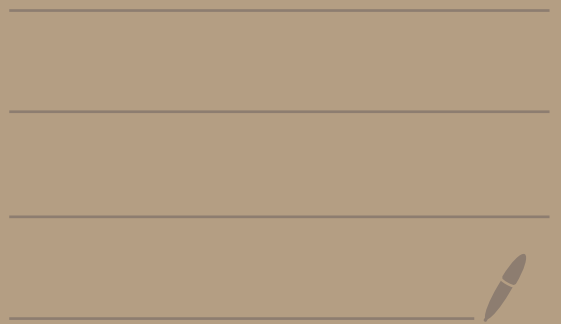


Why neutrons?

L M U Neutron Course

Spring 2021

13/4/2021



- cool - about;

$$\text{Sun} \rightarrow 10^{10} / \text{cm}^2 \text{ sec} \quad \text{neutrons}$$

mean free path:

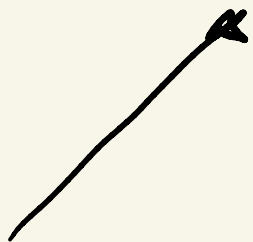
$$\lambda_{\text{nu}} \sim 1 \text{ m}$$

$$\lambda_e \approx 1 \text{ cm} \quad (\text{electron})$$

$$\lambda_{\nu} \approx 10^{20} \text{ cm}$$

$$E_{\nu} \approx 1 \text{ MeV}$$

$$(m_e \approx 0.5 \text{ MeV})$$



open the door to
new physics



missing energy

$E_e = \text{fixed}$



Postcard:

Pauli 1930

$\exists \nu_e$ ($Q_e = 0$)

neutral "light" particle

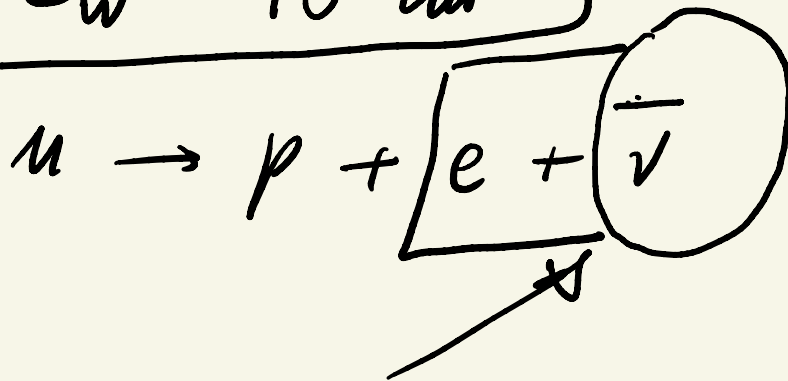
Fermi theory

'1934

"effective" theory of
weak interaction

$$\sigma_w \approx 10^{-44} \text{ cm}^2$$

$$E = 1 \text{ MeV}$$



(\rightarrow decay)

$e = \text{beta}$

anti-neutrino

$$e = \text{lepton}$$

(lepton = fermion)

$$p, n = \text{hadrons}$$

(hadrons = strong,
heavy)

$$\text{Lepton (L) \#} = \text{conserved}$$

$$(def.) \quad L(\nu) = 1$$

$$u \rightarrow p + e + \bar{\nu}_e$$
$$\Delta B \neq 0 \quad \Delta L \neq 0$$

$u, p =$ baryons (fermions) $s = 1/2$

$\pi, \mu =$ mesons (bosons) $s = 0$

• QED $e + e \rightarrow e + e$

$$\sigma_{QED} \approx 10^{-22} \text{ cm}^2 \quad E \approx \text{MeV}$$

Pauli: forgive me, I have sinned!

