

Neutrino Physics Course

Lecture VI

12 / 5 / 20 23

L M U

Spring 2023



SM Higgs mechanism:

the Weinberg model

$$G_{SM} = SU(2)_L \times U(1)_Y$$



Glashow '61

$$\exists A \longleftrightarrow Q_{em} = Q_A (e)$$



$$\exists Z \longleftrightarrow Q_Z = T_3 - \sin^2 \theta_w Q_{em}$$

$$\left(\frac{g}{\cos \theta_w} \right)$$

$$e = \sin \theta_w g$$

$$\tan \theta_w = \frac{g'}{g}$$

⇓ renormalisability

does not work!

• $M_A \neq 0$

⇓

$$\Delta_{\mu\nu}(A) = \frac{-i}{q^2 - M_A^2} \left[g_{\mu\nu} - \frac{q_\mu q_\nu}{M_A^2} \right]$$

BAD
(divergencies)

V (1) gauge theory Golobtune '61

$$\mathcal{L} = \frac{1}{2} (D_\mu \phi)^* (D^\mu \phi) - V(\phi)$$

$$\phi \rightarrow e^{i\alpha Q} \phi = e^{i\alpha} \phi$$

$(\alpha = \alpha(x))$

$$D_\mu = \partial_\mu - i g Q A_\mu$$

(source
of mass)

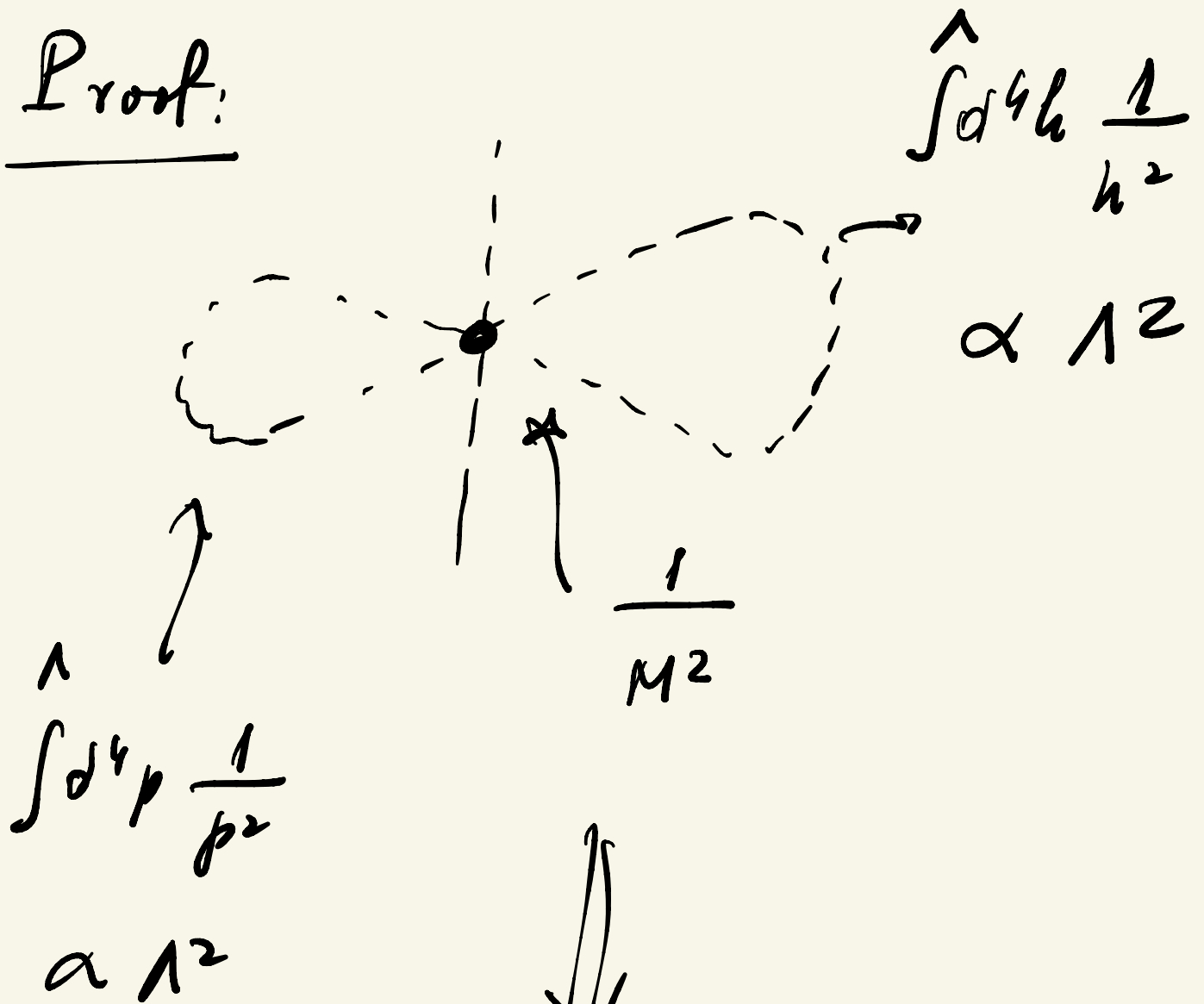
$$V = \frac{\lambda}{4} (|\phi|^2 - v^2)^2 \quad (1)$$

