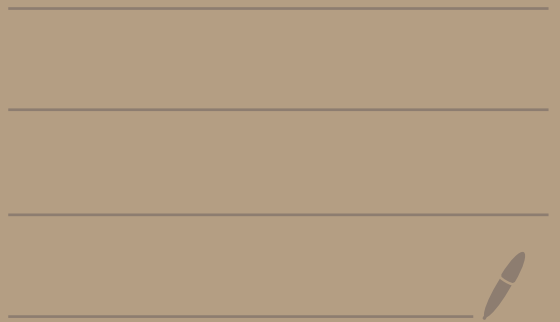


Lecture xxvii



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BB SM Neutrino  
Course



$M_0, M_u, M_d, M_e$

mass matrices in generation  
space

See saw

$$M_\nu = -M_0^T \frac{1}{M_N} M_0$$

$$iM_0 = \sqrt{M_N} O \sqrt{M_\nu}$$

$$O O^T = 1$$

$O \in C$

# LR theory

$$S^u = \frac{v_L}{v_R} - \frac{1}{\sqrt{M_N}} M_N^* \frac{1}{\sqrt{M_N}}$$

$$\boxed{H H^T = H H^* = S} \quad (1)$$

$$M_D = \sqrt{M_N} H \sqrt{M_N^*}$$

Solving (1)

$$\Rightarrow \boxed{S H = H S^*} \quad (2)$$

$\Downarrow$   $2 \times 2$  (2 eqn)

$$H_{2 \times 2} = \sqrt{S S^*} \frac{1}{\sqrt{S^*}}$$

We never solved for  $S^*$

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$$H_{3 \times 3} = ???$$

•  $S = V D V^T$

$$H H^T = S$$

did not work

$$H = \sqrt{S^T}$$

$$V V^T = I$$



$$S = O \Lambda O^T$$

$$O O^T = I$$

$\Lambda =$  Jordan form

is not always  
diagonal

Theorem:

$$U_L M U_R^T = d \text{ (diagonal)}$$

$$U_L U_L^T = U_R U_R^T = I$$

$$(MM^+)^+ = MM^+$$

$$U_L MM^+ U_L^+ = d^2$$

---

$$U_L S U_R^+ = \text{diagonal} = d$$

$$M = S \therefore S = S^T$$

$$\Rightarrow U_R = U_L^*$$

$\Downarrow$

$$V S V^T = d$$

$$M = H \therefore H^+ = H \Rightarrow V H V^+ = d$$

# Jordan

$$S = O \Lambda O^T$$

$$H = O \sqrt{\Lambda} O^* \quad (H = H^T)$$

$$HH^* = S$$

$$HH^* = O \sqrt{\Lambda} O^* O^{*} \sqrt{\Lambda^*} O^T$$

$$= O \sqrt{\Lambda} \sqrt{\Lambda^*} O^T$$

H  
S

does not work

$$H = O \sqrt{\Lambda} \underbrace{(E)}_{\text{correct}} O^T$$

$$H^* = O^* \sqrt{\Lambda^*} E^* O^{*T}$$

↓

$$HH^* = O \sqrt{\Lambda} \underbrace{E O^T O^*}_{I} \sqrt{\Lambda^*} E^* O^{*T}$$

$$= O \sqrt{\Lambda} E \sqrt{\Lambda^*} E^* O^{*T}$$

↓

$$= O \Lambda O^T$$

$$E \therefore \sqrt{\Lambda} E \sqrt{\Lambda^*} E^* =$$

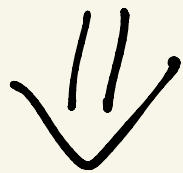
$$= \sqrt{\Lambda} \sqrt{\Lambda} (E^*)^{-1} E^*$$

$$= \Lambda$$

(if  $E$  exists)

