Fermionic PEPS

Fermion signs in 2D fermionic tensor networks can be kept track of using two 'fermionization rules'.
[Corboz2009] with Vidal and [Corboz2010b] with Evenbly, Verstrate, Vidal first introduced them, for MERA.
[Corboz2010b] with Orus, Bauer, Vidal adapted them to PEPS context.
This is the approach described in [Bruognolo2020] and presented in this lecture.

Key ingredients: (i) use only positive-parity tensors
(ii) replace line crossings by fermion SWAP gates

Equivalent formulations had also been developed by:
[Barthel2009] with Pineda, Eisert, [Pineda2010] with Barthel, Eisert
[Kraus2010] with Schuch, Verstraete, Cirac
[Shi2009] with Li, Zhao, Zhou
[Bultink2017a] with Williamson, Haegeman, Verstraete, building on [Bultinck2017] (same authors); these papers use the mathematical formalism of 'super vector spaces'.

1. Parity conservation

Fermionic Hamiltonians preserve parity of electron number: $\hat{P} = (-1)^{\hat{N}}$ (1)

$$\hat{H} = \hat{c}^\dagger \hat{c} + \hat{c}^\dagger \hat{c}\, \hat{c}^\dagger \hat{c} + \hat{c}^\dagger \hat{c}^\dagger + \hat{c}\hat{c}\,, \qquad [\hat{H}, \hat{P}] = 0 \tag{2}$$

$\Rightarrow$ all energy eigenstates are parity eigenstate, too, hence may be labeled by parity eigenvalue:

$$\hat{H}|\alpha, p\rangle = E_{\alpha,p}|\alpha,p\rangle\,, \quad \hat{P}|\alpha,p\rangle = p|\alpha,p\rangle\,, \qquad p = \pm \qquad (\text{'}\mathbb{Z}_2\text{-symmetry'}) \tag{3}$$

So, we may agree to work only with states of well-defined parity.

Example: state space of local fermions, $|n_\uparrow, n_\downarrow; p\rangle$ (4)

$$|0\rangle \equiv |0,0;+\rangle\,; \qquad |\uparrow\downarrow\rangle \equiv c_\downarrow^\dagger c_\uparrow^\dagger|0\rangle \equiv |1,1;+\rangle$$
$$|\uparrow\rangle \equiv c_\uparrow^\dagger|0\rangle \equiv |1,0;-\rangle\,, \qquad |\downarrow\rangle \equiv c_\downarrow^\dagger|0\rangle \equiv |0,1;-\rangle\,, \tag{5}$$

Every line in tensor network diagram also carries a parity index.

[When keeping track of abelian symmetries, parity label can be deduced from particle number: $p = (-1)^Q$]

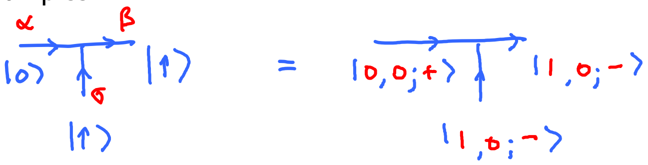
Enforcing $\mathbb{Z}_2$ symmetry [Corboz2010b, Sec.II.F]

To enforce $\mathbb{Z}_2$ symmetry on tensor network: choose all terms to be 'parity preserving'.

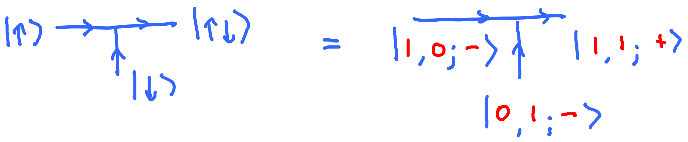
Rule (i): Total parity is positive for all tensors:

n-leg tensor: $A_{\alpha_1 \alpha_2 \dots \alpha_n} = 0$ if $P_{\alpha_1 \alpha_2 \dots \alpha_n} \equiv p(\alpha_1) p(\alpha_2) \dots p(\alpha_n) \neq 1$ (6)