$$+ \frac{(\phi^* \phi)^3}{M^2} \quad (?)$$

Q. why not?

A. Lose renormalizability.

Proof:



$$A(2) \propto \frac{\Lambda^4}{M^2}$$

Q.E.D.



$$M_0 = \left\{ \phi_0 \therefore V = V_{\min} = 0 \right\}$$

$$\parallel \\ \text{vacuum} = \left\{ |\phi_0|^2 = v^2 \right\} = S_1$$

$$(1) \phi_0 = v$$

$$(2) \phi = v + h + i G$$

$$\Rightarrow \begin{array}{l} m_h^2 = 2\lambda v^2 \\ m_G = 0 \end{array}$$

Neutrino - Goldstone
boson

$$V(h, \phi) = \dots + \frac{\lambda}{4} (h^4 + \phi^4) + \lambda v h^3$$

↑
no symmetry

⇓

SSB

$$\Phi_0 \therefore Q \Phi_0 \neq 0$$

$$(Q \phi = \phi) \quad \Downarrow$$

$$U \Phi_0 \neq \Phi_0$$

$v = \text{measure of } \mathbb{S}^2 \mathbb{B}$

$$u_n \propto v !$$



True ~~even~~ if $\alpha \neq \alpha(x)$

(Global symmetry)



$$\alpha = \alpha(x)$$

gauge symmetry

$$\mathcal{L}_{\text{kin}} = \frac{1}{2} |D_\mu \phi|^2$$

$$\rightarrow \sum \frac{1}{2} |D_\mu \phi_0|^2 = \frac{1}{2} |-ig A_\mu Q \phi_0|^2$$

$$(Q \phi_0 \neq 0)$$

$$= \frac{1}{2} |-ig A_\mu v|^2$$

$$= \frac{1}{2} g^2 v^2 A_\mu A^\mu$$



$$M_A = gv$$

Higgs
mechanism

• global case } $A + \phi \in C \Rightarrow 4 \text{ d.o.f.}$
 (2) (2) \uparrow
 ($u_A = 0$) ↓
 (h, G)

• local case } $A + \phi \in C \Rightarrow 4 \text{ d.o.f.}$
 (2) (2)
 ($u_A = 0$ (initially))

final (SSB) : $A (u_A \neq 0)$
 (3)
 + ϕ (2)
 ↑
 (h, G) 5 d.o.f.

Unitary gauge

$$\phi = v + h + iG \quad (2)$$

↑ why not?

$$\Rightarrow \phi = e^{iG/v} (v + h) \quad (3)$$

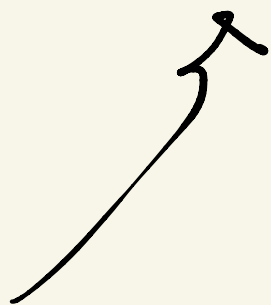
$$= v + h + iG + \dots$$

$$(3) \Rightarrow \phi \rightarrow e^{i\alpha(x)} \phi$$

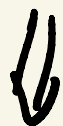
$$= e^{i \left(\alpha(x) + \frac{G(x)}{v} \right)} (v + h(x))$$

⊙ // choice of $\alpha(x)$

$$\Rightarrow \boxed{\phi \rightarrow \phi_{un} = v + h}$$



I gauged ϕ away!



massless

in.

$$A(z), \phi(z) \Rightarrow 4 \left. \vphantom{A(z)} \right\} \text{d.o.f.}$$

f.

$$A(z), h \Rightarrow 4$$



massive

Trouble

find theory:

$$h \leftrightarrow m_h = v \sqrt{2\lambda}$$

$$A \leftrightarrow M_A = g v$$

(Proca)

(no gauge symmetry)

↓

$\Delta_{\mu\nu}^{(P)}$ is bad

| A_μ (4 d.o.f.) i.e. $\partial_\mu A^\mu = 0$
 \Rightarrow 3 physical

Unitary gauge

only physical degrees
of freedom

Not QED



A_μ (4 d.o.f.)