Cowan, Reines 1956

Pauli: "Everybody comes to him
who knows how to wait."

reactor

Pontecovo

'40s

$$\bar{\Phi} (\text{flux}) \approx 10^{13} / \text{cm}^2 \text{sec}$$

1956

bomb shell



Lee, Yang (theory)

Wu et al (exp)

$P : L \leftrightarrow R$

\neq maximally

Lederman,

Garwin

• only ν_L exist

weak int. : only LH fermions

$\psi(e) : \psi_L, \psi_R$

$$\psi_{L,R} \equiv L(R) \psi$$

$$L \equiv \frac{1 + \gamma_5}{2}$$

(R)

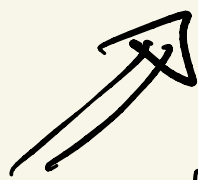
QED

$$e A_\mu [\bar{e}_L \gamma^\mu e_L + \bar{e}_R \gamma^\mu e_R] \quad (E)$$

weak

$$g W_\mu^+ [\bar{\nu}_L \gamma^\mu e_L + \bar{e}_L \gamma^\mu \nu_L] \quad (F)$$

1983
CERN



V-A

Marshall, Sudshan
1957

- W boson = messengers of weak int
- ↑ A photon = -11- em int.

$SU(2) \times U(1) = \text{ew S.M}$

↓

Feinman 1961

Standard Model 1967

Higgs ← Weinberg

$$\boxed{u_1 = 0}$$

blissly
⇓
"cure"

⇓
S.D.I is incomplete
(only incompleteness)

⇓
(the) door to New Physics (NP)

gauge theory \Leftrightarrow gauge bosons
as messengers of force

$$U_{em}^{(1)} = QED$$

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{f} \gamma^\mu D_\mu f - m \bar{f} f \quad (1)$$

$$\partial_\mu \quad (f \equiv \psi_f)$$

$$D_\mu = \partial_\mu - ie A_\mu \quad (2)$$

$$f \rightarrow e^{i\alpha(x) Q_{em}} f \quad (3)$$

$$Q_{em} f = Q_f f$$

↙ em charge

$$Q_e = -1, \quad Q_p = +1, \quad Q_n = 0$$

$$Q_\nu = 0$$

$$\rho_e + \rho_p = 0 \quad (\leq 10^{-40})$$

(neutrality of universe)

$$A_\mu \rightarrow A_\mu - \frac{1}{e} \partial_\mu \alpha(x) \quad (4)$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \quad (5)$$

$$F_{0i} \equiv E_i$$

$$F_{ij} = \varepsilon_{iju} B_u$$



$$e A_\mu \bar{\psi} \gamma^\mu Q_{em} \psi \quad \alpha = \alpha(x)$$



$$\mathcal{L}_{\text{Dirac}} = i \bar{f} \gamma^\mu \partial_\mu f - m \bar{f} f$$

$$f \rightarrow e^{i \alpha Q} f \quad \alpha = \text{const.}$$

↓ generalization

$$G_{\text{SM}}(\text{ew}) = SU(2)_L \times U(1)_Y$$

↓ (to be done)

$$W_\mu \left(\begin{pmatrix} \nu \\ e \end{pmatrix}_L \right) \rightarrow U \left(\begin{pmatrix} \nu \\ e \end{pmatrix}_L \right)$$