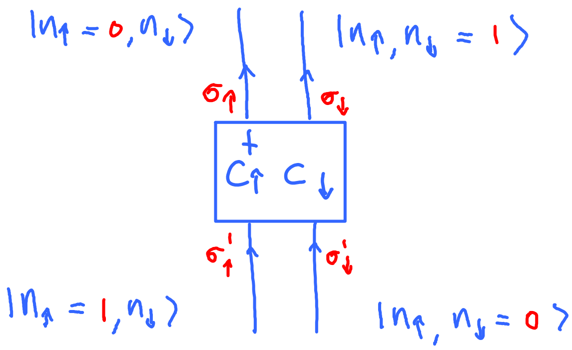
Examples:



$$P^{\alpha\sigma}{}_{\beta} = P_\alpha P_\sigma P_\beta$$
$$= (+)(-)(-) = 1$$



$$P^{\alpha\sigma}{}_{\beta} = (-)(-)(+) = +$$



$$P^{\sigma'_\uparrow \sigma'_\downarrow}{}_{\sigma_\uparrow \sigma_\downarrow} = \underbrace{(P_{\sigma'_\uparrow} P_{\sigma_\uparrow})}_{(-)} \underbrace{(P_{\sigma'_\downarrow} P_{\sigma_\downarrow})}_{(-)} = +$$

$c^\dagger_\uparrow$ and $c_\downarrow$ both change parity by $(-)$

so overall change is $(-)^2 = +$

2. Fermionic signs

$$c_i c_j = - c_j c_i \, , \quad c_i^\dagger c_j^\dagger = - c_j^\dagger c_i^\dagger \, , \quad c_i c_j^\dagger = \delta_{ij} - c_j^\dagger c_i$$

To keep track of these signs, we choose an ordering convention, say $1, 2, \dots, N$, and define:

$$|l_1, l_2, \dots, l_N\rangle = + c_N^\dagger \cdots c_2^\dagger c_1^\dagger \underbrace{|0_1, 0_2, \dots, 0_N\rangle}_{|0\rangle}$$

We have to keep this order in mind when evaluating matrix elements. Example: consider $N = 3$:

$$|\psi\rangle = |0,1,1\rangle = c_3^\dagger c_2^\dagger \mathbb{1}_1 |0\rangle , \qquad |\psi'\rangle = |1,1,0\rangle = \mathbb{1}_3 c_2^\dagger c_1^\dagger |0\rangle$$

$c_2^\dagger c_1^\dagger |0\rangle$

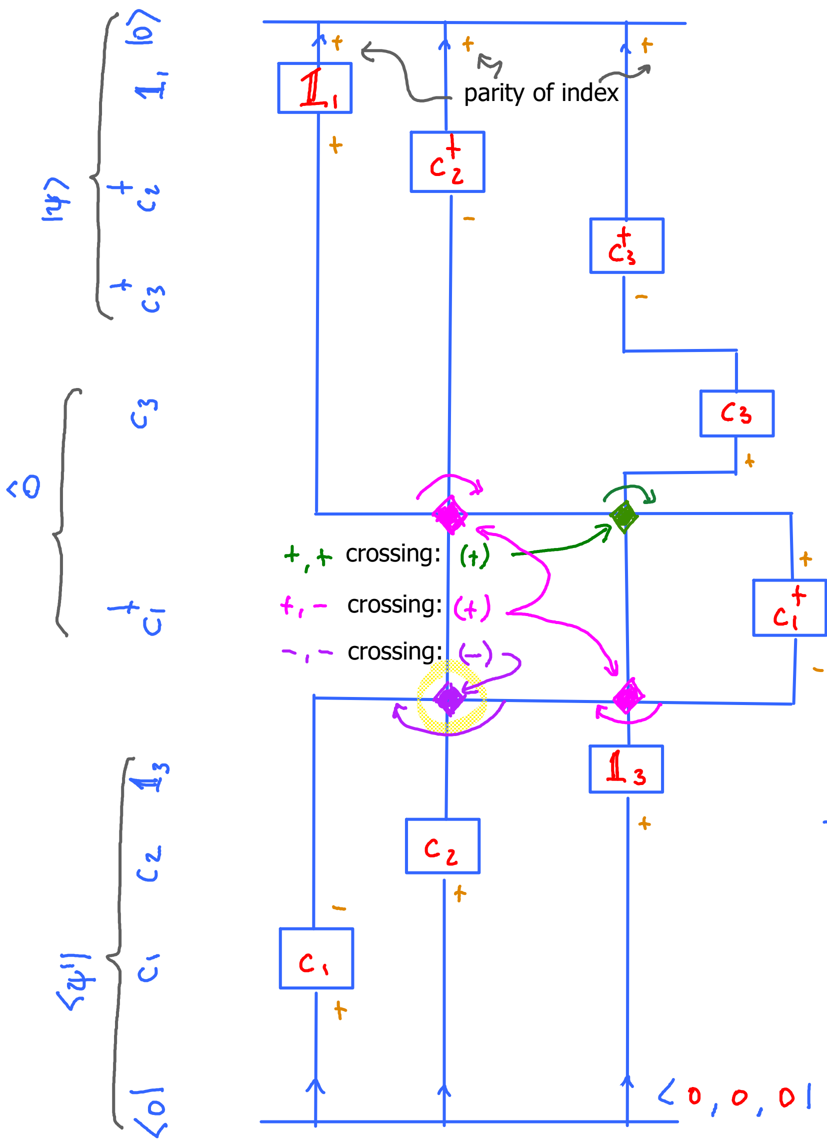
$$\langle \psi' | \underbrace{c_1^\dagger c_3}_{\hat{O}} |\psi\rangle = \langle 0| c_1 c_2 \overbrace{c_1^\dagger \underbrace{c_3 c_3^\dagger}_{=\mathbb{1}_3} c_2^\dagger}^{(-)} |0\rangle = - \langle 0 | \underbrace{c_1 \underbrace{c_2 c_2^\dagger}_{\mathbb{1}_2} c_1^\dagger}_{\mathbb{1}_1} |0\rangle = -1$$