•  $HH^* = \textcircled{S}$  constant

$$\left[ \begin{array}{l} \text{Tr } HH^* \in \mathbb{R} \quad (\text{real}) \\ \text{Tr } (HH^*)^2 \in \mathbb{R} \quad -||- \\ \text{Tr } (HH^*)^3 \in \mathbb{R} \quad -||- \end{array} \right]$$



$$\sqrt{\Lambda} E = E \sqrt{\Lambda}^* \quad E^T = E^* = E^{-1}$$



$$\boxed{\Lambda E = E \Lambda^*}$$

$$\boxed{E = E^+}$$

$$\boxed{\text{Take } \Lambda = \text{diagonal}}$$

$$\left. \begin{aligned} \text{Im Tr } \Lambda &= \text{Im Tr } \Lambda^2 = \text{Im Tr } \Lambda^3 = 0 \\ &\Downarrow \end{aligned} \right\}$$

$$(i) \Lambda = (\Lambda_1, \Lambda_2, \Lambda_3) \quad \Lambda_i \in \mathbb{R}$$

$$\Rightarrow E = 1$$

$$(ii) \mathcal{J} = (\Lambda, \Lambda_0, \Lambda^*), \quad \Lambda_0 \in \mathbb{R}$$

$$\mathcal{J}^* = (\Lambda^*, \Lambda_0, \Lambda)$$

$$\Rightarrow E = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix}$$

Summary

measure  
(v oscill)

$$S = \frac{v_L}{v_R} - \frac{1}{\sqrt{M_N}} M_{\nu}^* \frac{1}{\sqrt{M_N}}$$

$(\uparrow)$  HC measure

$$H H^* = S^*$$

$$H_D = \sqrt{M_N} H \sqrt{M_N^*}$$

$$S = O \Delta O^T$$

$\Downarrow$

$$H = O \sqrt{\Delta} E O^t$$

no ambiguity

$\Downarrow$

$$\Theta = \frac{1}{\sqrt{M_N}} H_D$$

particle

$$E^2 = p^2 + m^2$$

NO

$$M_D = i \sqrt{M_N} \text{Osc} \sqrt{M_V}$$

$$\text{Osc} = \text{fixed}$$

QCD

$$m_q = \text{input}$$

$$M_q = V_L^+ m_q V_R$$

4



diagonal, positive

$\mathbb{1} = \text{ambiguous}$

$$\begin{aligned} M_\nu &= m_\nu \mathbb{1} \\ M_N &= m_N \mathbb{1} \end{aligned} \Rightarrow \text{ambiguity}$$

Neutrino masses

$$\Delta m_A^2 = 10^{-3} \text{ eV}^2$$

$$\Delta m_\Theta^2 = 10^{-4} - 10^{-5} \text{ eV}^2$$

$$m_\nu \leq 0.1 \text{ eV}$$

(0.2 eV)