flavor =
 ν, e

$$U^\dagger U = 1, \quad \det U = 1$$

$$\mathcal{L}_{SM} = i \bar{f} \gamma^\mu D_\mu f \quad \text{-----}$$

$$D_\mu = \partial_\mu - ig T_a A_\mu^a \quad \text{-----}$$

$$U = e^{i \theta_a(x) T_a}$$

Euler

$$[T_a, T_b] = i \epsilon_{abc} T_c$$

$$T_a \equiv \frac{\sigma_a}{2} \quad a = 1, 2, 3$$

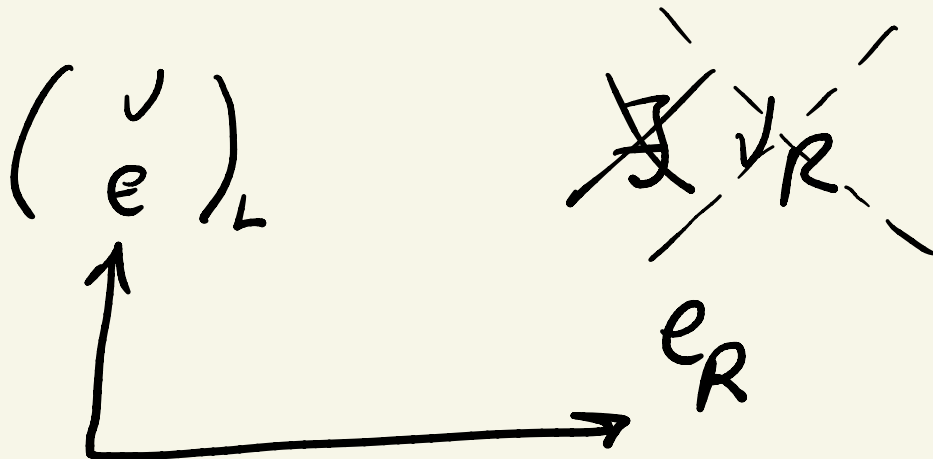
$$A_\mu^1, A_\mu^2, A_\mu^3$$

$$W_\mu^\pm = \frac{A_\mu^1 \mp i A_\mu^2}{\sqrt{2}}$$

$$\left[\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \right]$$

$$\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

SM



$$A_\mu \quad (QED)$$

$$\bar{f} \equiv f^\dagger \gamma^0$$

mass $m \bar{f} f = m f^\dagger \gamma_0 f$
 $= m(\bar{f}_L f_R + \bar{f}_R f_L)$

• $\{ \gamma^\mu, \gamma^\nu \} = 2 g^{\mu\nu}$

$g^{\mu\nu} \equiv \eta^{\mu\nu} = \text{diag}(1, -1, -1, -1)$

$\gamma^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix}, \gamma^0 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$\Sigma_{\mu\nu} = \frac{1}{4i} [\gamma_\mu, \gamma_\nu]$

$[\Sigma_{\mu\nu}, \Sigma_{\alpha\beta}] = \delta_{\mu\alpha} \Sigma_{\nu\beta} - \dots$
 (Lorentz)

$$d=4: \psi (\text{spinor}) \rightarrow \Lambda \psi$$

$$\Lambda \equiv \exp \left(i \theta_{\mu\nu} \Sigma^{\mu\nu} \right)$$

$\begin{matrix} \nearrow & \nearrow \\ 6 & 6 \end{matrix}$

(A5)

$$\gamma_5 \equiv \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \therefore$$

$$\{ \gamma_5, \gamma_\mu \} = 0, \quad \boxed{\gamma_5^2 = 1}$$

⇓

$$[\gamma_5, \Sigma_{\mu\nu}] = 0$$

$$\left(\gamma_5 = \begin{matrix} + \\ - \end{matrix} i \gamma_0 \gamma_1 \gamma_2 \gamma_3 \right)$$

\Downarrow

$$L(R) = \frac{1 \pm \gamma_5}{2}$$

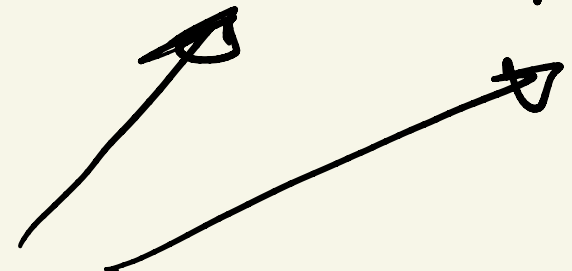
$$[L(R), \Sigma_{\mu\nu}] = 0$$

$$L^2 = L, R^2 = R, LR = 0$$

$$L \equiv \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}, R \equiv \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\psi_D = \psi_L + \psi_R = \begin{pmatrix} u_L \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ u_R \end{pmatrix}$$

(electron)



($D = \text{Dirac}$)

$u_L, u_R = \text{fundamental entities}$

$$\mathcal{L}_{\text{Dirac}} = i \bar{\psi}_0 \gamma^\mu \partial_\mu \psi_0 =$$

$$= i \bar{\psi}_L \gamma^\mu \partial_\mu \psi_L + i \bar{\psi}_R \gamma^\mu \partial_\mu \psi_R$$

L steps L

R steps R

1928

$$\mathcal{L}_0 = \mathcal{L}_{\text{Dirac}} - m \bar{\psi} \psi =$$

$$= i \bar{\psi} \gamma^\mu \partial_\mu \psi - m \bar{\psi} \psi$$

\Downarrow

$$\boxed{i \gamma^\mu \partial_\mu \psi = m \psi} \quad (\text{Dirac equation})$$