Let us repeat this computation in MPS language: [Corboz2009, App. A]

Order of vertical lines, from left to right, indicates order of operators acting on $|0\rangle$, from right to left.

Horizontal lines show how to move operators in $\hat{O}$ (here $c_1^\dagger c_3$) into appropriate 'slots' in $|\psi\rangle$ or $|\psi'\rangle$.

Line crossings indicate operator swaps. An overall minus sign arises whenever two odd-parity lines cross.



$$|0\rangle = |0,0,0\rangle$$

$$\mathbb{1}_1|0\rangle = |0,0,0\rangle$$

$$c_2^\dagger \mathbb{1}_1 |0\rangle = |0,1,0\rangle$$

$$c_3^\dagger c_2^\dagger \mathbb{1}_1 |0\rangle = |0,1,1\rangle$$

$$\underbrace{c_3 c_3^\dagger}_{=\mathbb{1}_3} c_2^\dagger \mathbb{1}_1 |0\rangle = |0,1,0\rangle$$

$$\mathbb{1}_1 \mathbb{1}_3 c_2^\dagger |0\rangle$$

$$c_1^\dagger \mathbb{1}_3 c_2^\dagger |0\rangle$$

$$- \mathbb{1}_3 c_2^\dagger c_1^\dagger |0\rangle$$

$$- \mathbb{1}_3 \mathbb{1}_3 c_2 c_2^\dagger c_1^\dagger |0\rangle$$

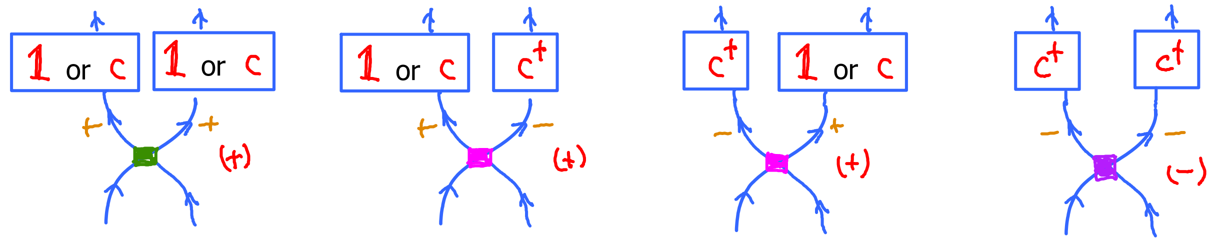
$$- \underbrace{c_2 c_2^\dagger}_{=\mathbb{1}_2} c_1^\dagger |0\rangle = -|1,0,0\rangle$$

$$- \underbrace{c_1 c_1^\dagger}_{=\mathbb{1}_1} |0\rangle = -|0,0,0\rangle$$

SWAP gates

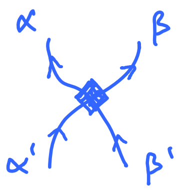
Line crossings keep track of operator orderings.

$(-)$ needed only for exchanging two lines which both host a fermion, i.e. which both have parity $(-)$.



To encode this compactly, introduce SWAP gate whose value depends on parity of incoming lines.