$$m_\nu \approx \frac{1}{25} \text{eV}$$

$$m_\nu \leq \frac{1}{10} \text{eV}$$

$$\theta = \frac{1}{M_N} M_D \leftarrow \text{crax}$$



$$\frac{g}{\sqrt{2}} \bar{N} \gamma_\mu \theta e W_\mu^+$$

$$N \rightarrow e + W$$

$$\Gamma \propto \theta^2$$

$$\Theta \in \mathbb{R} \quad \Theta = \frac{1}{M_N} M_D$$

$$\Theta^T = M_D^T \frac{1}{M_N}$$

$$\Theta \Theta^T = \frac{1}{M_N} M_D M_D^T \frac{1}{M_N}$$

up to signs  $\Rightarrow$

no ambiguity

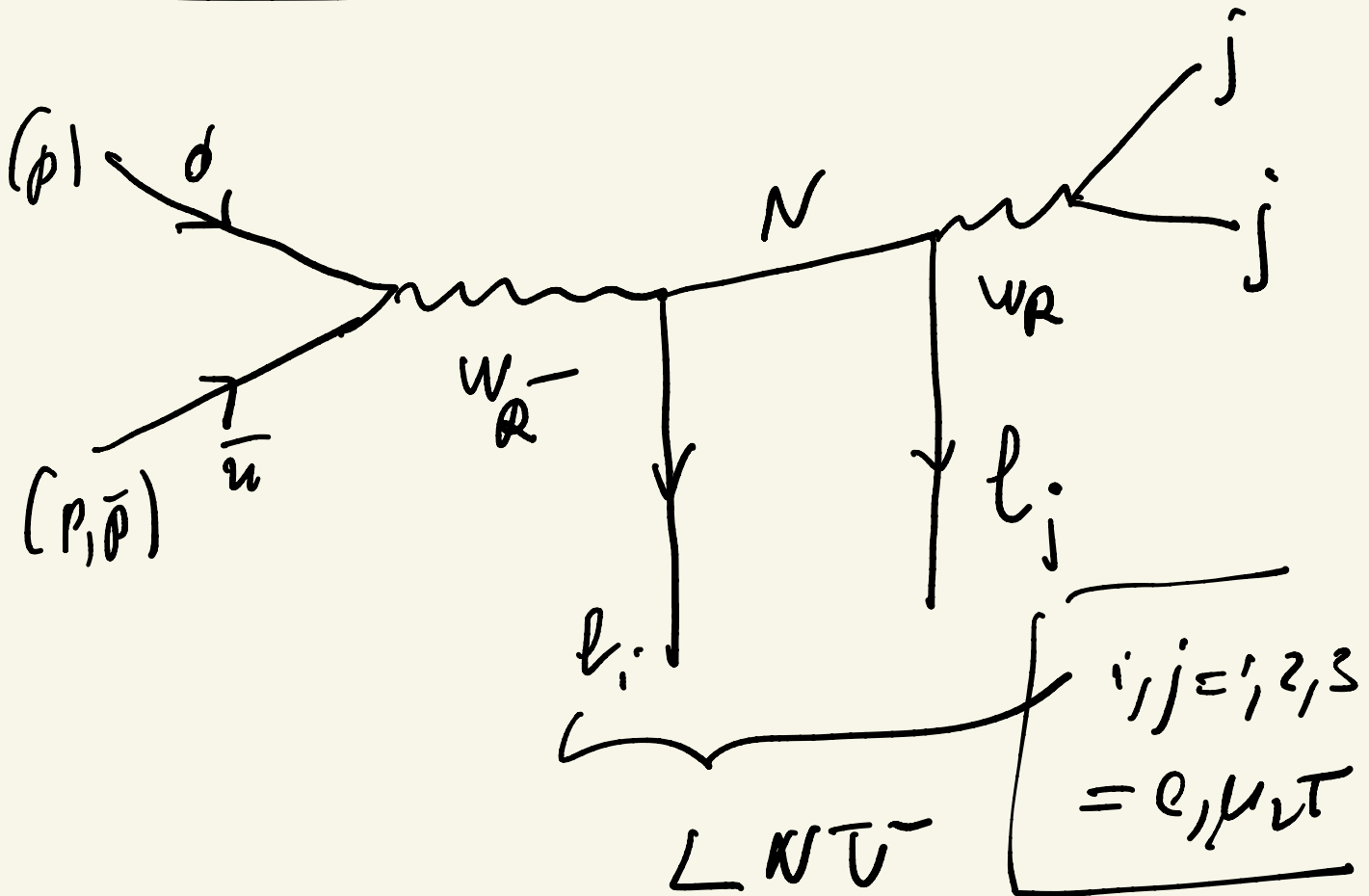
$$M_D \approx (\sqrt{M_{11}} \quad \sqrt{M_{22}})$$

( $v_1 = 0$ )

" see row gets solved "

# Summary of physics of probe of neutrino mass

(i) find New probe  $M_N$



$$N \rightarrow l + \nu_1 + \nu_2$$

$$N \rightarrow \bar{e} + j_1' + j_2'$$

$$\Rightarrow \begin{cases} 50\% = e \\ 50\% = \bar{e} \end{cases}$$

Probe of Majorana  
nature

$$\underline{M}_\nu = V_R m V_R^T$$

lepton mixings

$V_R = \sim 11 \sim$  in gen space

**LFV**

Lepton flavour violation

↓  
(a) neutrino oscillation

(ii)  $\mu \rightarrow e \gamma, \mu \rightarrow e e e \bar{e}$

$\mu + N \rightarrow e + N$

what happens in SM?

why is it suppressed?

$$B(\mu \rightarrow e \gamma) = \frac{\Gamma(\mu \rightarrow e \gamma)}{\Gamma(\mu \rightarrow e \nu \bar{\nu})} \leq 10^{-13}$$