\Downarrow

$$\psi(x) = e^{i p \cdot x} \tilde{\psi}(p)$$

$$p \cdot x = p x = p^\mu x_\mu$$

$$\Rightarrow p^\mu \gamma_\mu \tilde{\psi}(p) = m \tilde{\psi}(p) /$$

Isgere

$$p^\mu p^\nu \gamma_\mu \gamma_\nu \tilde{\psi}(p) = m^2 \tilde{\psi}(p)$$

||

$$\frac{1}{2} \{ \gamma^\mu, \gamma^\nu \}$$

⇓

$$p^\mu p_\mu \equiv p^2 = m^2$$

$E^2 = \vec{p}^2 + m^2$

 (def.)

$$\cdot m \bar{f} f = m (f_L^\dagger + f_R^\dagger) \gamma^0 (f_L + f_R)$$

$$= m (f_L^\dagger \gamma^0 f_L + L \leftrightarrow R +$$

$$+ f_L^\dagger \gamma^0 f_R + L \leftrightarrow R)$$

①

$$= m (f^\dagger L \gamma^0 L f + L \leftrightarrow R)$$

$$+ m (f^\dagger L \gamma^0 R f + L \leftrightarrow R) \text{ ②}$$

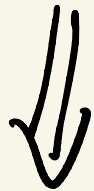
$$= m (f^\dagger \cancel{R L} f + L \leftrightarrow R) \text{ ①}$$

$$+ m (f^\dagger \gamma^0 R R f + L \leftrightarrow R) \text{ ②}$$

$$= m f^\dagger \gamma^0 (R+L) f = m \bar{f} f$$

$$m \bar{f} f = m (\bar{f}_R f_L + \bar{f}_L f_R)$$

Dirac mass



~~$\bar{\nu}_L$~~ ν_R in SM



$$m_\nu = 0$$

SM = a theory of ew processes
and particles that underwrite
them

$$SU_L(2) \times U_Y(1) \equiv G_{SM}$$



particle carry quantum
numbers under G_{SM}

$\nu_R = \text{phantom} -$

- knows nothing about $SU(2) \times U(1)$

(hypercharge) $\left. \begin{array}{l} T_a \nu_R = 0 \\ Y \nu_R = 0 \end{array} \right\} \left(\begin{array}{l} \nu_R \text{ does} \\ \text{not interact} \\ \text{with gauge} \\ \text{bosons} \end{array} \right)$

QED = they of em processes

(not at $v - Q_v = 0$)

↓ they of decay particles

Who is ν_R ?

Sha ↓

leesaw ↓

$(\nu)_L + (e)_R$

$(\nu)_L$ $(e)_L$ ν_R, e_R

weak

$G_{SM} = gauge group$

$G_{SM} = gauge group$

↓

$$g W_\mu^+ \bar{\nu}_R \gamma^\mu e_R = g W_\mu^+ \bar{\nu} \gamma^\mu e$$



we must drop the gauge group !!

Task = theory of neutrino mass

(Predictive, complete)

$$\Downarrow \left[G_{SM} = SU(2)_L \times U(1) \right]$$

$a=1,2,3$
 \Downarrow

$A_\mu^a (T_a)$

\Downarrow

$B(Y)$

G_{SM} = group theory of weak
and em processes

$$\partial_\mu \rightarrow D_\mu = \partial_\mu - ig_1 T_a^i A_\mu^{i a} -$$
$$- ig_2 T_\alpha^2 A_\mu^{2 \alpha} - \dots$$

