mu_{\Delta}^2}$$

"Higgsing" Δ



$$\overline{V}_\Delta = -\mu_\Delta^2 \Delta^2 + \lambda_\Delta \Delta^4$$

↓

$$m_H \simeq v_\Delta \leq O(\text{GeV})$$

$$\Rightarrow m_{f^{++}} \leq \text{GeV}$$

~~exp.~~

wrong