Rule (ii):



$$S^{\beta'\alpha'}_{\alpha\beta} = \delta^{\beta'}_{\alpha}\,\delta^{\alpha'}_{\beta}\,S(\alpha,\beta)$$

$$S(\alpha,\beta) = \begin{cases} -1 & \text{if} \quad p(\alpha) = p(\beta) = (-) \\ +1 & \text{otherwise} \end{cases}$$

Operators

[Corboz2010b, Sec. III.F]

Some matrix elements of operators involving fermions need minus signs.

Example: spinless fermions, consider two sites i, j , with local basis

$$|\sigma_i \sigma_j\rangle = (c_j^\dagger)^{\sigma_j}(c_i^\dagger)^{\sigma_i}|0_i, 0_j\rangle \quad , \quad \sigma_i \in \{0,1\}$$

Two-site operator: $\hat{O} = \sum |\sigma_i', \sigma_j'\rangle\, O^{\sigma_i'\sigma_j'}{}_{\sigma_i\sigma_j}\,\langle \sigma_i\,\sigma_j|$,

with matrix elements $(i<j)$

$$O^{\sigma_i'\sigma_j'}{}_{\sigma_i\sigma_j} = \langle \sigma_i'\sigma_j'|\hat{O}|\sigma_i, \sigma_j\rangle = \langle 0_i 0_j|(c_i^\dagger)^{\sigma_i'}(c_j^\dagger)^{\sigma_j'}\hat{O}\,(c_j^\dagger)^{\sigma_j}(c_i^\dagger)^{\sigma_i}|0_i, 0_j\rangle$$

Examples:

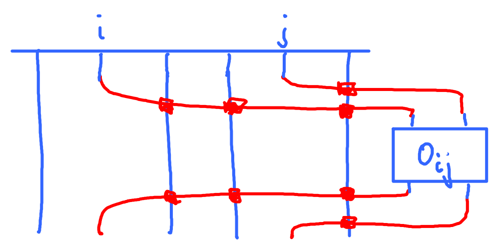
only non-zero element:

Hopping: $\hat{O} = c_i^\dagger c_j$, $\quad O^{1_i 0_j}{}_{0_i 1_j} = \langle 0_i 0_j|\, c_i\, c_i^\dagger\, c_j\, c_j^\dagger\,|0_i 0_j\rangle = +1$

$\hat{O} = c_j^\dagger c_i$, $\quad O^{0_i 1_j}{}_{1_i 0_j} = \langle 0_i 0_j|\, c_j\, c_j^\dagger\, c_i\, c_i^\dagger\,|0_i 0_j\rangle = -1$

Pairing: $\hat{O} = c_j c_i$, $\quad O^{0_i 0_j}{}_{1_i 1_j} = \langle 0_i 0_j|\, c_j c_i\, c_j^\dagger\, c_i^\dagger\,|0_i 0_j\rangle = -1$

When applying such an operator to a generic state, line crossings appear. These yield additional signs, which can be tracked using rule (ii).

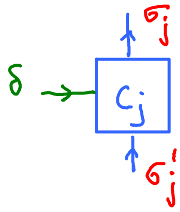


Parity changing tensors

$c^\dagger$ and c change parity; but rule (i) demands: use only parity-conserving tensors!

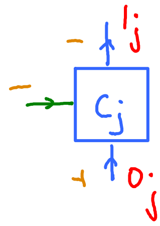
Remedy: add additional leg, with index taking just a single value, $\delta \equiv 1$ with parity $p(\delta) \equiv (-)$

which compensates for parity change induced by $c^\dagger$ or c :



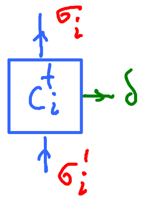
$= (c_j)^{\delta\sigma'_j}{}_{\sigma_j}$

only nonzero element:



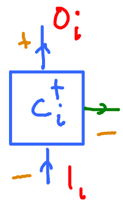
$= \langle 0_j | c_j c_j^\dagger | 0_j \rangle = 1$

Total parity: $P^{\delta\sigma'_j}{}_{\sigma_j} = p(\delta)\,p(\sigma'_j)\,p(\sigma_j) = (-)(+)(-) = (+)$ ✓