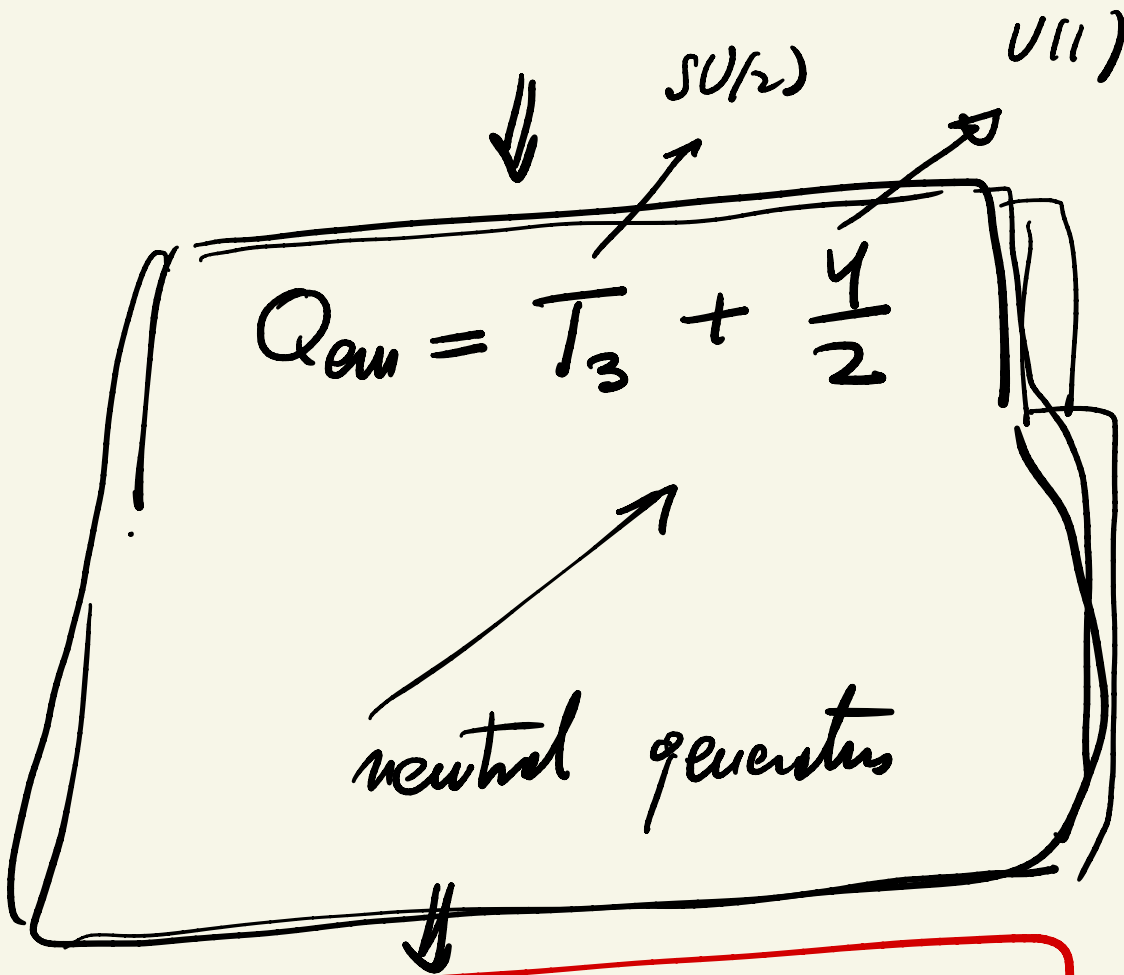
(repeated index =
= summed)

$$\Downarrow \quad W_\mu^\pm = \frac{A_\mu^1 \mp i A_\mu^2}{\sqrt{2}}$$

$A_\mu \Rightarrow$ lives in G_{SM}
= lower constituents of
= $f(A_\mu^a, B_\mu) =$

$$= f(A_\mu^s, B_\mu)$$

$\Leftrightarrow Q_{em} = \text{linear comb. of}$
 T_3, Y



$$Y = 2 [Q_{em} - T_3]$$

input input of theory

leptons: $\begin{pmatrix} \nu \\ e \end{pmatrix}_L$ $\left\{ \begin{array}{l} T_3 \nu_L = \frac{1}{2} \nu_L \\ T_3 e_L = -\frac{1}{2} e_L \end{array} \right.$

↑
doublet

$$T_3 = \frac{1}{2} \sigma_3 = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad \nabla$$

$$[T_a, T_b] = i \epsilon_{abc} T_c$$

$$e_R$$

$$\left. \begin{array}{l} T_3 e_R = 0 \\ Y e_R = -2 e_R \\ e_R = \text{particle} \end{array} \right\}$$

$$\begin{array}{l} Q e_R = -1 \\ Q e_L = -1 \end{array}$$

$$\nu_R : Q \nu_R = 0 \quad \text{neutral}$$

(definition)

mass of neutral particles

(Majorana mass)

$$W \left(\begin{array}{c} \nu \\ e \end{array} \right)_L \quad \left\| \quad e_R \parallel \nu_R \right.$$

$$Q \nu_R = 0, \quad T_3 \nu_R = 0$$

\Downarrow

$$Y \nu_R = 0$$

\Downarrow

$\nu_R \neq$ usual particles