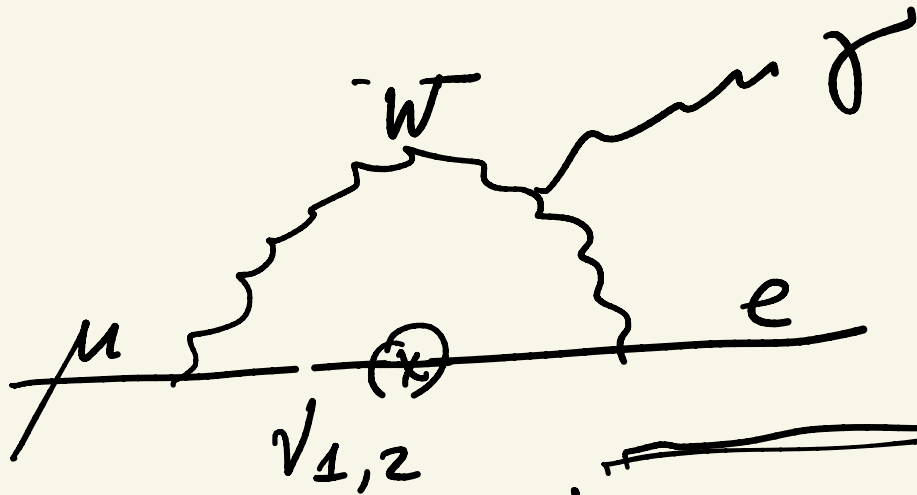
$$B(\mu \rightarrow e e e \bar{e}) = \frac{\Gamma(\mu \rightarrow e e e \bar{e})}{\Gamma(\mu \rightarrow e \nu \bar{\nu})} \leq 10^{-13}$$

$$\mathcal{B}(\mu \rightarrow e \gamma)_{SM} \stackrel{?}{=} ?$$

• we are out of  $\nu \Rightarrow$  oscillation

$$(m_\nu \neq 0)$$

$\Rightarrow$  LFV  $\therefore \mu \rightarrow e \gamma$



$$\mathcal{B}(\mu \rightarrow e \gamma) \approx \frac{\alpha}{\pi} \left( \frac{\Delta m^2}{M_W^2} \right)^2 \text{ mixing}$$

$$\Delta m^2 \approx 10^{-3} \text{ eV}^2 \approx 10^{-3} \cdot 10^{-18} \text{ GeV}^2$$

$\Downarrow$

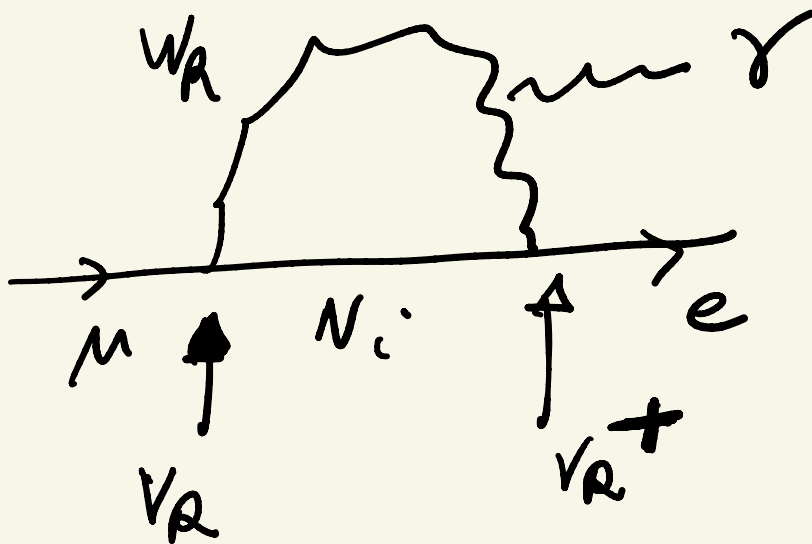
$$\frac{\Delta m^2}{M_{Pl}^2} = 10^{-21} / 10^4 \approx 10^{-25}$$

$$\Rightarrow B(\mu \rightarrow e \gamma) \approx 10^{-50}$$

"SM"