$= (c_i^\dagger)^{\sigma'_i}{}_{\sigma_i\delta}$

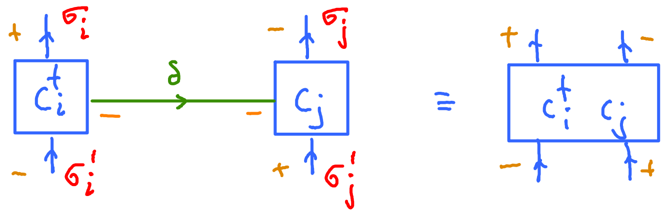
only nonzero element:



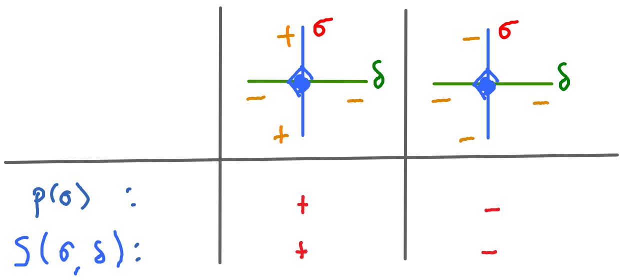
$= \langle 0_i | c_i c_i^\dagger | 0_i \rangle = 1$

Total parity: $P^{\sigma'_i}{}_{\sigma_i\delta} = p(\sigma'_i)\,p(\sigma_i)\,p(\delta) = (-)(+)(-) = (+)$ ✓

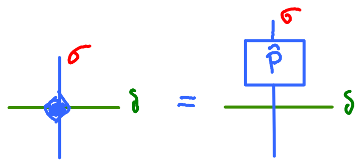
Two-site operator is represented as $c_i^\dagger c_j =$



Since δ carries just a single value, a SWAP gate involving crossing of δ-line and physical σ-line can be simplified to a parity operator acting on latter:



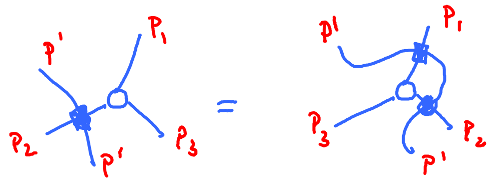
$\Rightarrow$



$\hat{P}(\sigma) = p(\sigma)$

Because all tensors by construction preserve parity, lines can be 'dragged over tensors':

(Shorthand: $p(\sigma_i) = p_i$)



before jump after jump

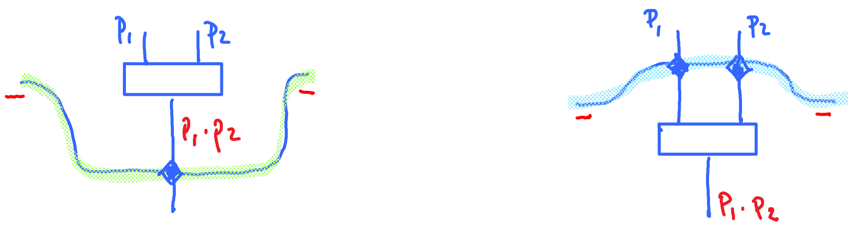
This is trivially true for $p' = (+)$ since then all swap signs are $+$: $S(p_i, +) = (+)$ for all p_i

Consider $p' = (-)$:

2-leg tensor:

SWAP sign: $S(p, -) = S(p, -)$

3-leg tensor:



SWAP sign:

(p_1, p_2)	$S(p_1 \cdot p_2, -)$	$S(p_1, -)\, S(p_2, -)$
$+, -$	$S(-,-) = (-)$	$S(+,-)\, S(-,-) = (+)(-) = (-)$
$-, +$	$S(-,-) = (-)$	$S(-,-)\, S(+,-) = (-)(+) = (-)$
$+, +$	$S(+,-) = (+)$	$S(+,-)\, S(+,-) = (+)(+) = (+)$
$-, -$	$S(+,-) = (+)$	$S(-,-)\, S(-,-) = (-)(-) = (+)$

General argument: parity-preserving tensor has <u>even</u> number of minus-parity lines:

$$(\text{sign})_{\text{before}} \cdot (\text{sign})_{\text{after}} = \prod_{\alpha \in \text{before}} S(p_\alpha, -) \prod_{\beta \in \text{after}} S(p_\beta, -) = (-)^{\text{even}} = (+)$$

