

Problem set 6 (Discussion on July 2)

Problem 1

DNA overstretching transition. Single-molecule stretching experiments in the 1990s revealed that DNA undergoes an overstretching transition if subjected to forces of ≈ 65 pN (Cluzel, *et al.*, *Science* 1996; Smith, *et al.*, *Science* 1996), where it lengthens about 1.7-fold compared to its B-DNA structure. A long-standing debate ensued about what exactly happens upon overstretching. The two possibilities usually considered are DNA melting (i.e. conversion of the double-stranded DNA to two single strands) and conversion of DNA to a double-stranded, but extended and underwound configuration called “S-DNA” (“S” for “stretched”). Van Mameren, *et al.*, *PNAS* 2009, investigated this question using a combination of optical tweezers force-spectroscopy and fluorescence imaging (Available online at <http://www.pnas.org/content/106/43/18231.full.pdf>).

- What do van Mameren, *et al.* conclude about what happens during the overstretching transition, in terms of S-DNA vs. melting?
- What evidence do they provide for their conclusion?
- If we assume that overstretching could involve *both* the formation of S-DNA and DNA melting, how conclusive is their evidence? In particular, does their work rule out the formation of S-DNA upon overstretching?

Problem 2

FJC, revisited. Here we will explicitly derive some identities involving the radius of gyration R_g , which is a very useful measure for the size of a polymer in solution. We assume an ideal FJC (without self-avoidance) with N identical segments of length a . The vectors \vec{r}_i point to segment i . One definition of the radius of gyration is

$$R_g^2 = \frac{1}{N} \sum_{i=1}^N \langle (\vec{r}_i - \vec{r}_{mean})^2 \rangle \quad (1)$$

Where $\langle \dots \rangle$ denotes the statistical average and \vec{r}_{mean} the center of mass position:

$$\vec{r}_{mean} = \frac{1}{N} \sum_{i=1}^N \vec{r}_i \quad (2)$$

a) Show that the R_g can also be expressed as

$$R_g^2 = \frac{1}{2N^2} \sum_{i,j=1}^N \langle (\vec{r}_i - \vec{r}_j)^2 \rangle \quad (3)$$

b) Use the fact that for a FJC

$$\langle (\vec{r}_i - \vec{r}_j)^2 \rangle = |i - j|a^2 \quad (4)$$

($|\dots|$ denotes the absolute value; note that this identity gives the end-to-end distance result for $i = 0$ and $j = N$) to show that for the FJC

$$R_g^2 = \frac{1}{6}Na^2 \quad (5)$$

Hints: Turn the two summations into two integrals, starting from zero. Adjust the integral limits of the inner integral such that the absolute value is taken into account.

Problem 3

3D Gaussian chain. The end-to-end vectors of an ideal chain with N segments of length a in 3D are Gaussian distributed:

$$P(N, \vec{R}) = \left(\frac{3}{2\pi \cdot Na^2} \right)^{3/2} \exp \left(-\frac{3R^2}{2 \cdot Na^2} \right) \quad (6)$$

The mean is $\langle \vec{R} \rangle = 0$ and the variance is $\langle \vec{R}^2 \rangle = Na^2$. The square root of the mean squared radius is interpreted as the “unperturbed end-to-end distance”.

- Use the Boltzmann relation (i.e. the connection between probability and entropy) to calculate the free energy change if an ideal chain is perturbed from an unperturbed end-to-end distance to an arbitrary end-to-end distance R .
- Calculate the force associated with the conformational change in b). *Hint:* You can use that $F = -\partial\Delta G/\partial R$.
- What is the effective “spring constant” of an ideal chain? Compare your result to the result obtained in class by expanding the Langevin function for low forces.