$$A_\mu \rightarrow A_\mu + \frac{1}{e} \partial_\mu \alpha(x)$$

$$\mathcal{L} + \mathcal{L}_{g.f.} = -\frac{1}{23} (\partial_\mu A^\mu)^2$$

↑ gauge parameter

QED $\therefore A_\mu + g_{\text{phys}}$
 $4^L - 2 = 2$
 d.o.f. physical

SSB QED $\therefore (A_\mu, \psi) + g_{\text{phys}}$
 $\overset{\rightarrow}{\psi} - 2 = 3$
 d.o.f. physical

↓

$\phi = v + h + iG \leftarrow e^{iG/v} (v+h)$
 \uparrow
not physical

$D_\mu \phi = i\gamma_\mu G - ig A_\mu v$

$$\frac{1}{2} |D_\mu \phi|^2 = \frac{1}{2} (\partial_\mu \phi - g A_\mu \phi)^2$$

$$= \frac{1}{2} (\partial_\mu \phi)^2 + \frac{1}{2} g^2 \phi^2 A_\mu A^\mu$$

$$- g (\partial_\mu \phi) A_\mu \phi$$

// Integrate by part

$$g \phi \partial_\mu A_\mu$$

~~$$= M_A \phi (\partial_\mu A_\mu)$$~~ not good



$$\mathcal{L}_{gf} = -\frac{1}{2\xi} (\partial_\mu A^\mu - \xi M_A \phi)^2$$

$$= -\frac{1}{2\xi} (\partial_\mu A^\mu)^2 - \frac{1}{2} \xi M_A^2 \phi^2$$

$$\underline{\underline{-\mu_A G (2\mu_A/\mu)}}$$

⇓

$$\mathcal{L}_G = \frac{1}{2}(\dot{\mu}_G)^2 - \frac{1}{2}3\mu_A^2 G^2$$

$$\Rightarrow \boxed{\mu_G = \sqrt{3} \mu_A}$$

$G \neq \text{physical}$

(gauge dep. mass)

⇓

$$D(\omega) = \frac{i}{\omega^2 - 3\mu_A^2} \quad (4)$$

$$\Delta_{\mu\nu}(A) = \frac{-i}{\ell^2 - m_A^2} \left[g_{\mu\nu} + (\beta - 1) \frac{\partial_\mu \partial_\nu}{\ell^2 - m_A^2} \right]$$

(5)

⇓

$$D(\partial) \xrightarrow{\ell \rightarrow \infty} \frac{1}{\ell^2}$$

$$\Delta_{\mu\nu}(A) \xrightarrow{\ell \rightarrow \infty} \frac{1}{\ell^2}$$

} for any finite }

unitary gauge?

⇓

} → ∞



$$D(\phi) \rightarrow 0 \quad (\text{no } \phi)$$

$$\Delta_{\mu\nu}(A) \rightarrow \frac{-1}{\xi^2 - m_A^2} \left[g_{\mu\nu} - \frac{\xi_\mu \xi_\nu}{m_A^2} \right]$$

Proca!

Cross of it all:

$$Q \Phi_0 \neq 0$$

$$Q \phi_0 = \xi \phi_0$$



$$\boxed{|m_A| = |g \xi| = \xi \xi |\phi_0|}$$

global case

$$\mathbb{Q} \phi_0 \neq 0$$

$$\phi = a + b + i c$$

↓

$$\boxed{\mu_G = 0}$$

$N \in$
locus

local case

$$\mathbb{Q} \phi_0 \neq 0$$

$$\Rightarrow \mu_A \neq 0$$

but $N \in$ is "eaten" by A

but we cannot do loops
without \in

⇓

$$\phi \rightarrow e^{i d(x)} \phi \quad \therefore \phi_{\text{em}} = (v+h)$$

($G = \text{gauge group}$)



$$M_G \rightarrow \infty \quad \text{in} \quad \phi_{\text{ren}} = v+h+iG$$

(renormalizable)

Choosing variables

global case $d = \text{const.}$

$$\phi_{un} = e^{i\sigma/a} (a+h)$$

$$\Rightarrow |\phi_{un}|^2 = (a+h)^2$$

$$\Rightarrow \mu_G = 0 \quad (G \text{ not in } V)$$

$$\partial_\mu \phi_{un} = \partial_\mu h e^{i\sigma/a} +$$

$$(a+h) \partial_\mu i\sigma/a e^{i\sigma/a}$$

$$= (\partial_\mu h + i \partial_\mu \sigma (1+h/a)) e^{i\sigma/a}$$



non-trivial int.