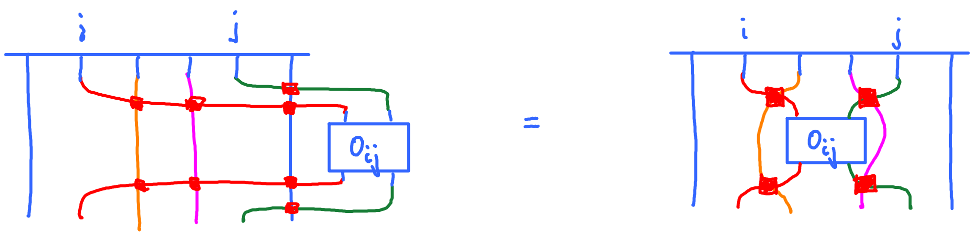
all minus-parity legs cut by 'before' line

all minus-parity legs cut by 'after' line

total number of minus-parity lines, which is even

$\Longrightarrow$ $(\text{sign})_{\text{before}} = (\text{sign})_{\text{after}}$ ✓

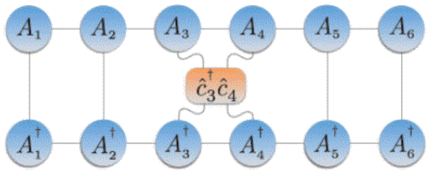
Jump move allows tensor network diagrams to be rearranged according to convenience:



4. MPS examples [Bruognolo2017]

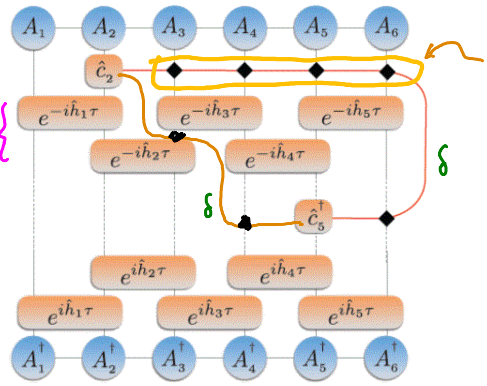
Nearest-neighbor expectation value needs no swap gates:

$\langle\psi|c_3^\dagger c_4|\psi\rangle =$



Time evolution of non-nearest-neighbor hopping operator:

$\langle\psi|e^{iHt} c_5^\dagger e^{-iHt} c_2|\psi\rangle =$

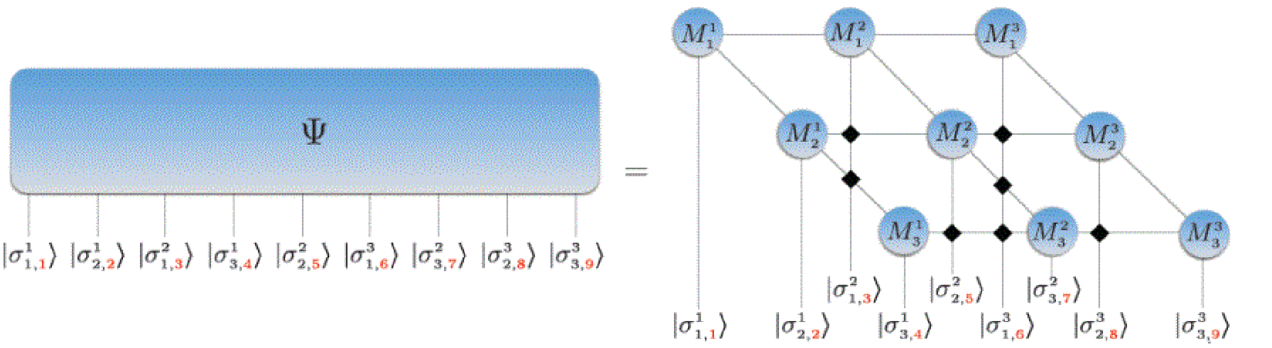


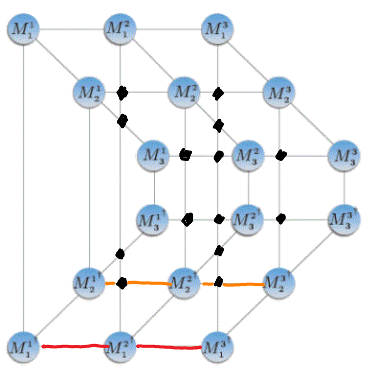
has same effect as Jordan-Wigner string

Due to jump moves, the red line and light brown lines connecting c_2 and $c_5^\dagger$ are equivalent (use one or the other)