"SM" = SM + neutrino mass

$$B_{LR}(\mu \rightarrow e \gamma) \approx \frac{\alpha}{\pi} \left( \frac{M_{WL}}{M_{WR}} \right)^4 \frac{\Delta m^2}{M_{WR}^2}$$





$$l_i \rightarrow l_j + \gamma$$

$$M_N - M_N^+$$

LHC

$$M_{W_R} \approx 4 \text{ TeV}$$

$$M_{W_R} \approx 8 \text{ TeV (limit)}$$

$$M_{W_L} \approx 80 \text{ GeV}$$

$$\Rightarrow M_{W_R} \approx 100 M_{W_L}$$

$$B(\mu \rightarrow e\gamma) \approx \frac{\alpha}{\pi} \cdot 10^{-2} \cdot 10^{-8}$$

$$\frac{\Delta \kappa_N^2}{M_{W_R}^2}$$

$$GIM \text{ in SM} \Rightarrow \frac{m_c^2}{M_W^2} \leq 10^{-3}$$

$$\mu \rightarrow e e \bar{e}$$

$$\delta^{++}$$

also enters

$$\mathcal{L}_Y = l_R^T \Delta_R Y_R l_R + h.c.$$

$$M_N^* = Y_R \langle \Delta_R \rangle$$

$$e_i (M_N^*)_{ij} l_j \delta^{++}$$



$$\delta^{--} \rightarrow l_i l_j (\alpha M_N^{ij})$$

Tello, G.S. '2020

# Correlations

LHC

$0\nu 2\nu$ , LFV ( $\mu \rightarrow e \gamma$ ...)

$\nu$  oscillations,  $m_\nu$  (KATRIN)

LFV,  $0\nu 2\nu$   $\Rightarrow$  new physics

• LFV

$$\text{LR } B(\mu \rightarrow e) \approx \frac{\alpha}{\pi} \left( \frac{M_L}{M_R} \right)^4 \frac{\Delta m_N^2}{M_R^2}$$

$$\begin{aligned} & \parallel \\ 10^{-14} & = \frac{\alpha}{\pi} 10^{-1} \left( \frac{M_L}{M_R} \right)^4 & \parallel \\ & & \downarrow \\ & & O(1/10) \end{aligned}$$

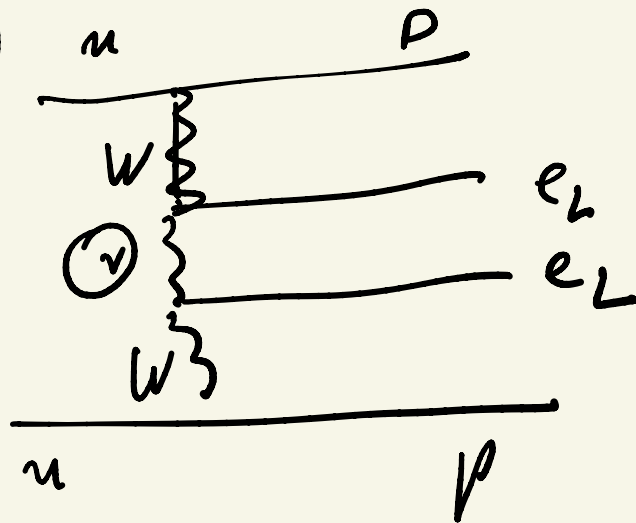
$$\approx 10^{-3} \left( \frac{M_L}{M_R} \right)^4$$

$$\left(\frac{M_L}{M_R}\right)^4 \approx 10^{-11}$$

$$M_R \approx 10^3 M_L \dots$$

OV2/s

"SM"



$$\propto G_F^2 \left(\frac{1}{M_L^4}\right) \frac{M_V^4}{p^2}$$

$$W_R : \propto G_F^2 \left(\frac{M_L}{M_R}\right)^4 \frac{M_N}{\cancel{p^2 - M_N^2}}$$

$$\approx G_F^2 \left(\frac{M_L}{M_R}\right)^4 \frac{1}{M_N}$$

"SM"

$$G_F^2 \frac{10^{-10} \text{ GeV}}{(p=100 \text{ MeV})^2} \approx G_F^2 10^{-8} \frac{1}{\text{GeV}}$$

LR

$$G_F^2 \left(\frac{M_L}{M_R}\right)^4 \frac{1}{M_{\nu\nu}} \approx G_F^2 10^{-8} \frac{1}{\text{GeV}}$$

$$\left(\frac{M_L}{M_R}\right)^4 \frac{1}{M_{\nu\nu}} \approx 10^{-8} \frac{1}{\text{GeV}}$$

tight constraint

$\epsilon_R$  !!!

$M_N$  @ LHC

( $N = \text{MeV}$ )



$$\left\{ \begin{array}{l} -M_D = i \sqrt{\mu_N^2} \sqrt{M_V} \\ \theta = \frac{1}{\mu_N} M_D \end{array} \right. "$$