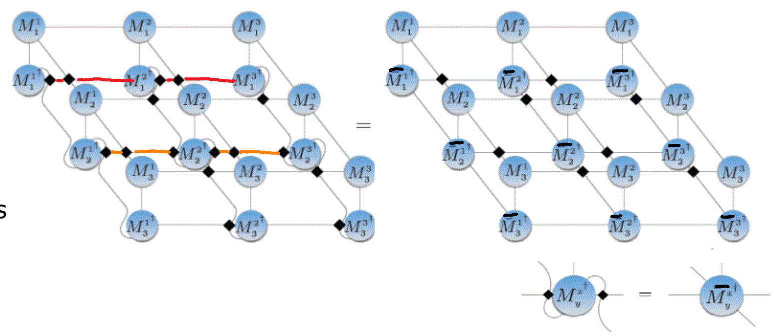
Fermionic order in a PEPS

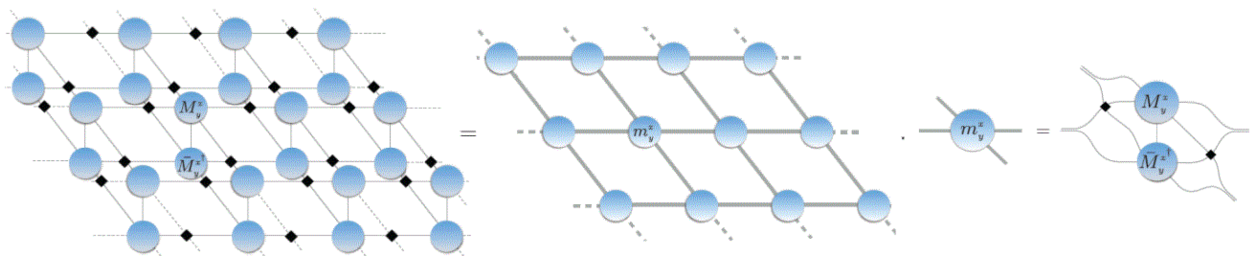
Choose some ordering for open indices and stick to it!



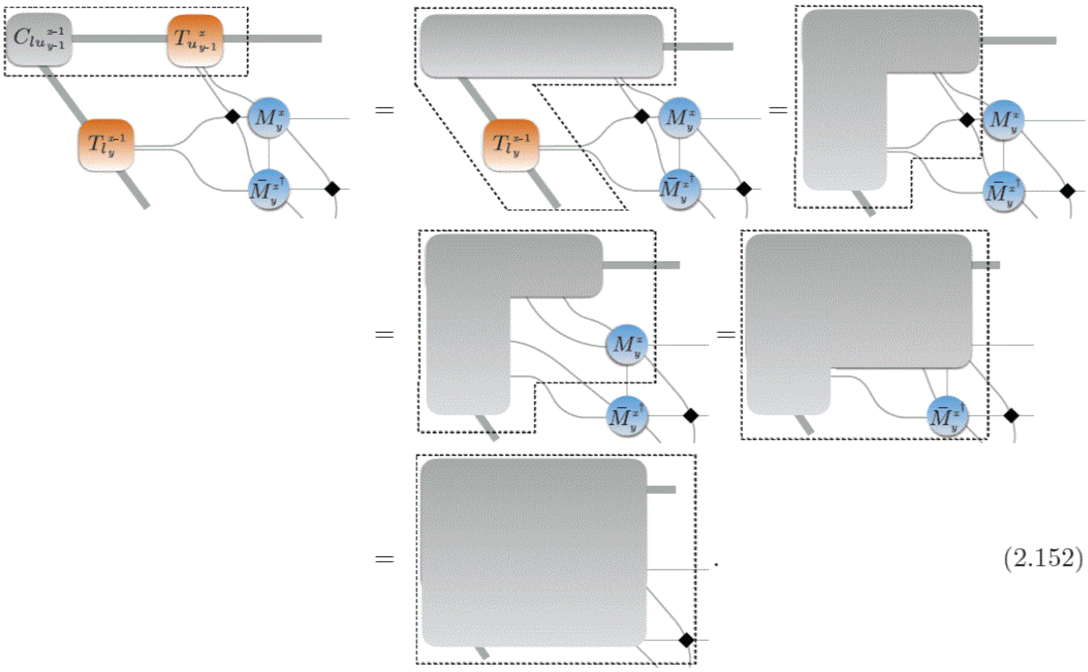


= jump moves





Absorbing SWAP gates



(2.152)